

# Mathematical Model Analysis for Hg Flow

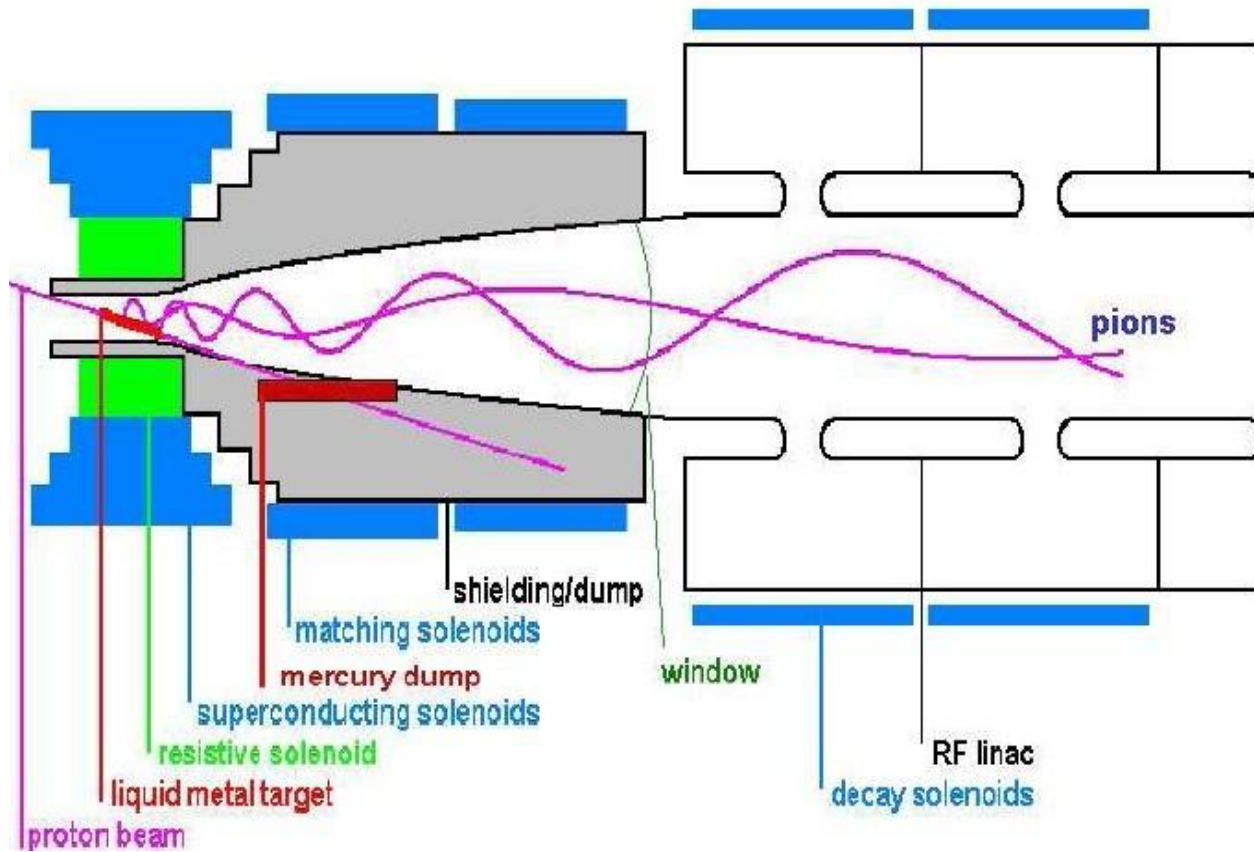
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- Equation of State

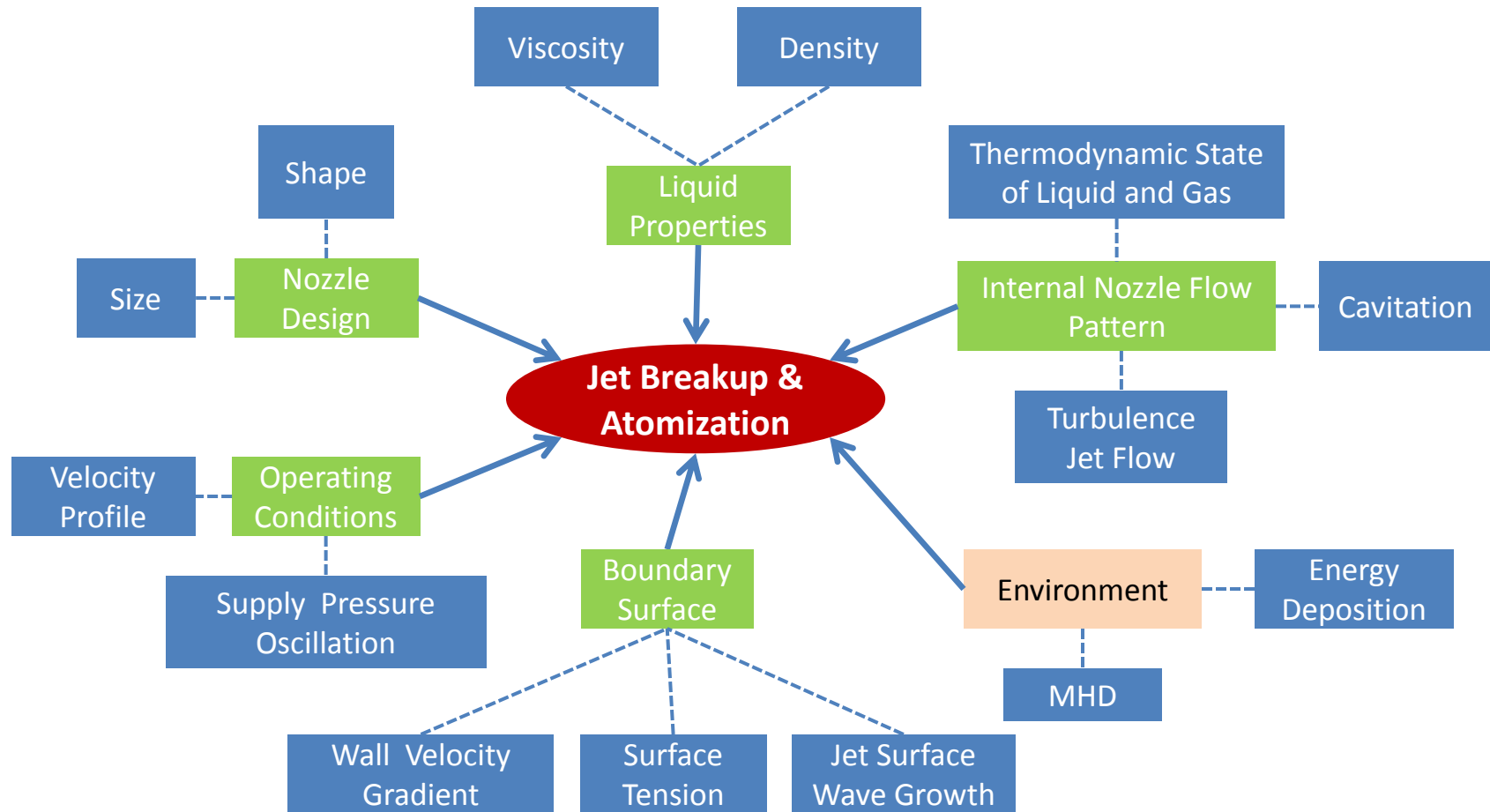
# Introduction



Neutrino Factory/Muon Collider

**The shape of the jet is important!**

# Introduction—Mechanism of Jet Breakup & Atomization



# Introduction—Dynamic Problems in Hg Target Flow

- **Internal Flow**

- Dynamics of MHD pipe flow
- The influence of nozzle design on jet exit conditions

- **Jet**

- Dynamics of free jet under MHD & energy deposition
- Jet breakup & instability mechanism and the effects of MHD & energy deposition
- Comparison with known classical free jet dynamics
- Inlet conditions to the Dump

# Mathematical Model for Hg Flow —Parameters (1)

Mercury Properties (25 °C )	
Density	13.546 kg/L
Sound Speed	1451 m/s
Viscosity	1.526E-3 kg/s·m
Dynamic Viscosity	1.127 m <sup>2</sup> /s
Thermal Conductivity	8.69 W/m·K
Electrical Conductivity	1E6 Siemens/m
Specific Heat	0.139 J/kg·K
Prandtl Number	0.025
Surface Tension	465 dyne/cm
Permeability	4*PI*1E-7

Variable Values	
Velocity in Pipe	3.4 m/s
Beam Energy	24 GeV
Beam Intensity	10 TP
Beam Pulse Length	2 ns
Magnetic Strength	20 Tesla
Pipe Diameter	2.54 cm (1")
Jet Diameter	1 cm
Beam Diameter	0.15 cm RMS
Static Pressure inside Pipe	18.5 Bar (kPa)
Dynamic Pressure inside Pipe	0.7 Bar (kPa)
Air Density Outside Jet	0.0013 kg/L
Jet Tilts w.r.t. Magnetic Axis	33 mrad
Beam Tilts w.r.t. Magnetic Axis	67 mrad

## Mathematical Model for Hg Flow —Parameters (2)

$$R_{e_{pipe}} = \frac{\rho v_{pipe} d_{pipe}}{\mu} = \frac{13.6 \times 10^3 \times 20 \times 0.0254}{1.5 \times 10^{-3}} = 7.666 \times 10^5$$

$$R_{e_{mag,pipe}} = \beta v_{jet} d_{jet} \sigma_{Hg} = 4 \times \pi \times 10^{-7} \times 3.4 \times 0.0254 \times 1 \times 10^6 = 0.108$$

$$M_{a_{pipe}} = \frac{v_{pipe}}{v_{sound}} = \frac{3.4}{1451} = 2.343 \times 10^{-3}$$

$$R_{e_{jet}} = \frac{\rho v_{jet} d_{jet}}{\mu} = \frac{13.6 \times 10^3 \times 20 \times 0.01}{1.5 \times 10^{-3}} = 1.8 \times 10^6$$

$$R_{e_{mag,jet}} = \beta v_{jet} d_{jet} \sigma_{Hg} = 4 \times \pi \times 10^{-7} \times 20 \times 0.01 \times 1 \times 10^6 = 0.2512$$

$$M_{a_{jet}} = \frac{v_{jet}}{v_{sound}} = \frac{20}{1451} = 0.014$$

$$W_{e_{jet}} = \frac{\rho d_{jet} v_{jet}^2}{\tau} = \frac{13.6 \times 10^3 \times 0.01 \times 20^2}{465 \times 10^{-3}} = 115000$$

Small Ma  Incompressible

Re\* for Curved Pipe?

Re\* effected by B Field?

 Turbulence ?

(Re\* for Curved Tube is Larger than Straight Pipe\*)

# Mathematical Model for Hg Flow — Incompressible Flow (No MHD)

## Governing Equations

$$\nabla \cdot \vec{V} = 0$$

$$\rho \left( \frac{\partial \vec{V}}{\partial t} + \vec{V} \cdot \nabla \vec{V} \right) = -\nabla P + \mu \nabla^2 \vec{V} + \rho \vec{g}$$

$$\rho C_p \frac{dT}{dt} = \frac{DP}{Dt} + \nabla \cdot (k \nabla T) + \Phi$$

$$\text{Where } \Phi = \mu \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

## Nondimensional Governing Equations

$$\nabla^* \cdot \vec{V}^* = 0$$

$$\rho^* \frac{D\vec{V}^*}{Dt^*} = -\nabla^* P^* + \frac{1}{Re} \nabla^* \cdot \left( \frac{\partial u_i^*}{\partial x_j^*} + \frac{\partial u_j^*}{\partial x_i^*} \right) + \frac{1}{Fr} \rho^*$$

$$\rho^* \frac{DT^*}{Dt^*} = Ec \frac{DP^*}{Dt^*} + \frac{1}{PrRe} \nabla^* \cdot (k^* \nabla^* T^*) + \frac{Ec}{Re} \Phi^*$$

$$V, r, t, \rho, P, \mu, g, C_p, \Phi, T, k \quad (11)$$

$$d_0, V_0, \rho_0, \mu_0, g, C_p, \Delta T \quad (7)$$

$$Re, Fr, Ec, Pr \quad (4)$$

$$\text{Reynolds Number } Re = \frac{\rho_0 V_0 d_0}{\mu_0}$$

$$\text{Froude Number } Fr = \frac{V_0}{\sqrt{d_0 g}}$$

$$\text{Mach Number } Ma = \frac{V_0}{c}$$

$$\text{Magnetic Reynolds Number } Re_{mag} = \beta_0 V_0 d_0 \sigma_0$$

$$\text{Eckert Number } Ec = \frac{V_0^2}{c_p \Delta T} = (\gamma - 1) Ma^2$$

$$\text{Prandtl Number } Pr = \frac{\mu_0}{\rho_0 k_0}$$



# Mathematical Model for Hg Flow — Incompressible Flow (MHD)

## Governing Equations

$$\nabla \cdot \vec{V} = 0$$

$$\rho \left( \frac{\partial \vec{V}}{\partial t} + \vec{V} \cdot \nabla \vec{V} \right) = -\nabla P + \mu \nabla^2 \vec{V} + \rho \vec{g} + \boxed{\vec{J} \times \vec{B}} \quad \text{Lorentz Force}$$

$$\rho C_p \frac{DT}{Dt} = \frac{DP}{Dt} + \nabla \cdot (k \nabla T) + \Phi + \boxed{\frac{1}{\sigma} \vec{J}^2}$$

$$\frac{\partial \vec{B}}{\partial t} = \nabla \times (\vec{V} \times \vec{B}) - \nabla \times \left( \frac{c^2}{4\pi\sigma} \nabla \times \vec{B} \right)$$

$$\vec{J} = \frac{c}{4\pi\beta} \nabla \times \vec{B}$$

$$\nabla \cdot \vec{B} = 0$$

## Nondimensional Governing Equations

$$\nabla^* \cdot \vec{V}^* = 0$$

$$\rho^* \frac{D\vec{V}^*}{Dt^*} = -\nabla^* P^* + \frac{1}{Re} \nabla^* \cdot \left( \frac{\partial u_i^*}{\partial x_j^*} + \frac{\partial u_j^*}{\partial x_i^*} \right) + \vec{J}^* \times \vec{B}^* + \frac{1}{Fr} \rho^* \vec{g}$$

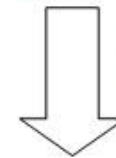
$$\rho^* \frac{DT^*}{Dt^*} = Ec \frac{DP^*}{Dt^*} + \frac{1}{PrRe} \nabla^* \cdot (k^* \nabla^* T^*) + \frac{Ec}{Re} \Phi^* + \frac{Re^2 Ec}{Re_{mag}} \vec{J}^{*2}$$

$$\frac{\partial \vec{B}^*}{\partial t^*} = \nabla^* \times (\vec{V}^* \times \vec{B}^*) - \frac{1}{Re_{mag}} \nabla^{*2} \vec{B}^*$$

$$\vec{J}^* = \frac{c}{4\pi Re} \nabla^* \times \vec{B}^*$$

$$\nabla^* \cdot \vec{B}^* = 0$$

$$V, r, t, \rho, P, \mu, g, J, B, C_p, \Phi, T, k, \sigma, c, \beta \quad (16)$$



$$d_0, V_0, \rho_0, \mu_0, g, C_p, \Delta T, \sigma_0, c, \beta_0, k_0 \quad (11)$$

$$Re, Re_{mag}, Fr, Ec, Pr \quad (5)$$

# Mathematical Model for Hg Flow — Proton Beam Energy Conversion

- Assume as a  $\delta$  function (A. Hassanein)

$$E_{pulse}(r, t) = q_b \delta(t - t_b) = q_0 \exp(-(r/r_b)^2) \delta(t - t_b) \text{ where } q_0 = \frac{Q_b}{\pi r_b^2}$$

- Instantaneous energy deposition (P. Sievers and P. Pognat)

$$E \rightarrow \Delta T_0 \xrightarrow{P = \kappa \alpha_V \Delta T} P \xrightarrow{\frac{dE_c}{dV} = \frac{(\alpha_V \Delta T(r))^2}{2\kappa}} E_c \xrightarrow{\frac{\kappa(\alpha_V \Delta T)^2}{2} = \frac{\rho v^2}{2}} v$$

- 1D parabolic energy distribution of the beam long the propagation direction of the Hg flow. Sin<sup>2</sup> Envelop with  $\tau=2\text{ns}$  pulse length (I.F. Barna et al.)

$$E_{pulse}(x, t) = E_0 \sin^2\left(\frac{\pi t}{\tau}\right) (1 - (x/x_s)^2)$$

- Monte Carlo Code (MARS, GRAN, FLUKA)
- Using Sergei's calculation for a 24-GeV, 10-TP proton beam (W.Bo, R. Samulyak)

# Mathematical Model for Hg Flow — Boundary Conditions

## Pipe Flow

- No slip at the pipe wall;
- Non-conducting wall;
- Velocity profile and initial values;
- Inlet and outlet pressure;
- Magnetic field profile;

## Jet Flow

- The normal component of the velocity field is continuous across the interface;
- The pressure jump at the interface is defined by the surface tension  $\tau$  and main radii of curvature:

$$\Delta P_{\Gamma} = \tau \left( \frac{1}{r_1} + \frac{1}{r_2} \right) \quad \longrightarrow \quad p^* = \frac{1}{W_a} \frac{s^{*2} + 2s_{\theta^*}^{*2} - s^* s_{\theta^*}^*}{(s^{*2} + s_{\theta^*}^{*2})^{3/2}}$$

- The normal component of the current density vanishes at the interface giving rise to the Neumann boundary condition for the electric potential:

$$\frac{\partial \varphi}{\partial \vec{n}_{\Gamma}} = \frac{1}{c} (\vec{V} + \vec{B}) \cdot \vec{n}$$

- Heat transfer balance at the free surface;

## Next Step

- Study the energy deposition conversion for proton beam on jet to set up proper mathematical model for energy equation;
- Jet exit condition:
  - The influence of discontinuity boundary conditions transiting from pipe flow to jet flow;
  - Nozzle design on jet exit condition
- Dynamics of MHD on pipe flow

# Equation of State (for compressible flow)

- Continuous Model
- Two Phase EOS
- Homogeneous Model
- Heterogeneous Model
- ISM Model
- SESAME Library
- Summation Method

# Equation of State — Continuous Model

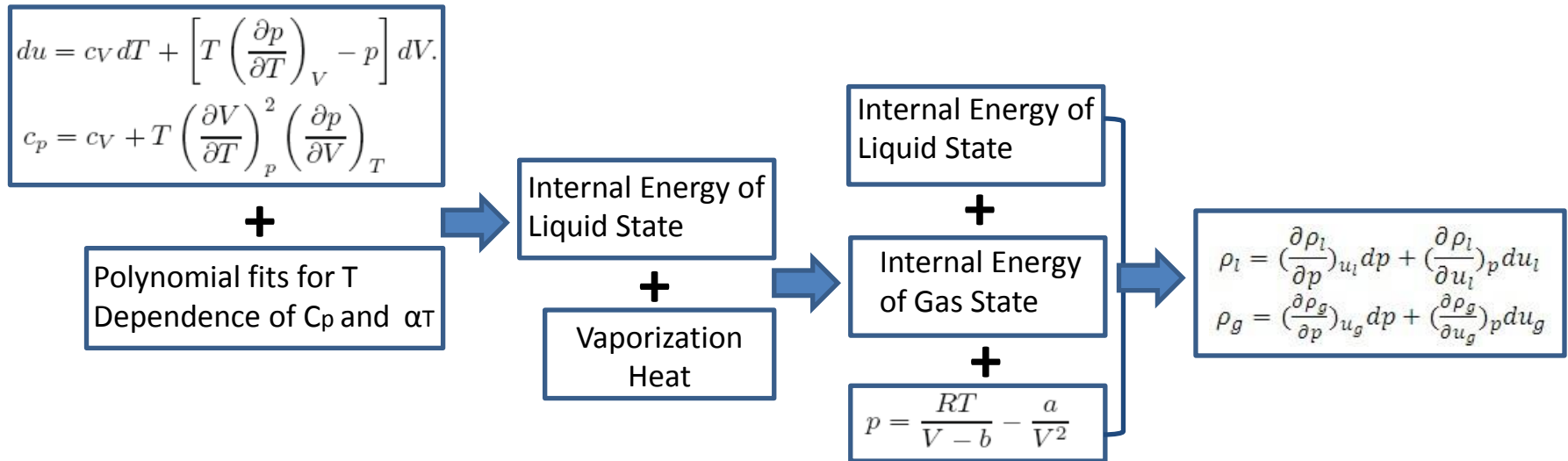
- Continuous Model (David P. Schmidt)

$$p = p_l^{sat} + p_{gl} \cdot \log \left[ \frac{\rho_g \cdot a_g^2 \cdot (\rho_l + \alpha \cdot (\rho_g - \rho_l))}{\rho_l \cdot (\rho_g \cdot a_g^2 - \alpha \cdot (\rho_g \cdot a_g^2 - \rho_l \cdot a_l^2))} \right]$$
$$p_{gl} = \frac{\rho_g \cdot a_g^2 \cdot \rho_l \cdot a_l^2 \cdot (\rho_g - \rho_l)}{\rho_g^2 \cdot a_g^2 - \rho_l^2 \cdot a_l^2} \alpha = \frac{\rho - \rho_l}{\rho_g - \rho_l}$$

$a_g$  and  $a_l$  are the sound speeds of the pure phases.  $\omega$  is zero in cartesian coordinates and unity in polar coordinates.

# Equation of State — Two Phase Model

- Two Phase Model (I.F. Barna et al.)



# Equation of State — Homogeneous Model

- Homogeneous Model (Bruggeman)

Liquid/vapor mixture is treated as a pseudo-fluid that obeys an equation of state of a single component flow.

Conductivity of the mixed phase is calculated by

$$\sigma_m = \begin{cases} 0 & \text{if } 1 \geq \beta \geq \frac{2}{3} \\ \frac{1}{2}(2 - 3\beta) & \text{if } \beta \leq \frac{2}{3} . \end{cases}$$



# Equation of State — Heterogeneous Model

- Heterogeneous Model (Jian Du, Roman Samulyak)

(1) Pure Vapor (Polytropic EOS):  $P = (\gamma_v - 1)E\rho$

(2) Liquid (Stiffened Polytropic EOS):

Stiffened EOS Model (incompressible single Liquid Phase )

$$P + \gamma P_\infty = (\gamma - 1)\rho(\epsilon + E_\infty)$$

Adiabatic exponent  $\gamma=3.19$ ;

Stiffening constant  $P_\infty=3000\text{bar}$ ;

Energy translation  $E_\infty=4.85\times 10^4 \text{ erg/sgK}$

(3) Liquid-Vapor Mixture:

$$P = P_{sat,l} + P_{vl} \log \left[ \frac{\rho_{sat,v} a_{sat,v}^2 (\rho_{sat,l} + \beta (\rho_{sat,v} - \rho_{sat,l}))}{\rho_{sat,l} (\rho_{sat,v} a_{sat,v}^2 - \beta (\rho_{sat,v} a_{sat,v}^2 - \rho_{sat,l} a_{sat,l}^2))} \right]$$

Where  $\beta = \frac{\rho - \rho_{sat,l}}{\rho_{sat,v} - \rho_{sat,l}}$

$$P_{vl} = \frac{\rho_{sat,v} a_{sat,v}^2 \rho_{sat,l} a_{sat,l}^2 (\rho_{sat,v} - \rho_{sat,l})}{\rho_{sat,v}^2 a_{sat,v}^2 - \rho_{sat,l}^2 a_{sat,l}^2}$$

## Equation of State — ISM Model

- ISM Model (Ihm, Song, Mason)

$$\frac{P}{\rho kT} = 1 + \frac{(B - \alpha) \rho}{1 + 0.22\lambda b\rho} + \frac{\alpha\rho}{1 - \lambda b\rho} \quad \Rightarrow \quad G(b\rho)^{-1} = \alpha\rho \left[ Z - 1 + \frac{(\alpha - B) \rho}{1 + 0.22\lambda b\rho} \right]^{-1} = 1 - \lambda b\rho$$

Where

$$B\rho_m = 0.403891 - 0.076484(\Delta H_{\text{vap}}/RT)^2 - 0.0002504(\Delta H_{\text{vap}}/RT)^4$$

$$\alpha\rho_m = a_1 \exp[-c_1(RT/\Delta H_{\text{vap}})] + a_2\{1 - \exp[-c_2(\Delta H_{\text{vap}}/RT)^{1/4}]\}$$

$$b\rho_m = a_1[1 - c_1(RT/\Delta H_{\text{vap}})] \exp[-c_1(RT/\Delta H_{\text{vap}})] \\ + a_2\{1 - [1 + 1/4c_2(\Delta H_{\text{vap}}/RT)^{1/4}]\} \exp[-c_2(\Delta H_{\text{vap}}/RT)^{1/4}]\}$$

$$a_1 = -0.1053, \quad a_2 = 2.9359$$

$$c_1 = 5.7862, \quad c_2 = 0.7966$$

Predict the density of Hg from the melting point up to 100 degree above the boiling Temperature.

# Mathematical Model for Hg Flow — Equation of State(4)

- [SESAME Library](#)

The SESAME opacity tables are compatible with the equation of state SESAME tables

- [Summation Method](#) (Ahmed Hassanein)

Sum of cold compression, ion thermal, and electron thermal terms

$$E = E_c + E_I + E_e$$

$$P = P_c + P_I + P_e$$

Where

$$E_c(\text{Energy of cold compression}): E_c = E_{c0} \left[ \left( \frac{\rho}{\rho_0} \right)^{\gamma-1} \right] \frac{\rho}{\rho_0} + P_{c0} \left( 1 - \frac{\rho}{\rho_0} \right); E_{c0} = \frac{P_{c0}}{\gamma-1};$$

$$P_c(\text{Pressure of cold compression}): P_c = P_{c0} \left[ \left( \frac{\rho}{\rho_0} \right)^{\gamma} - 1 \right]; P_{c0} = \frac{\rho c_s^2}{\gamma};$$

$$E_I(\text{Ion thermal energy}): E_I = 3nk(T - T_0);$$

$$P_I(\text{Ion thermal pressure}): P_I = G_I E_I;$$

$$E_e(\text{Electron thermal energy}): E_e = \frac{1}{2} \beta (T - T_0)^2, \beta = \beta_0 \left( \frac{\rho_0}{\rho} \right)^{2/3};$$

$$P_e(\text{Electron thermal pressure}): P_e = G_e E_e;$$

$G_I, G_e$ : the Gruneisen coefficients for the ion and electron;

$c_s$ : the speed of sound in the liquid target;