

Futures in Accelerator-Based Physics

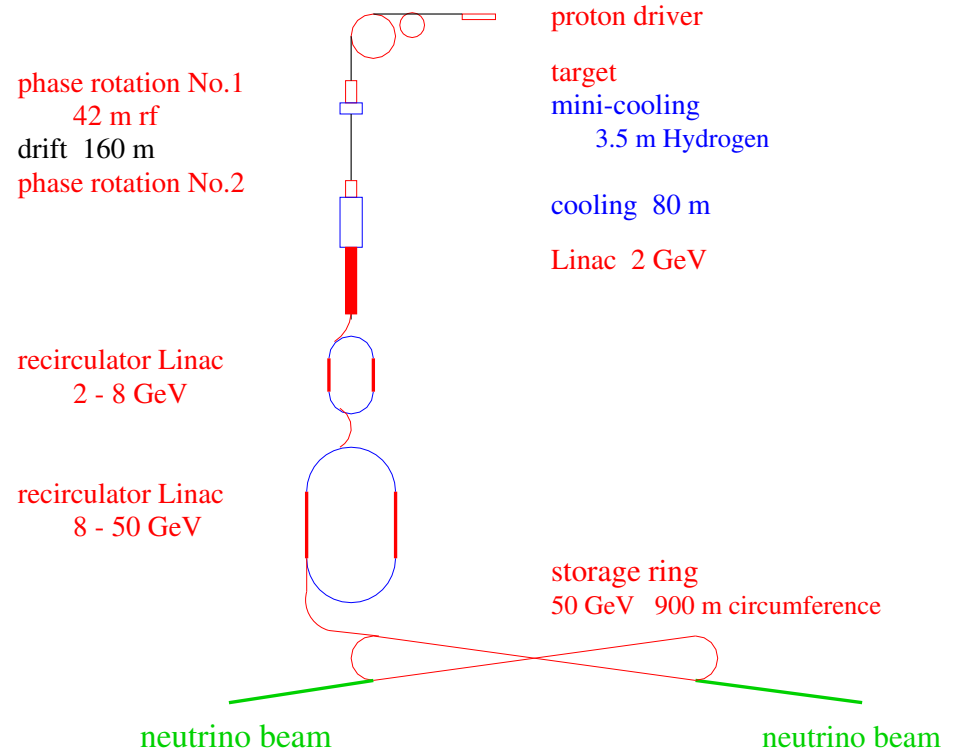
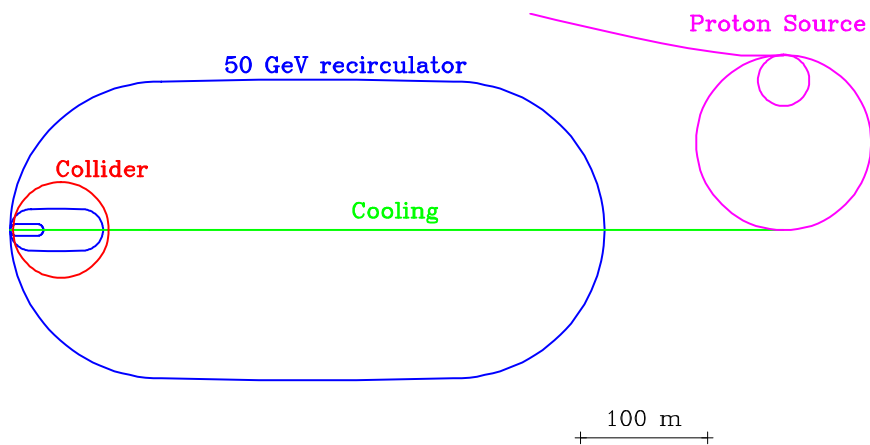
K.T. McDonald

Princeton U.

U. Ioannina, Greece

May 21, 2004

<http://puhep1.princeton.edu/~mcdonald/accel/>



From the Cathode Ray Tube to a Muon Collider

1897 Braun: Cathode ray tube (~ 20 keV,television sets).

1930 Cockcroft & Walton, Van de Graaff, Wideroe: Linear accelerators to ~ 1 MeV.

1930 Lawrence: Cyclotron = first circular accelerator, ~ 10 MeV.

1940 Kerst: Betatron = circular induction accelerator, ~ 10 MeV.

1945 McMillan, Veksler: Synchrotron = circular ring accelerator, ~ 1 GeV.

1950 Christofilos: Alternating gradient focusing, \Rightarrow synchrotrons to 10 GeV.

1955 Kerst: Proton colliding rings (first built as electron rings).

1956 Wigner: Liouville's theorem applies to particle beams.

1956 O'Neill: Ionization cooling for storage rings.

1960 Melissinos: Muon storage ring.

1962 Milburn: Laser-electron Compton backscattered photon beam.

1966 Budker: Electron cooling.

1968 Van der Meer: Stochastic cooling of antiproton rings.

1968-69 Tikhonin, Budker: Muon colliding rings with ionization cooling.

1993 Mikhailichenko, Zolotarev: Optical stochastic cooling.

Nicholas Christofilos



Developed the concept of alternating gradient focusing (inverted pendulum stability) in 1950 while working as an elevator operator.

Later invented the induction linac.

Pioneered studies of electromagnetic pulses and aurorae producing by high altitude explosions of nuclear bombs.

Tribute by A.C. Melissinos: <http://puhep1.princeton.edu/~mcdonald/accel/christofilos.pdf>

United States Patent Office Patented Feb. 28, 1956 2,736,799

2,736,799
FOCUSING SYSTEM FOR IONS AND ELECTRONS
 Nicholas Christofilos (or Philos), Athens, Greece
 Application March 10, 1950, Serial No. 148,920
 8 Claims. (Cl. 250-27)

The present invention relates to a new focusing system for ions and electrons and application thereof in particle accelerators. A major problem in the design of particle accelerators is the provision of suitable means for focusing the accelerated particles towards a predetermined orbit and compensating the mutual electrostatic repulsive forces. An ideal focusing system must accelerate the moving particles towards a predetermined orbit from all directions and the focusing forces must increase as the distance from said orbit increases. If we consider an orthogonal coordinate system, x, y, z , and suppose that the particle's orbit coincides with the x -axis and considering as P_x, P_y, P_z the x, y, z , components of the focusing forces, then, in an ideal focusing system the equations of the P_x, P_y, P_z would be

$$\begin{aligned} P_x &= 0 & (1) \\ P_y &= -\epsilon y & (1a) \\ P_z &= -\epsilon z & (1b) \end{aligned}$$

But simultaneously the Laplace equation $\Delta\psi=0$ must be satisfied so that it must be

$$\frac{\partial^2 P_x}{\partial x^2} + \frac{\partial^2 P_y}{\partial y^2} + \frac{\partial^2 P_z}{\partial z^2} = 0 \quad (1c)$$

or

$$\epsilon_y = -\epsilon_z \quad (1d)$$

From the above equations it is shown that a focusing field capable to accelerate ions or electrons towards a predetermined orbit from all directions simultaneously is impossible. Therefore the focusing system proposed herein is based in a new principle, namely: If, along a predetermined orbit of ions or electrons an electrostatic or electromagnetic field is produced by means of suitably arranged conductors (connected to a high voltage source or energized by high intensity current) exerting on the moving, along said orbit, particles (ions or electrons) forces directed normally to said orbit and varying periodically, in direction and magnitude along said orbit, and increasing in magnitude as the distance from said orbit increases, then the mean value of the focusing forces is negative (directed towards the orbit) and the particles are focused towards the orbit from all directions. The focusing forces, acting on the particles, resulting from the field which is produced electrostatically or electromagnetically, increase as the distance from the orbit increases. The particles move at some finite distance from the orbit and in a direction substantially parallel to the orbit by virtue of the periodically varying exciting focusing forces due to the field. The particles undergo forced oscillations and are subject to the alternately converging and diverging forces from the field. The electrically produced force field, electromagnetic or electrostatic, imposed upon the orbit and the path of the particles exerts forces on the particles within a plane whose normal is substantially parallel to the velocity vector of each of the particles. The path of the particles becomes concave towards the orbit in a converging section and convex to-

wards the orbit in the diverging section. Since the forces are greater as the distance from the orbit becomes greater, the mean value of the converging and diverging forces along the converging section is greater than the mean value of the forces along the diverging section. The resultant force and net effect of the mean value of these alternating forces causes the particles in the path to be forced towards the orbit from all directions and focusing is thereby obtained.

10 In a focusing system based upon this principle the x, y, z , components of the focusing forces are

$$P_x = 0 \quad (2)$$

$$P_y = -\epsilon y \sin \frac{2\pi x}{\lambda} \quad (2a)$$

$$P_z = \epsilon z \sin \frac{2\pi x}{\lambda} \quad (2b)$$

As it is obvious

$$\Delta\psi = \frac{\partial^2 P_x}{\partial x^2} + \frac{\partial^2 P_y}{\partial y^2} + \frac{\partial^2 P_z}{\partial z^2} = 0 \quad (1d)$$

so that the Laplace equation is satisfied and therefore the production of such a field is possible. If we consider a particle moving parallel to the x -axis and at a distance $x=0, y=0$ the force exerted on said particle is

$$P_x = \epsilon \cdot \epsilon_y \sin \frac{2\pi x}{\lambda} \quad (3)$$

As the force P_x varies periodically as the particle moves along the x -axis, said particle undergoes forced oscillations of frequency

$$f = \frac{v}{\lambda} \quad (4)$$

where v is the velocity of the particle.

The result of these oscillations is that the distance from the orbit oscillates around the mean value z_0 according to the equation

$$z = z_0 \left(1 - \mu \sin \frac{2\pi x}{\lambda} \right) \quad (5)$$

where

$$0 < \mu < 1 \quad (6)$$

In the region where

$$\sin \frac{2\pi x}{\lambda}$$

is negative the mean value of the distance from the orbit is greater than z_0 while in the region where

$$\sin \frac{2\pi x}{\lambda}$$

is positive the mean value of the distance is less than z_0 , so that the mean value of the force in the first region is greater than the mean value in the second region, with the result that the mean value of the force in a length λ is negative, focusing the particle towards the x -axis, from all directions.

If the maximum value of the force is

$$P_{max} = \epsilon \cdot \epsilon \cdot \epsilon_0 \quad (7)$$

then the mean value P_m is

$$P_m = \epsilon \cdot \epsilon \cdot \epsilon_0 \quad (8)$$

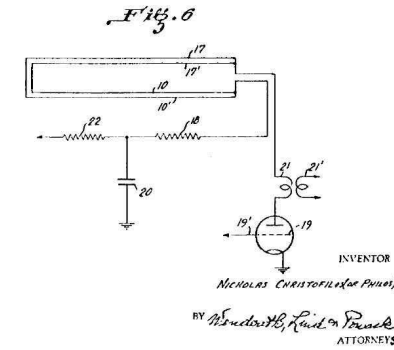
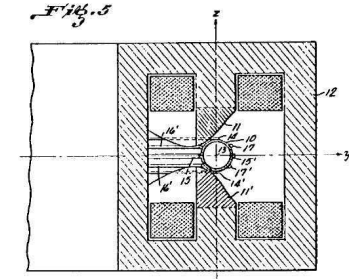
where

$$\epsilon = \epsilon \cdot \frac{\mu}{2} \quad (9)$$

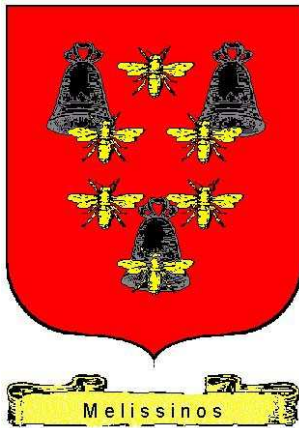
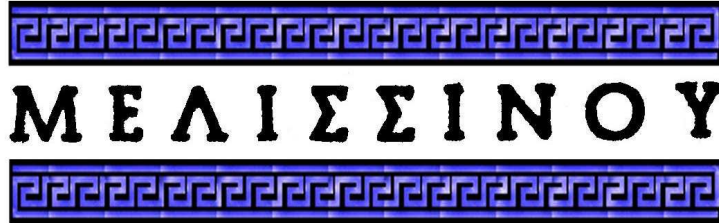
and

$$\mu = \frac{\epsilon \lambda^2}{4 \cdot v^2 \cdot |P|} \quad (10)$$

Feb. 28, 1956 2,736,799
 NICHOLAS CHRISTOFILOS (OR PHILOS)
 FOCUSING SYSTEM FOR IONS AND ELECTRONS
 Filed March 10, 1950 4 Sheets-Sheet 2



Adrian Melissinos



1960: First proposal for a muon storage ring. Later implemented for studies of $g - 2$ of the muon.

1997: Production of positrons in nonlinear light-by-light scattering (with KTM).
“Sparking the Vacuum”.

Tribute by KTM: <http://puhep1.princeton.edu/~mcdonald/adrianfest/adrianfesttrans.pdf>

The University of Rochester
Rochester 20, N.Y.

1960

Department of Physics

It is proposed to build a strong focusing ring to contain μ -mesons in a given momentum band for several τ -meson lifetimes and then eject them.

Orbit radius 1.9 m
Mean radius 2.2 m

Aperture width 7 cm
Aperture height 4 cm

Maximum field 13 Kg.

Field index $n_1 - n_2 - n \approx 15$

Betatron osc freq. $Q_x = 2.25$
 $Q_y = 1.75$

Rotation frequency ≈ 22 Mc/sec

Weight Fe 10 tons Cu 2 tons

Sectors: 4 each double
Focusing order: $1/2F$ D $1/2F$

At the Cosmotron, a fairly well focused π -beam of 10^7 /pulse could be obtained for a $\Delta p/p \approx 20\%$. Assuming that for such a beam the capture efficiency is 50%, the μ -meson beam becomes

$$2.2 \times 10^4/\text{pulse} \quad \text{with } \Delta p/p = 10\%$$

At the AGS one could expect a fairly well focused beam of π 's 2×10^8 so that under the same considerations the μ -mesons are 4.4×10^5 /pulse with $\Delta p/p = 10\%$

The cost is estimated to less than \$100,000.

If all this does not seem too unreasonable I will proceed to calculate orbits.

The Sandal-Making Poet of Athens

“Take away the Glories and the Honors
The granite palaces of this vain world
And only give me the smile of Pain
The tear of Joy and I will erect
A thousand palaces in me in which to live.”

—Stavros Melissinos,

Known among sandal-makers as “The Poet”
(and among poets as “the Sandal-Maker”).

89 Pondrossou Street in Monastiraki, Athens



Melissos of Samos

Student of Parmenides.

Admiral of Samian navy, defeated Pericles 441 BC.

“Nor is anything empty. For what is empty is nothing.
What is nothing cannot be.”

“Accordingly, being was not generated, nor will it be
destroyed;

So it always was and always will be.”



Why Muons?

- **A muon is a heavy electron.**
⇒ Fundamental interest in the properties of the muon and of its decays.
- **Muons live 2.2 μs when at rest.**
⇒ Muons of any energy live \approx 1,000 turns in a 2-T magnetic field.
⇒ Can use rings to accelerate, store and collide muons.

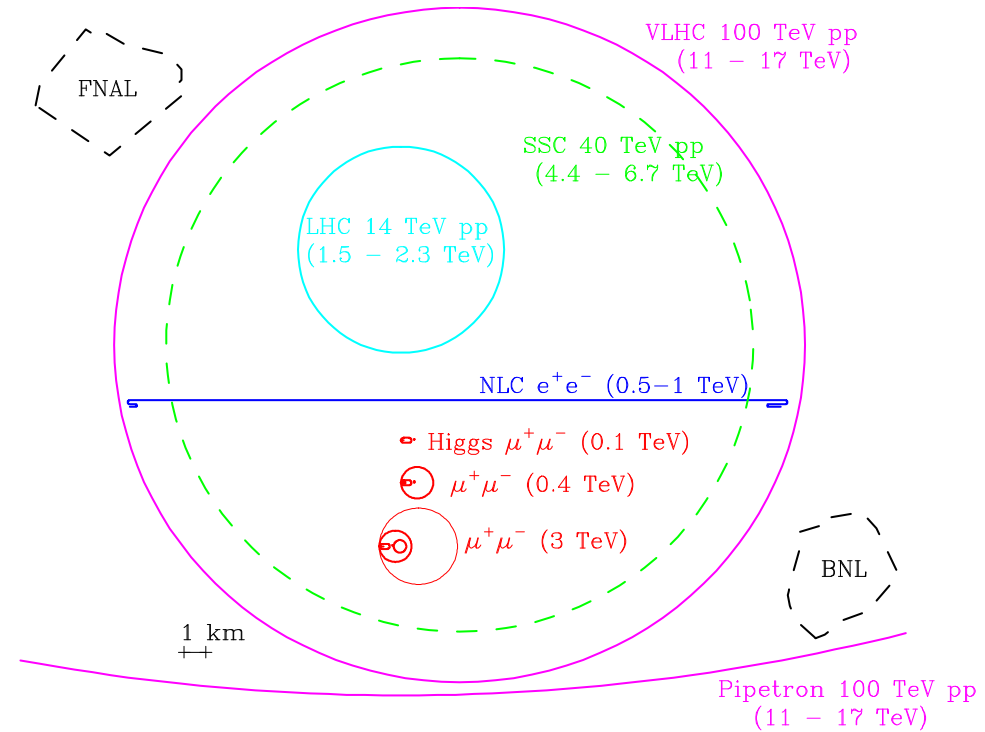
Why Now?

- $m_\mu = 205m_e$
⇒ **Initial state radiation suppressed in $\mu^+\mu^-$ collisions.**
⇒ Precision leptonic initial states up to 100 TeV.
- **Muon decay, $\mu \rightarrow \nu_\mu e \bar{\nu}_e$, provides well-known fluxes of $\nu_\mu, \bar{\nu}_e$ ($\bar{\nu}_\mu, \nu_e$) in equal amounts.**
⇒ **Neutrino factory.**

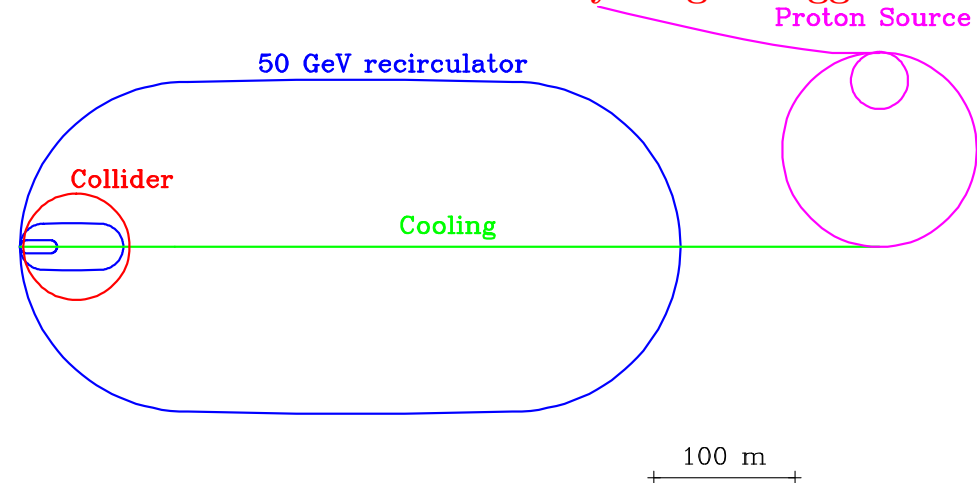
What is a Muon Collider?

An accelerator complex in which

- Muons (both μ^+ and μ^-) are collected from pion decay following a pN interaction.
- Muon phase volume is reduced by 10^6 by ionization cooling.
- The cooled muons are accelerated and then stored in a ring.
- $\mu^+\mu^-$ collisions are observed over the useful muon life of ≈ 1000 turns at any energy.
- Intense neutrino beams and spallation neutron beams are available as byproducts.

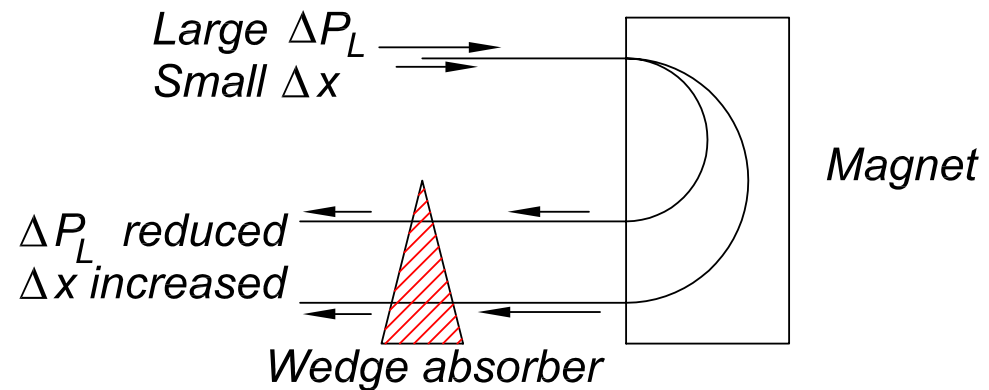
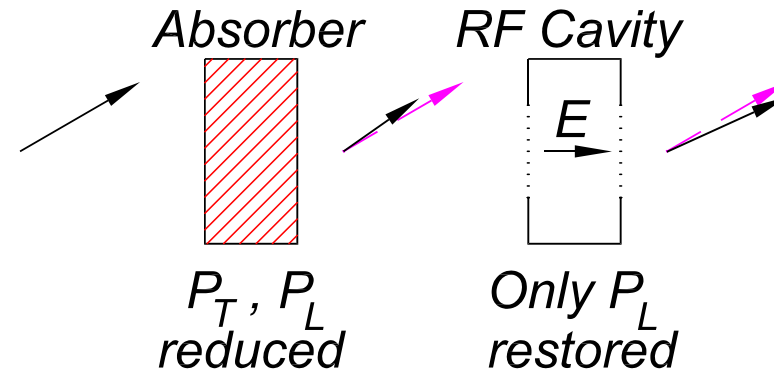


A First Muon Collider to study a light-Higgs



Fast Ionization Cooling of Muon Beams

- dE/dx loss cools both P_T and P_L .
- Multiple scattering heats P_T , straggling heats P_L .
- With low- Z absorber can have net cooling of P_T , but P_L is heated.
- A magnet + wedge absorber can exchange transverse and longitudinal phase space.
- Then cool transversely again....



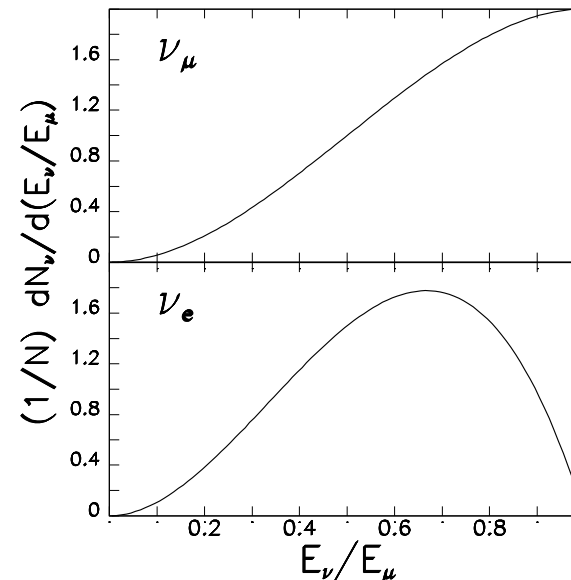
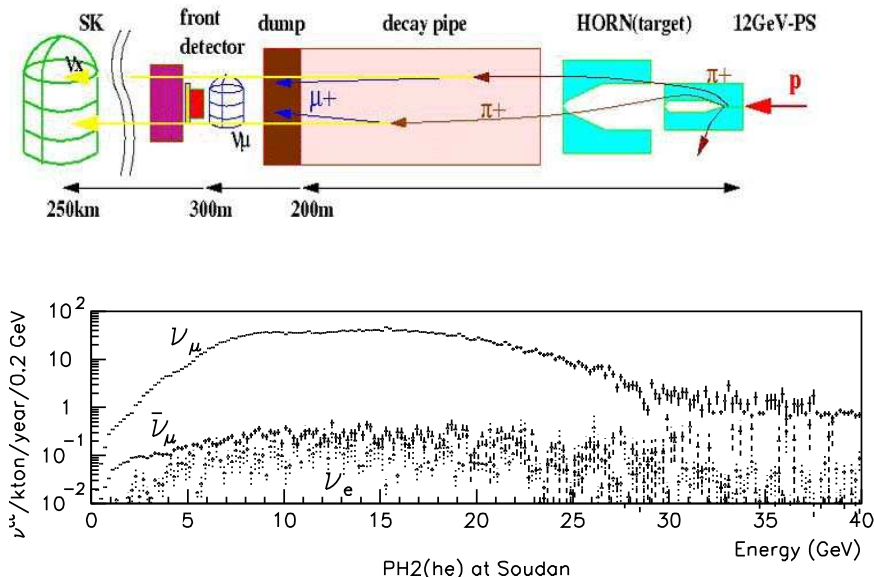
An affordable implementation of fast ionization cooling is the major R&D challenge for a muon collider.

The Neutrino Factory and Muon Collider Collaboration:

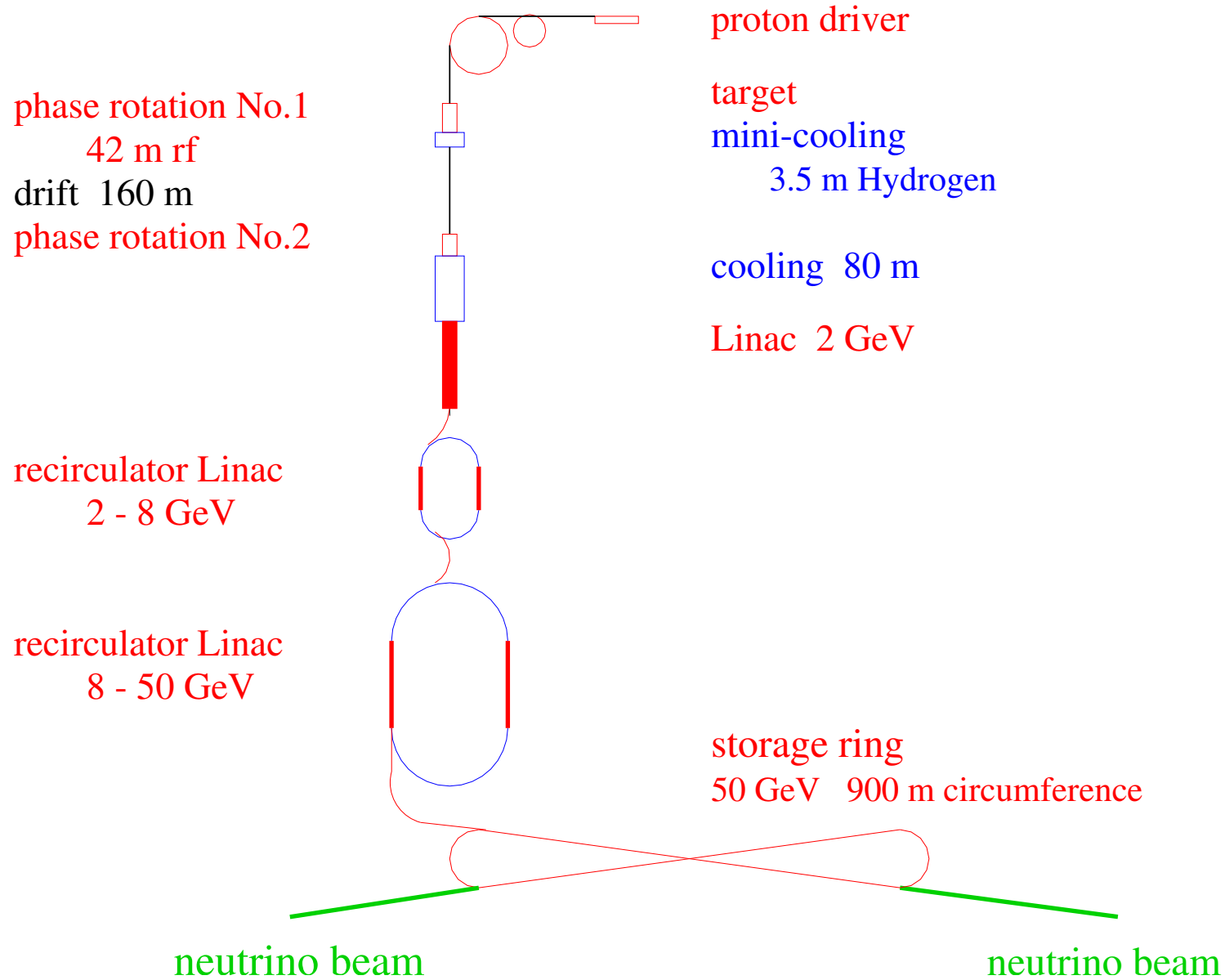
http://www.cap.bnl.gov/mumu/mu_home_page.html

A Neutrino Factory Based on a Muon Storage Ring

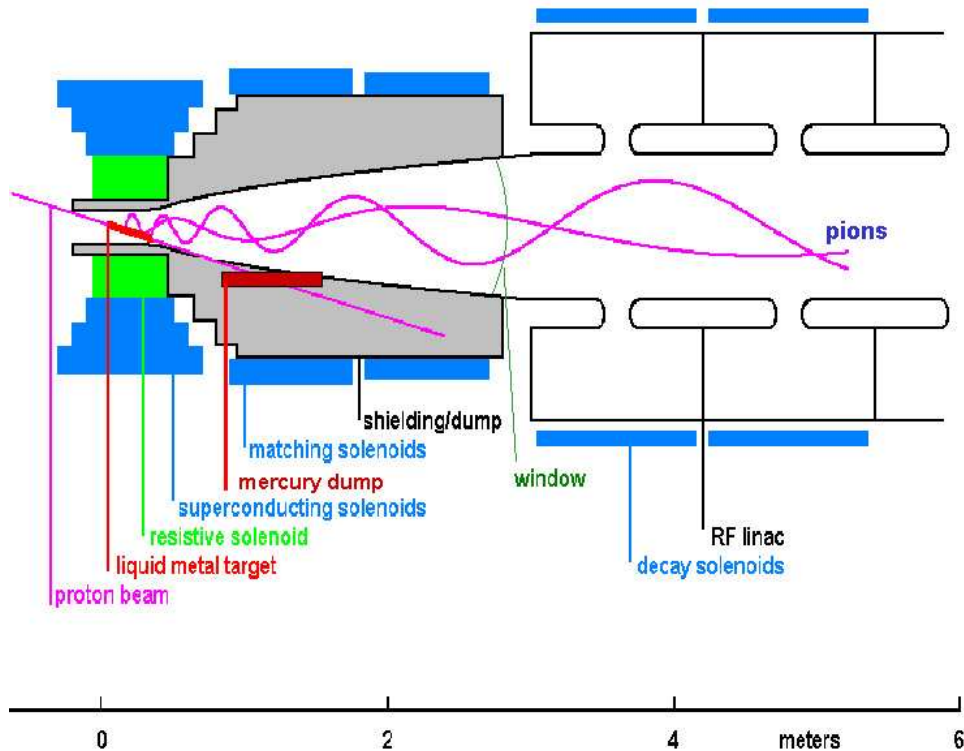
- Many of the neutrino oscillation solutions permit study of the couplings between 2, 3, and 4 neutrinos in accelerator based experiments.
- More neutrinos are needed!
- Present neutrino beams come from $\pi, K \rightarrow \mu\nu_\mu$ with small admixtures of $\bar{\nu}_\mu$ and ν_e from μ and $K \rightarrow 3\pi$ decays.
- Higher (per proton beam power) and better characterized neutrino fluxes are obtained from μ decay.
- Collect low-energy μ 's from π decay, Cool the muon bunch, Accelerate the μ 's to the desired energy, Store them in a ring while they decay via $\mu^- \rightarrow e^- \nu_\mu \bar{\nu}_e$. [Of course, can use μ^+ also.]



A Neutrino Factory Based on a Muon Storage Ring



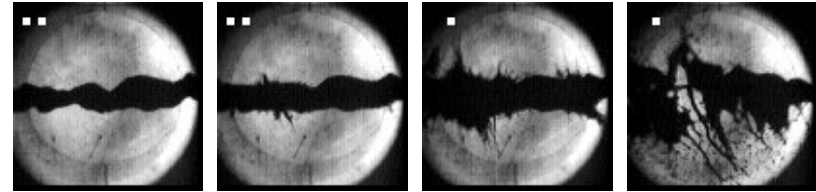
Multimawatt Target for a Muon Collider/Neutrino Factory



Use a free mercury jet as target for a 4-MW proton beam to produce π 's, μ 's and ν 's.

Capture the charged π 's in a 20-T solenoid magnet.

Targetry R&D:



Mercury jet dispersed by a proton beam.



Nick Simos performing solid target studies in a BNL hot cell.

Can Study CP Violation of Neutrinos at $L/E = (2n + 1)500 \text{ km/GeV}$

[Marciano, hep-ph/0108181, Diwan *et al.*, hep-ph/0303081]

The n th maximum of ν_2 - ν_3 oscillations occurs at $L/E \approx (2n + 1)400 \text{ km/GeV}$.

The CP asymmetry grows with distance:

$$A = \frac{P(\nu_\mu \rightarrow \nu_e) - P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)}{P(\nu_\mu \rightarrow \nu_e) + P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)} \approx \frac{2s_{12}c_{12}c_{23}\sin\delta}{s_{23}s_{13}} \left(\frac{\Delta m_{12}^2}{\Delta m_{23}^2} \right) \frac{\Delta m_{23}^2 L}{4E_\nu}$$

$$\Rightarrow \frac{\delta A}{A} \approx \frac{1}{A\sqrt{N}} \propto \frac{E_\nu}{L\sqrt{N}} \approx \text{independent of } L \text{ at fixed } E_\nu.$$

$N_{\text{events}} \propto 1/L^2$,

\Rightarrow Hard to make other measurements at large L .

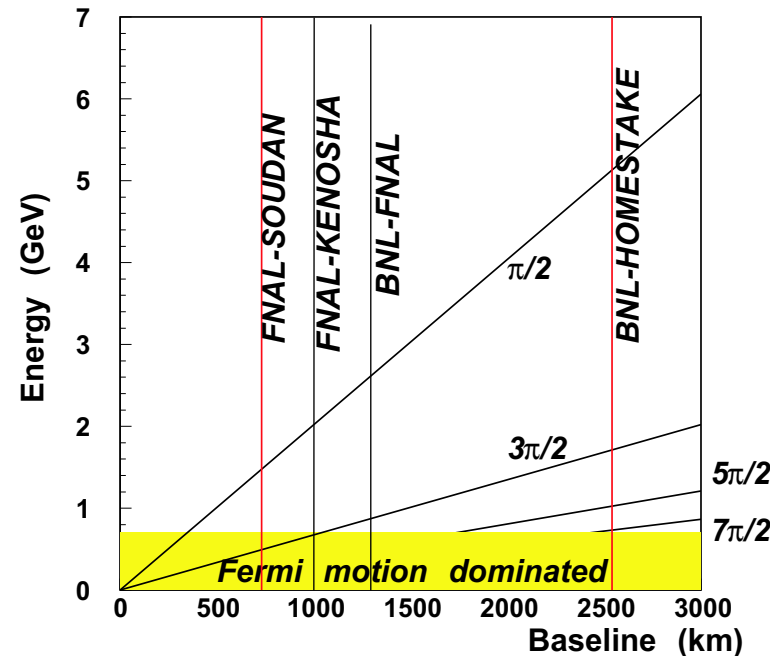
Low E_ν favorable for CP violation measurements.

If (still) need to disentangle matter effects from CP asymmetries, use the $n = 0$ and 1 oscillation maxima with E_1 as low as possible,

Ex: FNAL-Kenosha (986 km),

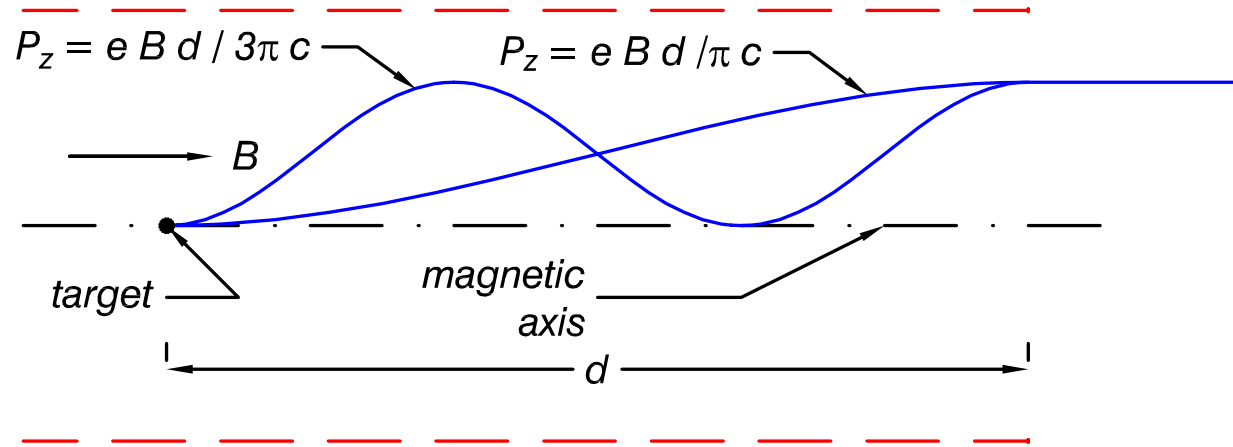
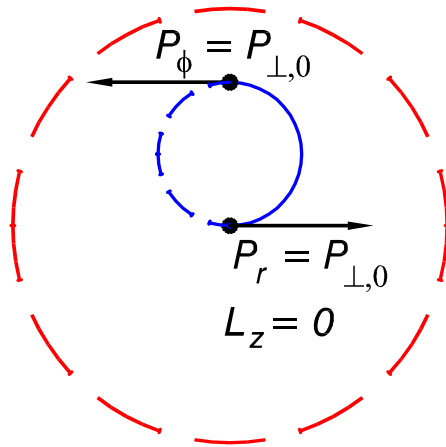
BNL-FNAL (1286 km).

Oscillation Nodes for $\Delta m^2 = 0.0025^2 \text{ eV}$



Narrowband Beam via Solenoid Focusing

(physics/0312022)



- Point-to-parallel focusing occurs for $P_\pi = e B d / (2n + 1) \pi c$.
- \Rightarrow Narrowband neutrino beam with multiple peaks at

$$E_\nu \approx \frac{4}{9} \frac{e B d}{(2n + 1) 2 \pi c}.$$

- \Rightarrow Can study several neutrino oscillation peaks at once, at

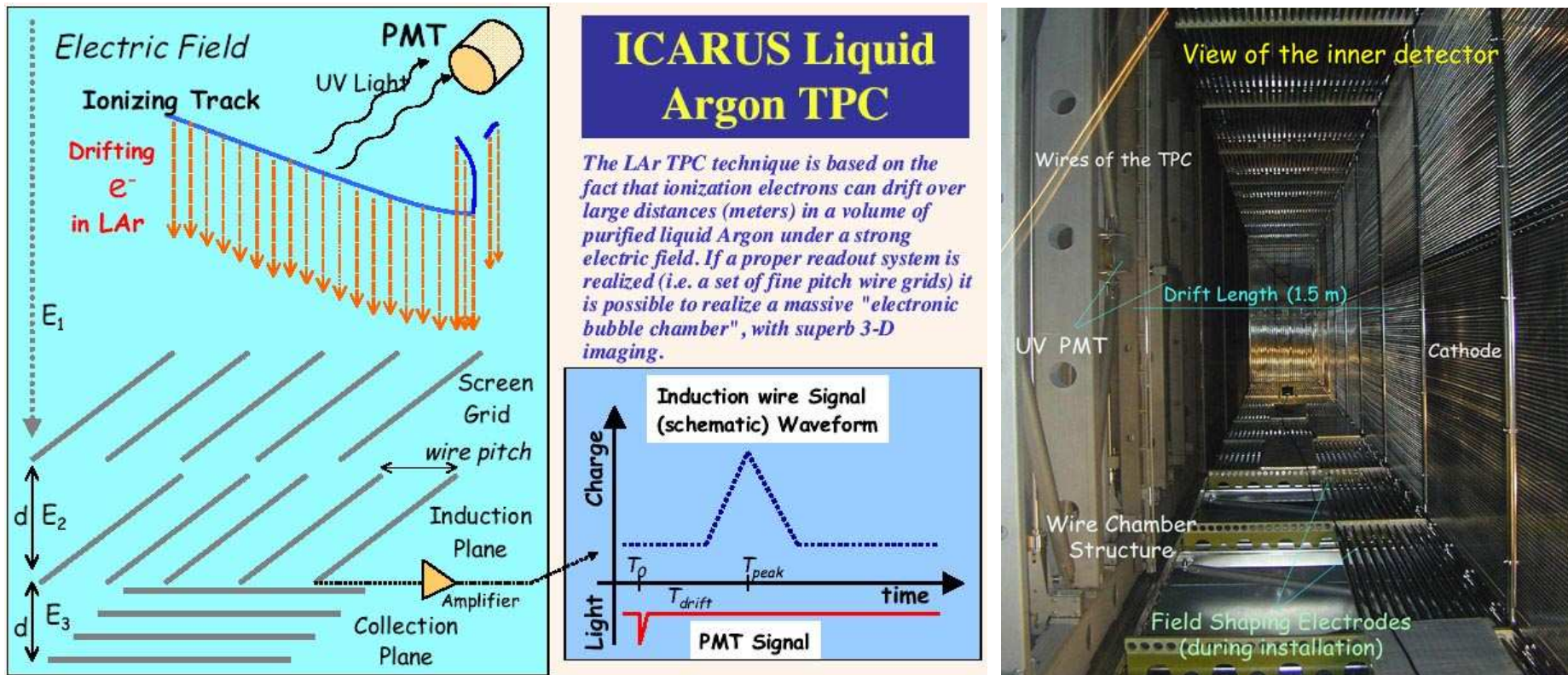
$$\frac{1.27 M_{23}^2 [\text{eV}^2] L [\text{km}]}{E_\nu [\text{GeV}]} = \frac{(2n + 1) \pi}{2}.$$

- Get both ν and $\bar{\nu}$ at the same time (while ν_e and $\bar{\nu}_e$ suppressed),
 \Rightarrow Must use detector that can identify sign of μ and e ,
 \Rightarrow Magnetized liquid argon TPC.

Liquid Argon TPC Overview

- A liquid argon time-projection chamber is a total-absorption tracking calorimeter = An electronic bubble chamber.
- It's efficiency for detection of ν_e appearance events will be greater than 90% for GeV energies.
(This is $\gtrsim 3$ times the efficiency of low- Z sampling detectors.)
- A large (> 10 kton) liquid argon TPC, if in a single cryostat, will cost very nearly the same as a low- Z sampling detector of the same mass.
(There is highly competitive industry support for production, purification and storage of large quantities of liquid argon.
Liquid scintillator costs 2.5 times as much as liquid argon, per unit mass.)
- The hardware of a liquid argon TPC is in a mature state, and readily scalable to large masses.
- More in need of further development is the software
– in the style of bubble chambers.
(Human scanning of event displays if necessary.)

ICARUS Liquid Argon TPC



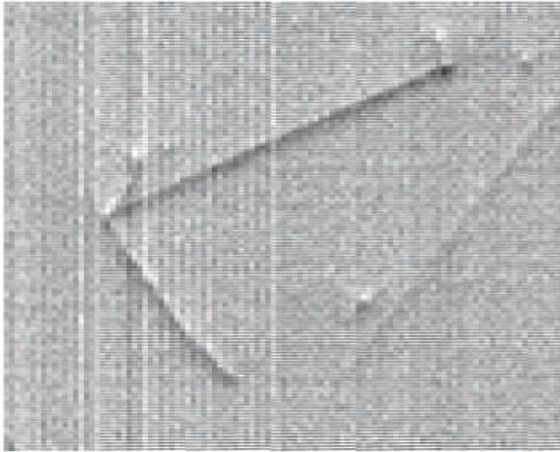
Liquid argon time projection chamber conceived by C. Rubbia (1977).

Largest implementation to date is the ICARUS T600 (600 ton) module, on the surface in Pavia, Italy.

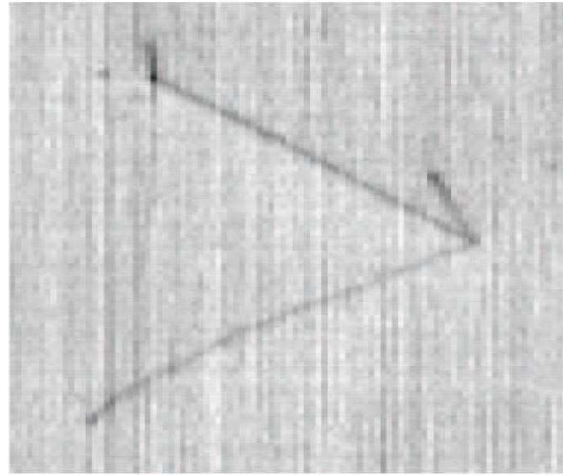
<http://www.aquila.infn.it/icarus/>

Events from the ICARUS T300 Cosmic Ray Test

Induction I



Induction II



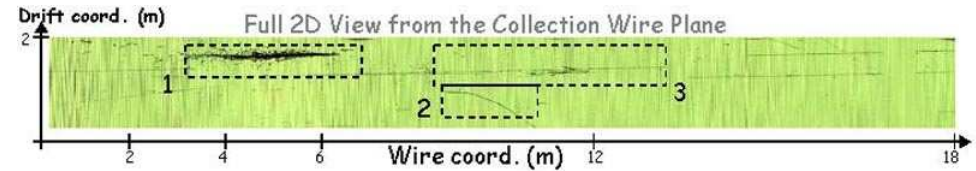
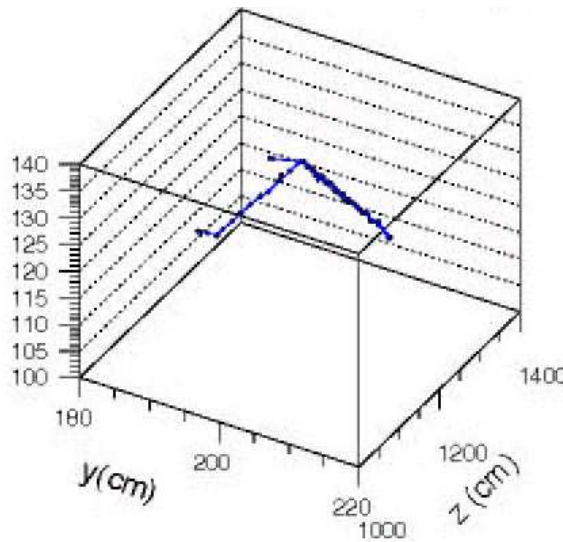
Collection



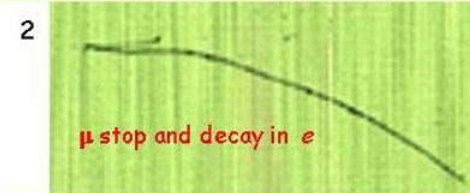
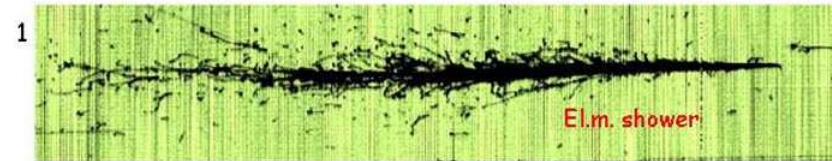
Above: 3 views of a low-energy hadronic interaction.

Right: Computer reconstruction.

Far right: Cosmic ray shower that includes a muon with a δ -ray, a stopping muon, and an electromagnetic shower.



Zoom details



Liquid Argon TPC Properties

- **3D tracking + total-absorption calorimetry.**
- **Pixel size: 3 mm × 3 mm (wire planes) × 0.6 mm (via 400 ns time sampling).**
- **$\rho = 1.4 \text{ g/cm}^3$, $T = 89\text{K}$ at 1 atm., $X_0 = 14 \text{ cm}$, $\lambda_{\text{int}} = 80 \text{ cm}$.**
- **A minimum ionizing particle yields 50,000 e/cm .**
- **Drift velocity of 1.5 m/msec at 500 V/cm \Rightarrow 5 m drift in 3 msec.**
- **Diffusion coef. $D = 6 \text{ cm}^2/\text{s} \Rightarrow \sigma = 1.3 \text{ mm}$ after 3 msec.**
- **Can have only 0.1 ppb of O_2 for a 5 m drift,
 \Rightarrow Purify with Oxisorb.**
- **Liquid argon costs \$0.7M/kton – and is “stored” not “used”.**

- Large modules ($\gtrsim 100$ kton) can be built using technology of liquid methane storage. (Total cost of a 100-kton detector is estimated to be \$200M.)

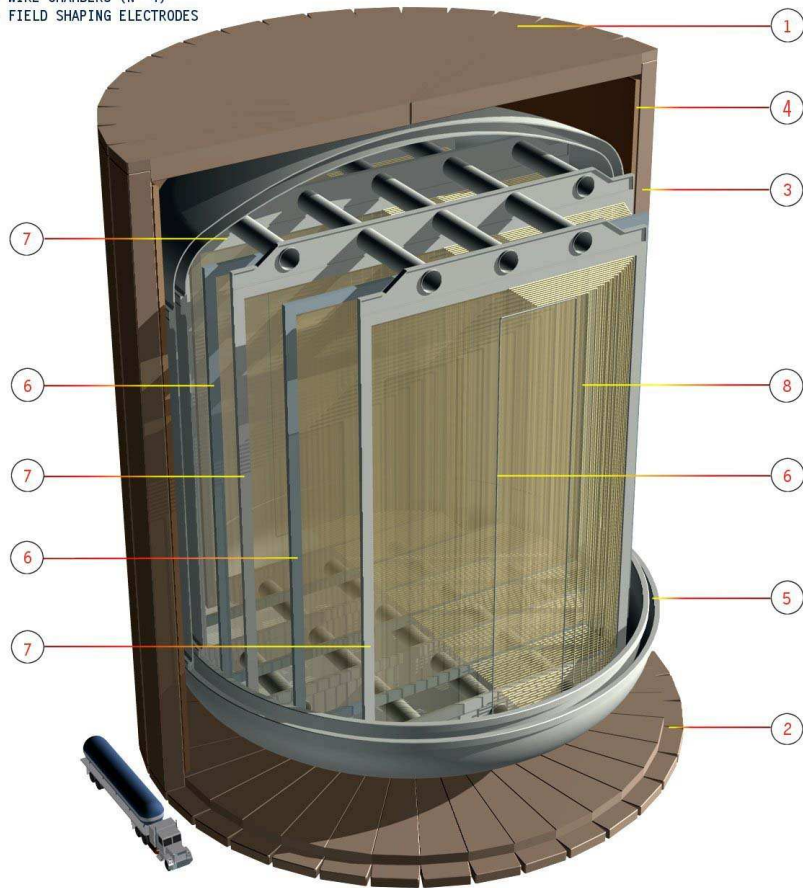


- Detector is continuously “live” and can be “self-triggered” using pipelined, zero-suppression electronics.
- Operates at the Earth’s surface with near zero overlap of cosmic ray events.
- Detector is compatible with operation in a magnetic field.

LANNDD – 100 kton Liquid Argon Neutrino and Nucleon Decay Detector

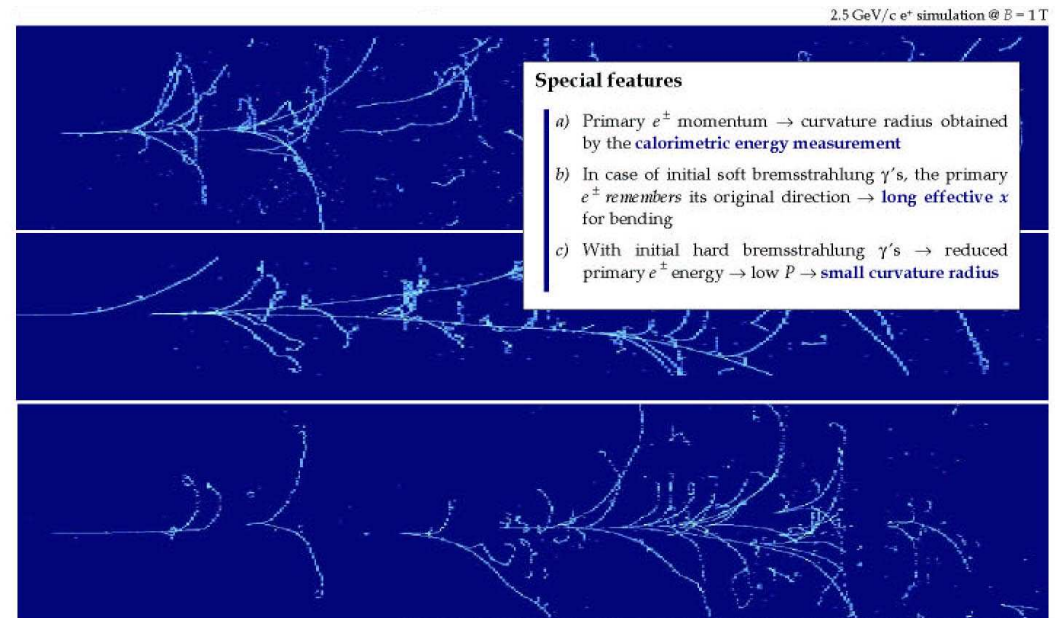
(astro-ph/0105442, Nucl. Instr. and Meth. A503, 136 (2003))

- 1- TOP END CAP IRON YOKE
- 2- BOTTOM END CAP IRON YOKE
- 3- BARREL IRON RETURN YOKE
- 4- COIL
- 5- CRYOSTAT
- 6- CATHODES (N° 5)
- 7- WIRE CHAMBERS (N° 4)
- 8- FIELD SHAPING ELECTRODES



LANNDD
Liquid Argon Neutrino and Nucleon Decay Detector

F. Sergiampietri-August 2000



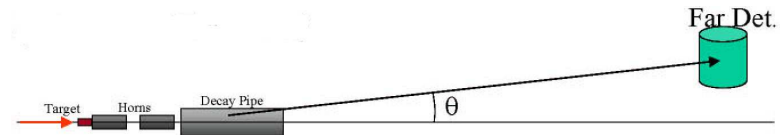
A. Bueno, M. Campanelli, A. Rubbia, IX International Workshop on "Neutrino Telescopes", VENICE, 2001

Can resolve sign of electron up to ≈ 2.5 GeV in a 0.5-T magnetic field.

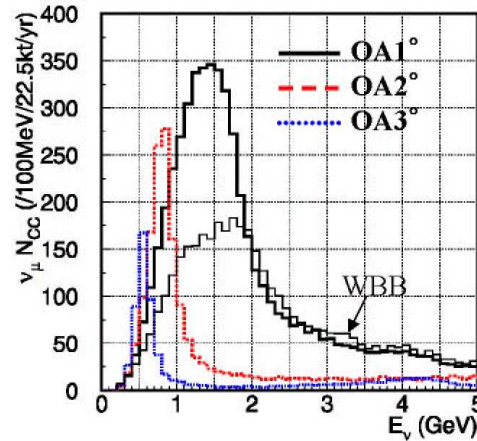
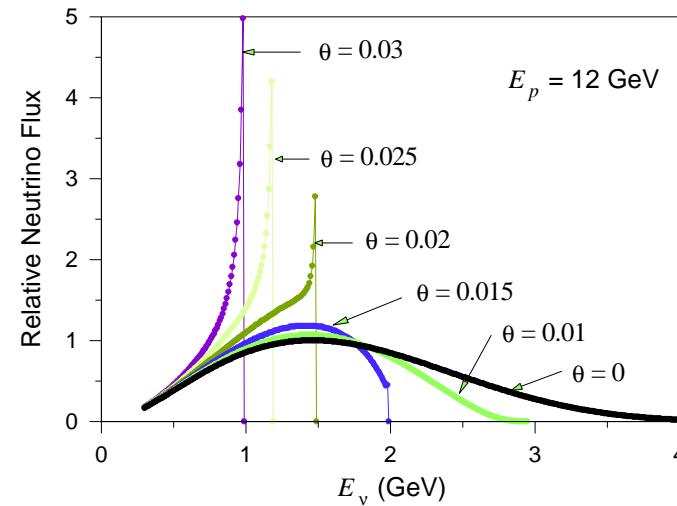
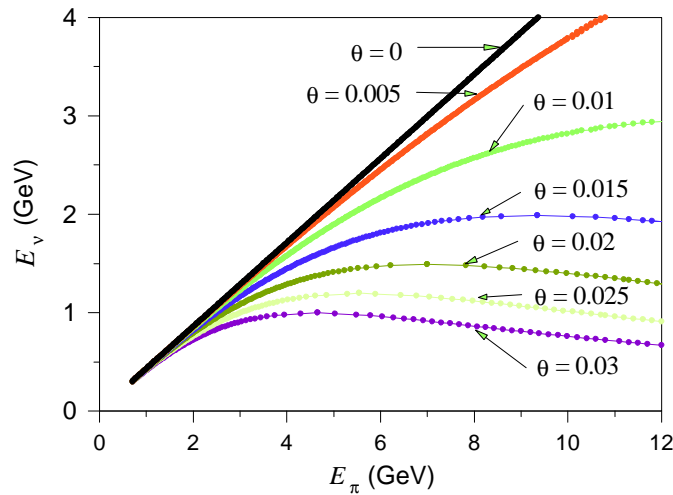
Max drift length of 5 m (limited by O_2 purity), \Rightarrow Several drift cells.

Before a Neutrino Factory, a Neutrino Superbeam

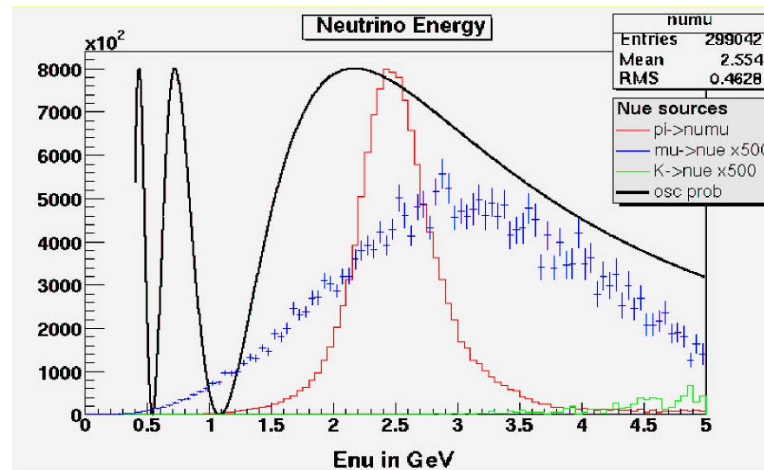
Use an Off-Axis Neutrino Beam (BNL E-889).



$\pi \rightarrow \mu\nu$ decay kinematics has a Jacobian peak: $\theta \approx 2^\circ / \text{GeV}$. (Sternheimer, 1955)



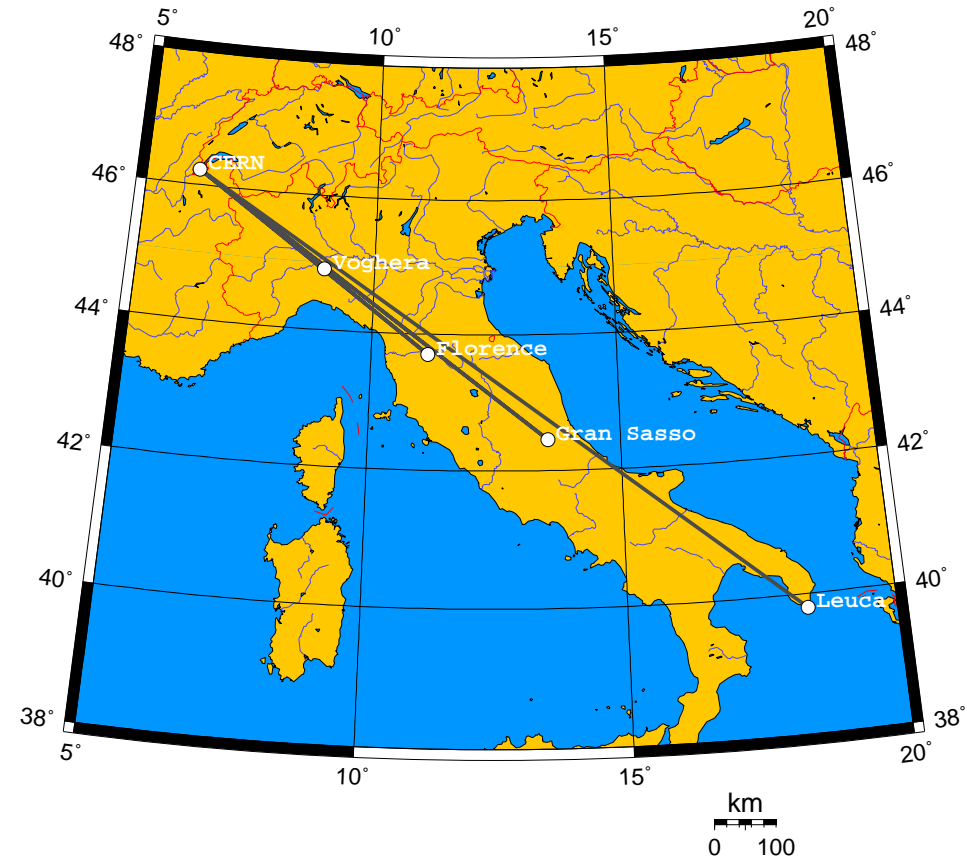
J2K (Kobayashi)



NUMI (Wojcicki)

Off-Axis Neutrino Beams from CERN

Site	Distance to CERN (km)	Lat.	Long.	$\angle_{\text{CERN to Leuca}}$
Voghera	270	44.9°	8.95°	4.4°
Florence	490	43.7°	11.15°	3.9°
Gran Sasso	730	42.45°	13.57°	2.5°
Leuca	1225	39.8°	18.35°	—

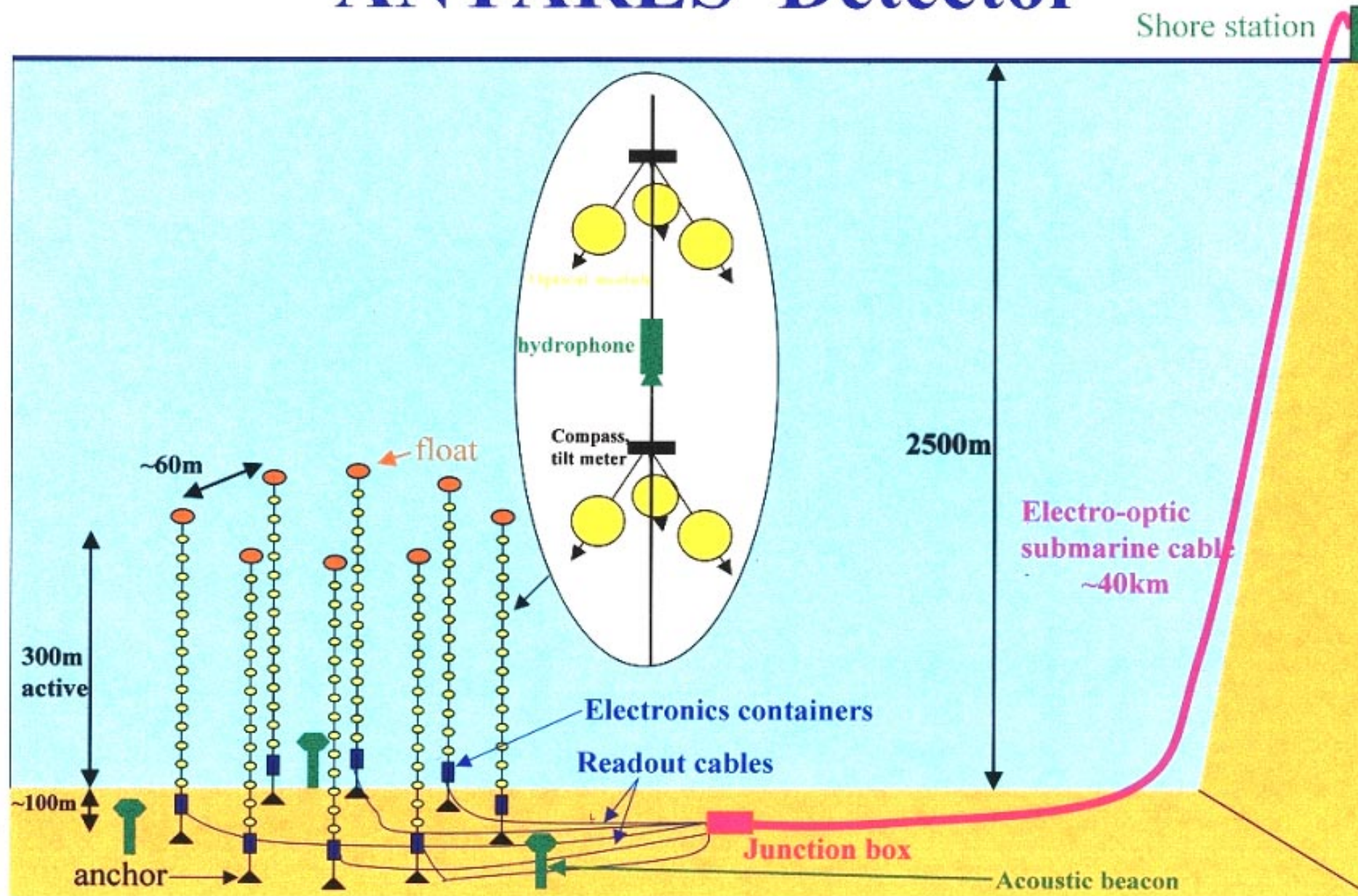


Could also use converted LNG tanker in the Gulf of Taranto:





ANTARES Detector



Destruction of Nuclear Bombs Using Ultra-High Energy Neutrino Beam

— dedicated to Professor Masatoshi Koshiba —

Hiroataka Sugawara* Hiroyuki Hagura† Toshiya Sanami‡

Abstract

We discuss the possibility of utilizing the ultra-high energy neutrino beam ($\simeq 1000$ TeV) to detect and destroy the nuclear bombs wherever they are and whoever possess them.

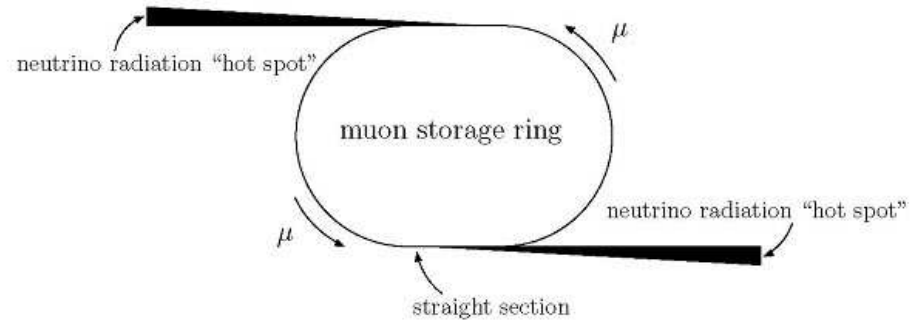


Figure 1: Neutrino radiation from a race track shaped muon storage ring. The decay of muons will produce the neutrino radiation emanating out tangentially everywhere from the ring. In particular, the straight sections in the ring will cause radiation "hot spots" where all of the decays line up into a pencil beam.

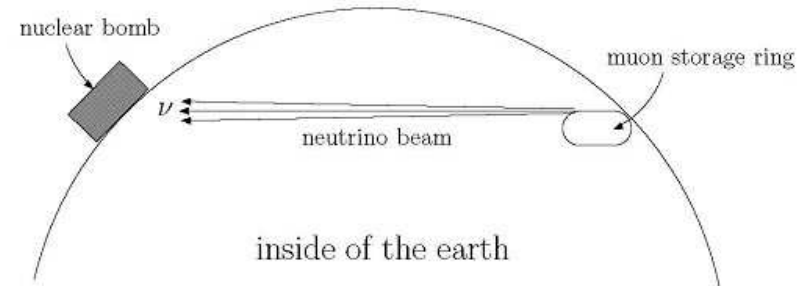


Figure 2: Neutrino beam is aimed at the nuclear bomb that is placed on the opposite side of the earth. The beam is emitted downstream from one of the straight sections of the muon storage ring (see fig. 1), and reaches the bomb after passing through the inside of the earth.