



Kirk T McDonald

Princeton U

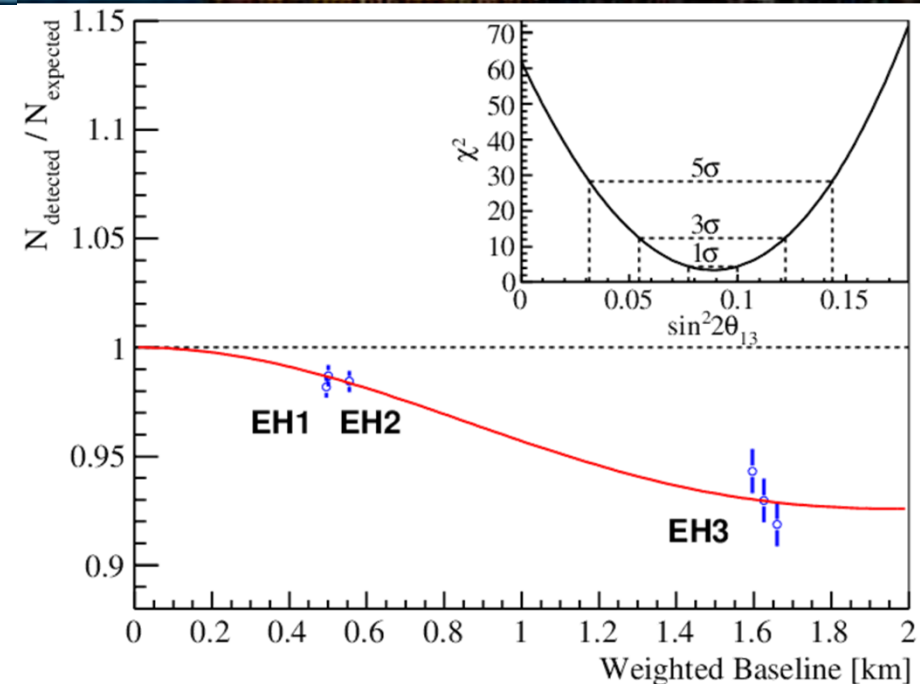
(updated June 1, 2012)

on behalf of the Daya Bay Collaboration

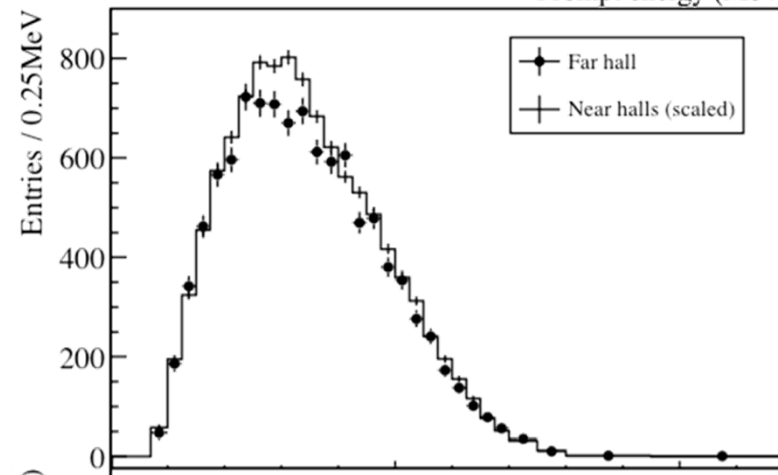
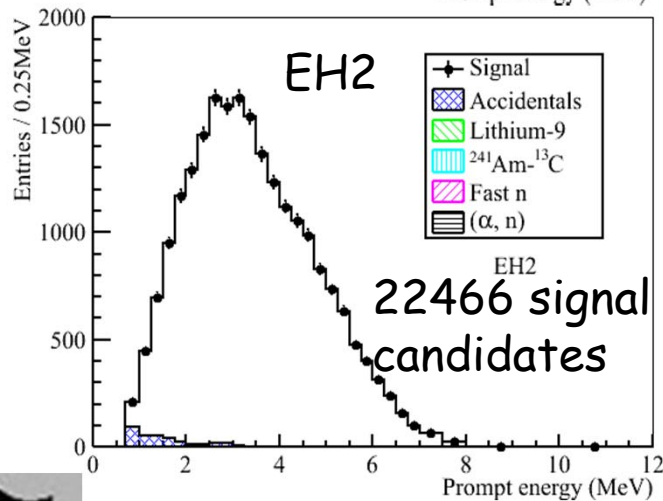
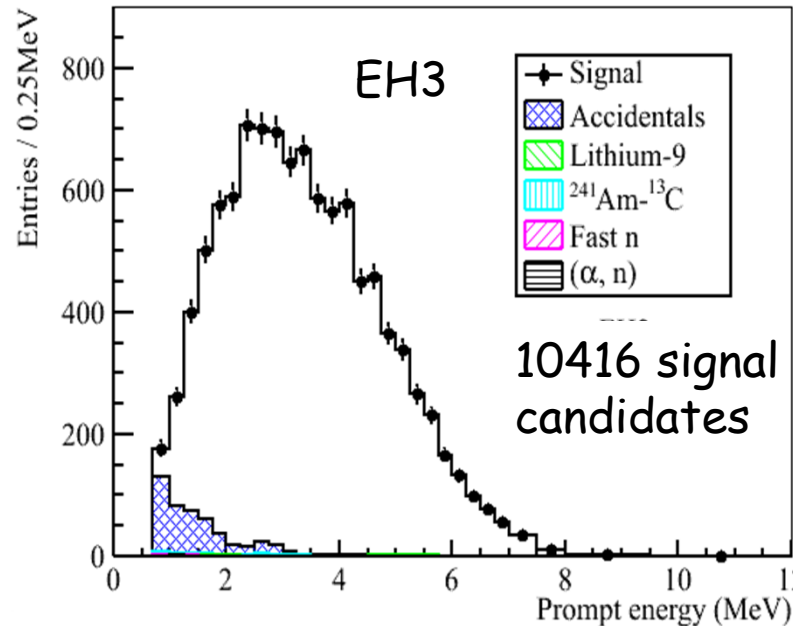
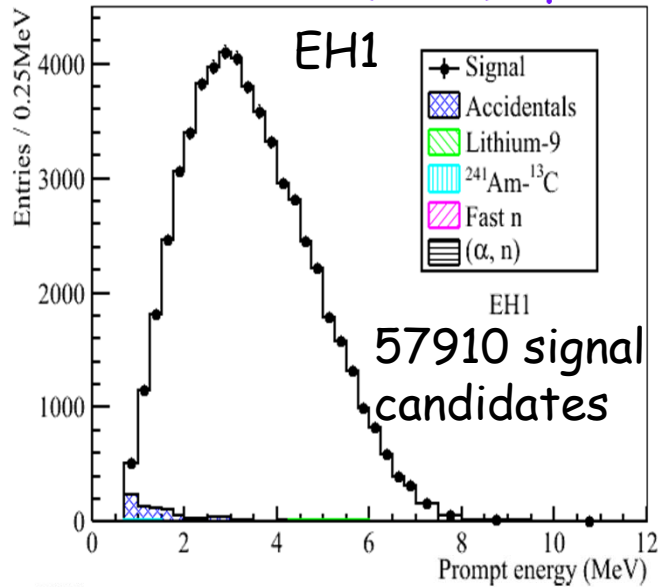
We observe that
 $\sin^2 2\theta_{13} = 0.092 \pm 0.016$ (stat.) ± 0.005 (syst.)
 after 55 days of operation with 6 detectors
 at 3 sites close to 3 pairs of ~ 3 GW reactors.

F.P. Ahn *et al.*, Phys. Rev. Lett. **108**, 171803 (2012).

After 110 days of running,
 $\sin^2 2\theta_{13} = 0.089 \pm 0.010$ (stat.) ± 0.005 (syst.)



Neutrino oscillations at Daya Bay deplete the e^+ energy spectrum near 3 MeV,
 \Rightarrow Far-detector (EH3) spectrum will be "flatter" if oscillations have occurred.



In predicting the Far (EH3) spectrum from the Near (EH1 and EH2) detectors, EH2 is weighted 7 times EH1.

$\sim 3\text{-}\sigma$ difference in near-far spectral shapes (preliminary).



Backup Slides



- The MNS matrix relates the mass eigenstates (ν_1, ν_2, ν_3) to the flavor eigenstates (ν_e, ν_μ, ν_τ):

MNS = Maki-Nakagawa-Sakata, who (1962) extended Pontecorvo's prediction (1957) from 2 to 3 neutrinos.

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

U_{e3} is last unknown matrix element

$$\begin{aligned} \sin^2 \theta_{12} &\sim 0.31 \\ \sin^2 \theta_{23} &\sim 0.43 \end{aligned}$$

- It can be described by three 2D rotations:

$$U_{\text{MNS}} = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta_{23} & \sin \theta_{23} \\ 0 & -\sin \theta_{23} & \cos \theta_{23} \end{pmatrix}}_{\text{Atmospheric}} \underbrace{\begin{pmatrix} \cos \theta_{13} & 0 & e^{-i\delta} \sin \theta_{13} \\ 0 & 1 & 0 \\ -e^{i\delta} \sin \theta_{13} & 0 & \cos \theta_{13} \end{pmatrix}}_{\text{Reactor}} \underbrace{\begin{pmatrix} \cos \theta_{12} & \sin \theta_{12} & 0 \\ -\sin \theta_{12} & \cos \theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{Solar}} \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha} & 0 \\ 0 & 0 & e^{i\beta} \end{pmatrix}}_{\text{Majorana Phases}}$$

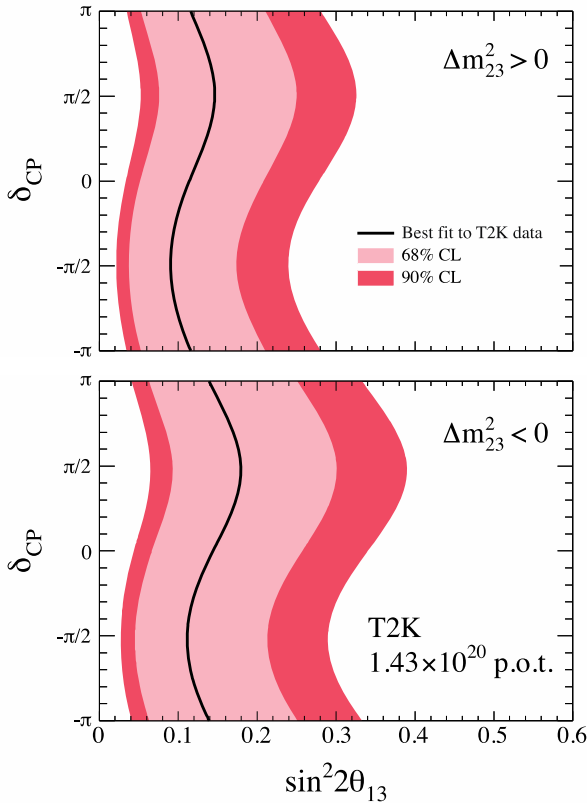
Measurement of the last unknown mixing angle, θ_{13} , is going to have a substantial impact on the future of neutrino physics.



Accelerator-based appearance expts.

Reactor-based disappearance expt.

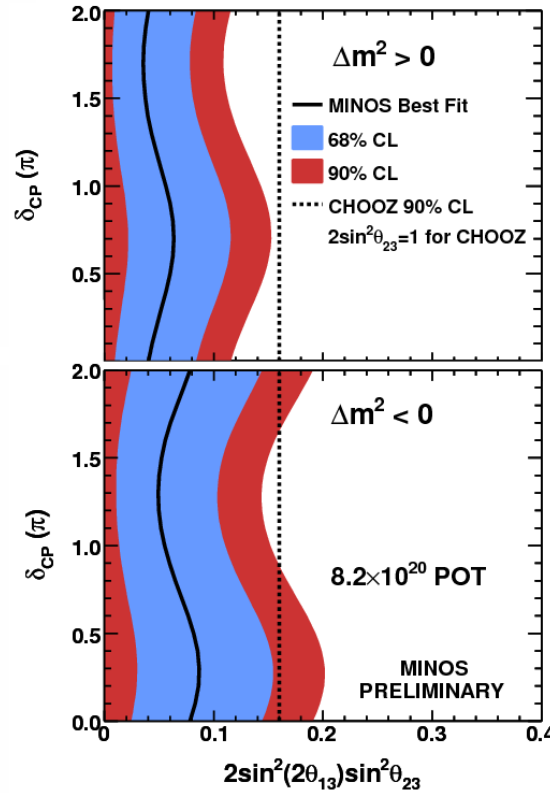
T2K



PRL 107, 041801 (2011)

$0.03(0.04) < \sin^2 2\theta_{13} < 0.28(0.34)$ at 90% CL

MINOS



PRL 107, 181802 (2011)

Double Chooz

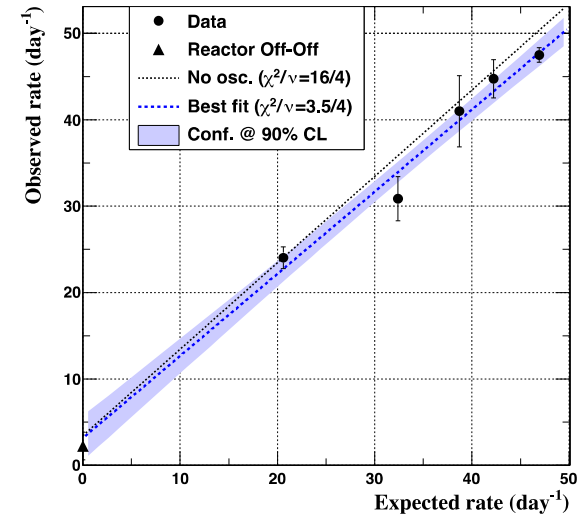


FIG. 2. Daily number of $\bar{\nu}_e$ candidates as a function of the expected number of $\bar{\nu}_e$. The dashed line is a fit to the data, the band is the 90% CL of this fit. The dotted line is the expectation in the no-oscillation scenario. The triangle indicates the one day measurement with both reactors off.

arXiv:1112.6353v1

$\sin^2 2\theta_{13} = 0.086 \pm 0.041(\text{stat}) \pm 0.030(\text{sys})$
 $(0.015 < \sin^2 2\theta_{13} < 0.16$ at 90% CL)

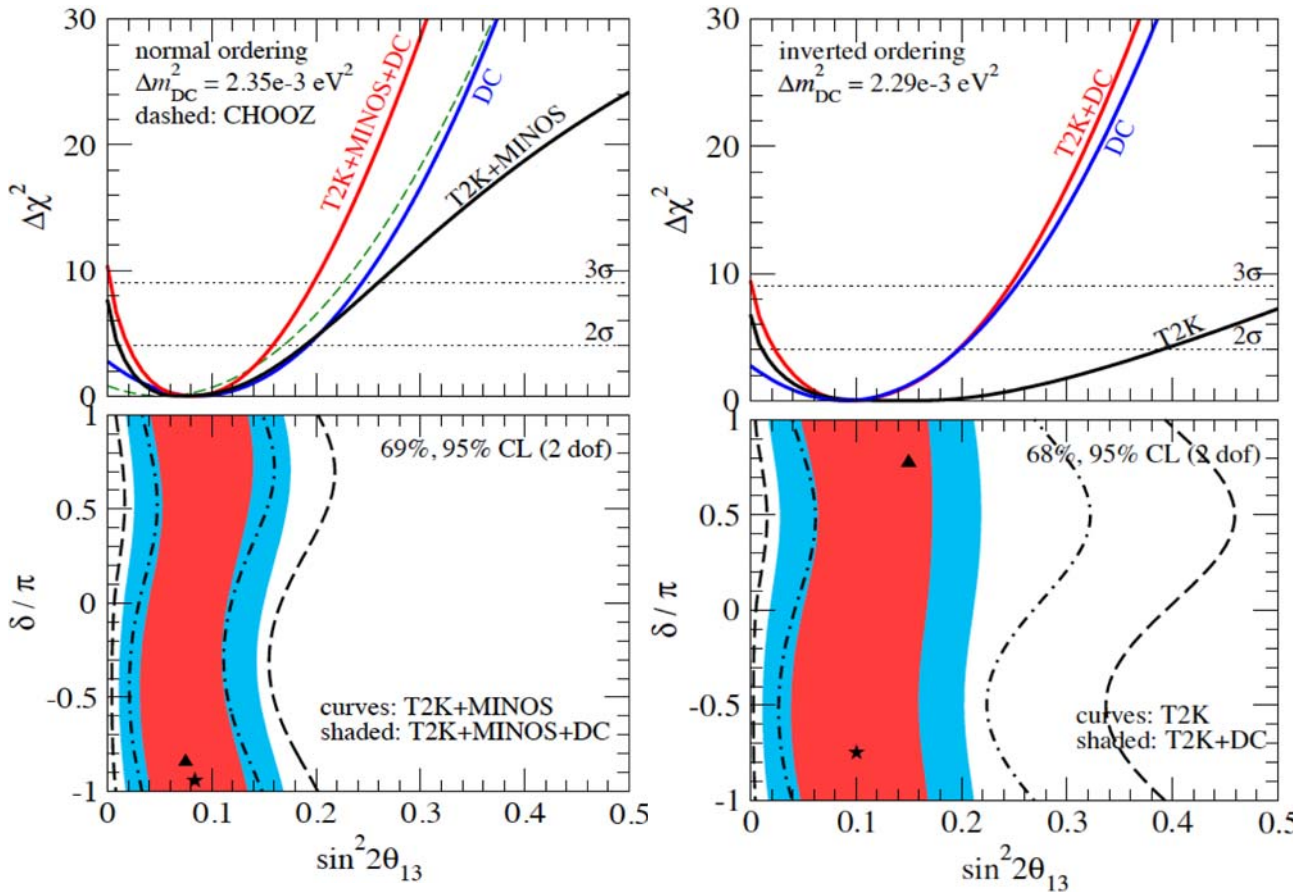


T2K has 6 candidate electron-appearance events, all close to the beam-entry side of the Super-K detector.

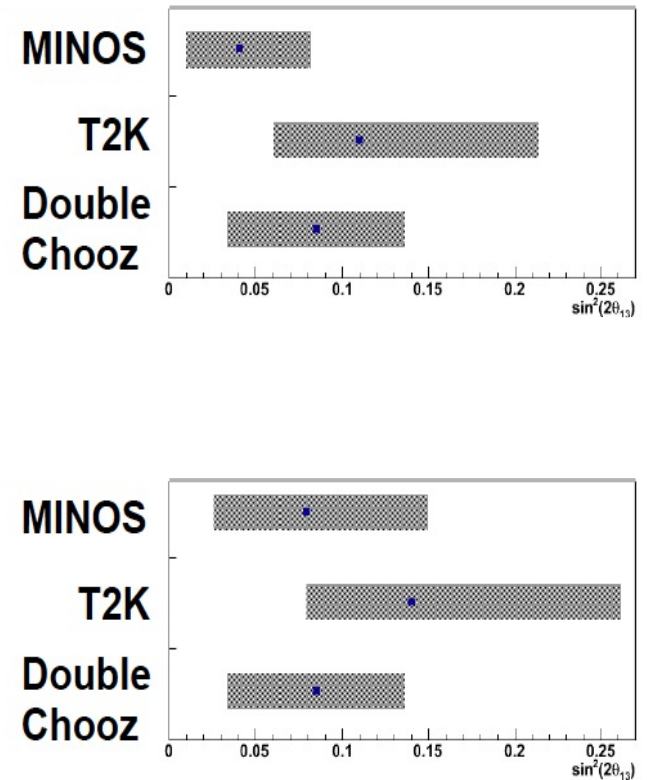
4/24/2012

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Seminar at Oxford U

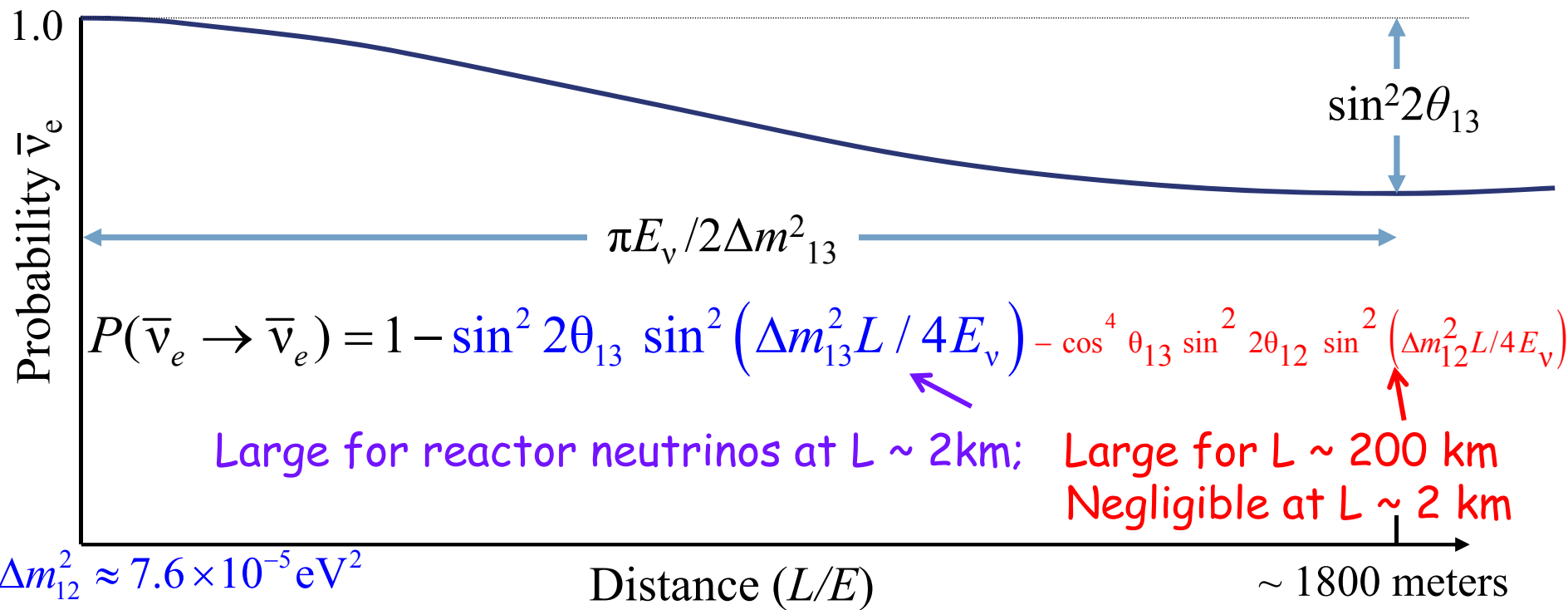


<http://arxiv.org/abs/1111.3330> Machado *et al.*



$\sin^2 2\theta_{13}$	best fit	1σ	2σ	3σ	$\Delta\chi^2(\theta_{13} = 0)$
normal ordering:	0.092	0.051–0.140	0.021–0.186	0.002–0.233	9.50
inverted ordering:	0.092	0.056–0.146	0.024–0.198	0.002–0.246	9.43

Combined measurements exclude $\theta_{13} = 0$ at greater than 3σ , but no single experiment has this sensitivity.



$$\Delta m_{12}^2 \approx 7.6 \times 10^{-5} \text{ eV}^2$$

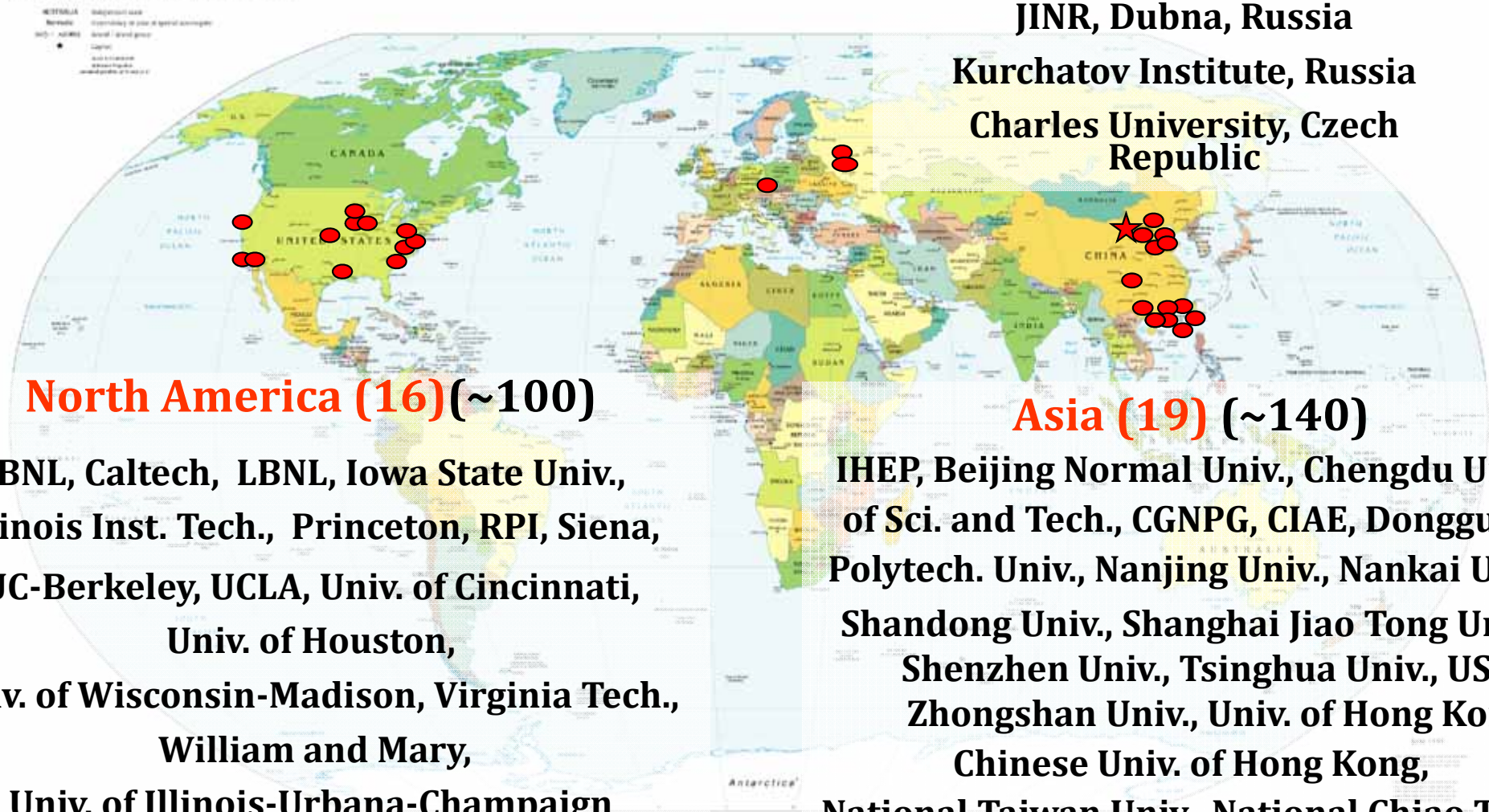
$$\Delta m_{13}^2 \approx \Delta m_{23}^2 \approx 2.4 \times 10^{-3} \text{ eV}^2$$

~ 1800 meters
(at 3 MeV)



~ 250 collaborators, 38 institutions

Political Map of the World, June 1999



- Identical near and far detectors cancel many systematic errors
- Multiple detectors at 3 locations boost statistics while reducing systematic errors with multiple independent measurements
- Three-zone detector design eliminates the need for spatial cuts which can introduce systematic uncertainties
- Shielding from cosmic rays and natural radioactivity reduces background rates
- Movable detectors allow possible cross calibration between near and far detectors to further reduce systematic errors

Experimental Layout

	Overburden	$R_{\nu\mu}$	E_{μ}	D1,2	L1,2	L3,4
EH1	250	1.27	57	364	857	1307
EH2	265	0.95	58	1348	480	528
EH3	860	0.056	137	1912	1540	1548

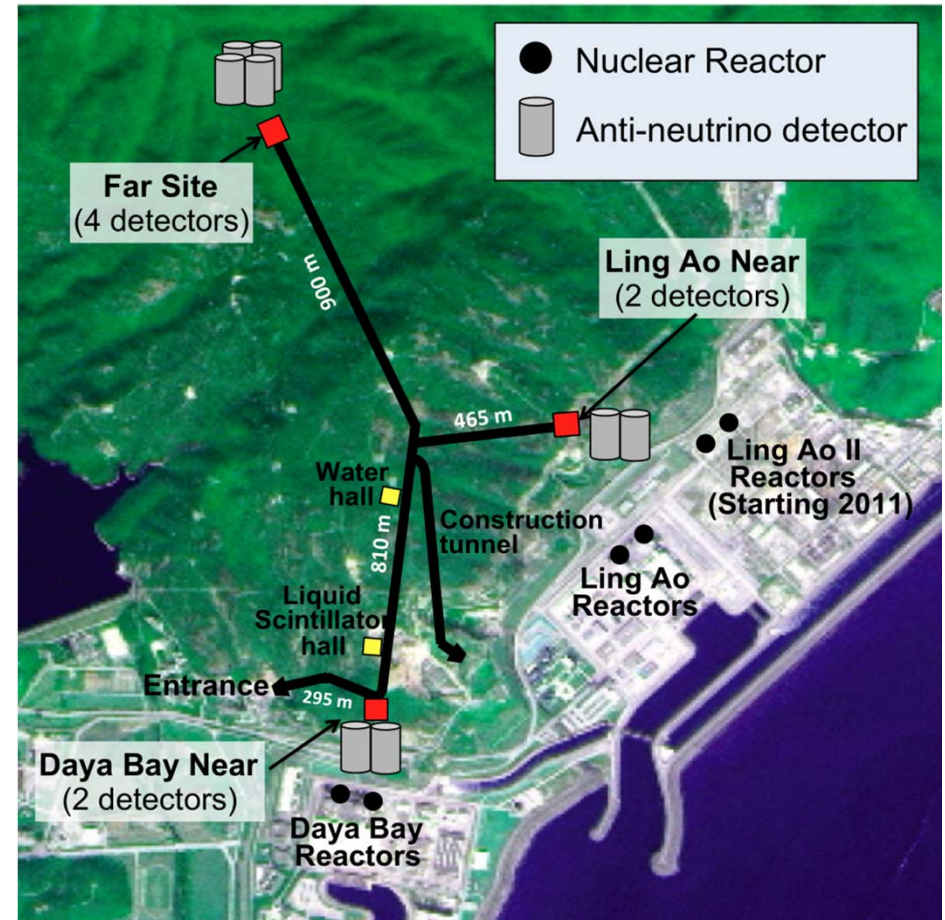
mwe Hz/m² GeV m m m

EH = Experimental Hall

D.B. = Daya Bay reactor cores

L.A. = Ling Ao reactor cores

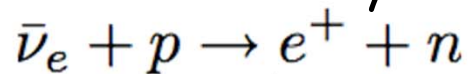
L.A. II = Ling Ao II reactor cores



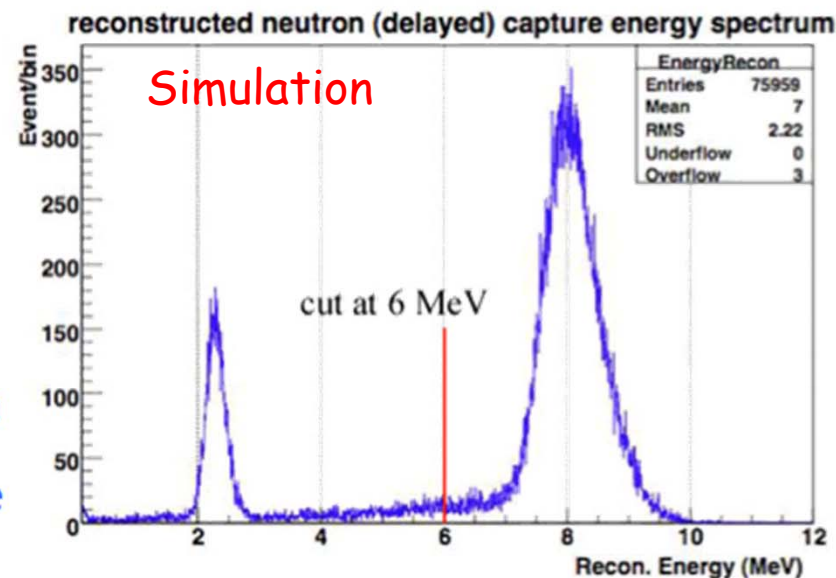
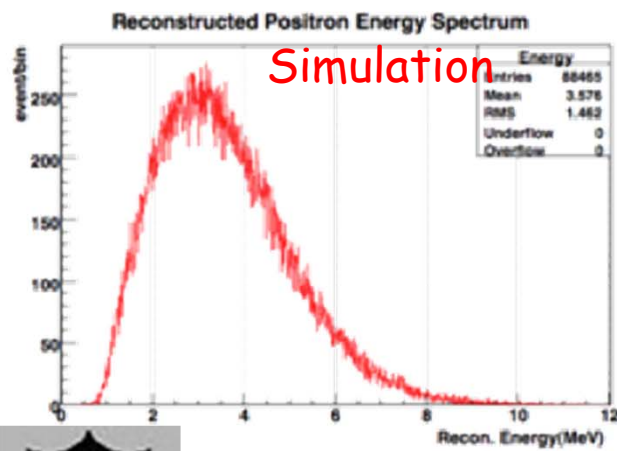
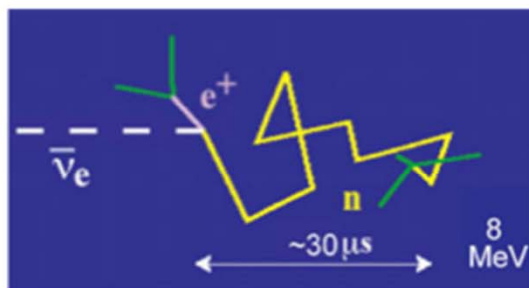
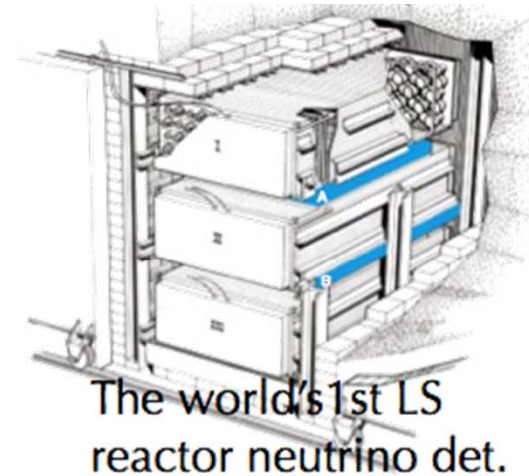
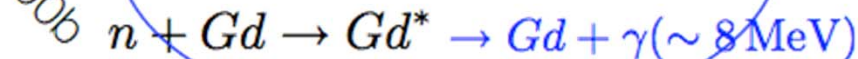
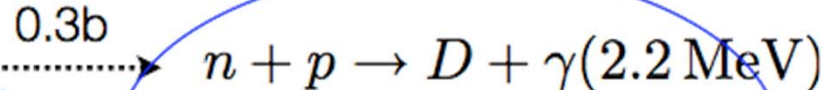
Daya Bay: 0.1% Gd doped liquid scintillator as target

Reines *et al.*, Phys. Rev. **117**, 159 (1960)

Inverse beta decay = IBD



$$E_{\bar{\nu}_e} \approx E_{e^+} + m_n - m_p$$



Correlated Signals

- background suppression
- well-defined target zone

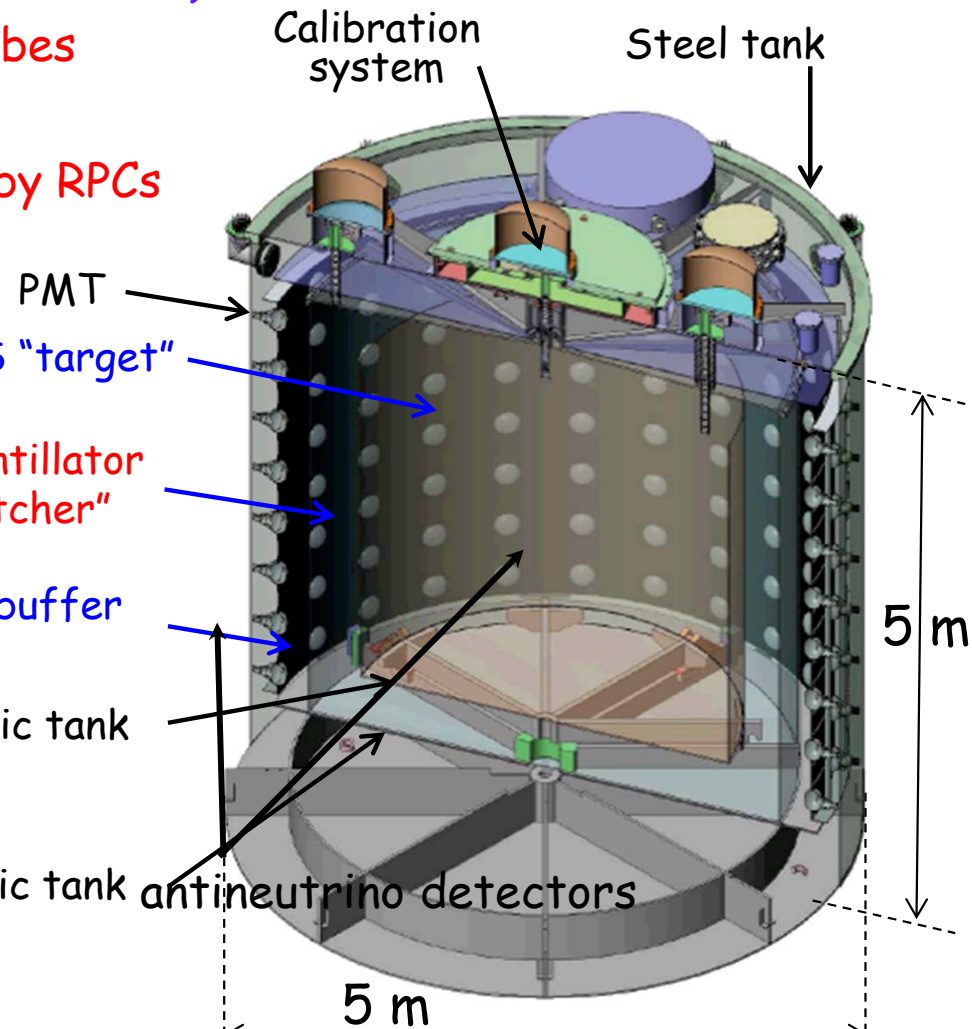
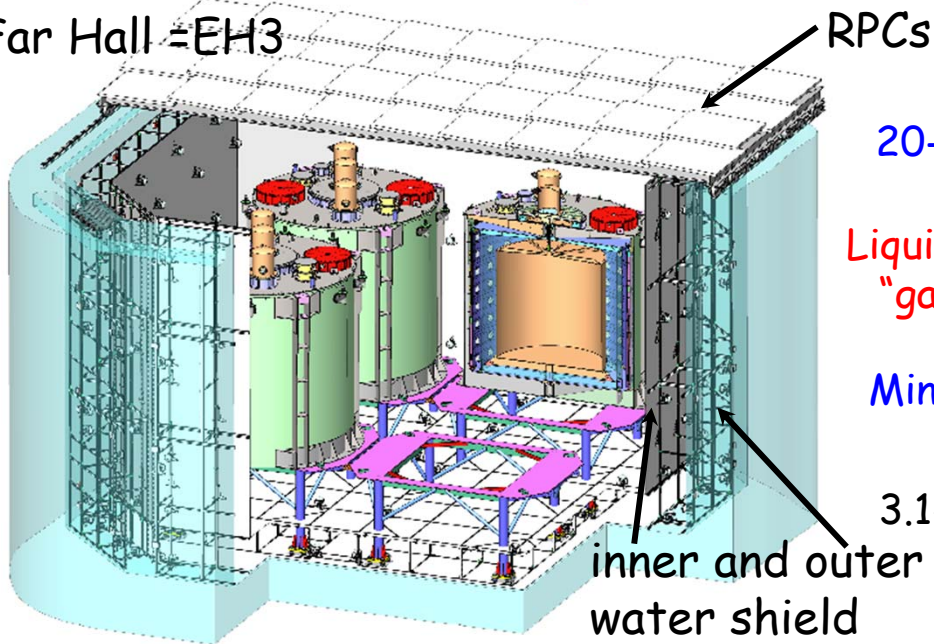


Daya Bay: 8 Identical Antineutrino Detectors

Only 6 detectors installed at present.

- Three-zone cylindrical design:
 1. Target: 20 t (0.1% Gd-based liquid scintillator)
 2. Gamma catcher: 20-t LS (no Gd, so no 8-MeV neutrons detected here)
 3. Buffer : 40-t (mineral oil; no scintillator, so no e, n or γ detected here)
- Detector mass ~ 110 ton (\Rightarrow need big crane to lift!)
- 192 low-background 8" photomultiplier tubes
- Reflectors at top and bottom
- Detectors sit in a pool of water, covered by RPCs

Far Hall =EH3



PMT = PhotoMultiplier Tube
 AD = Antineutrino Detector
 RPC = Resistive Plate Chamber

Automated calibration system

→ Routine weekly deployment of sources, lowered into any of the 3 liquid zones.

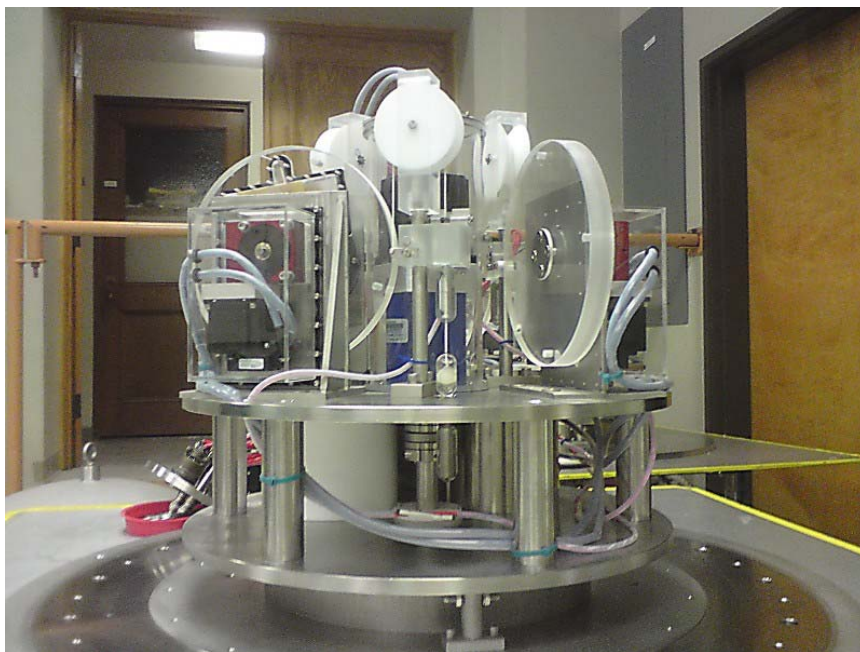
LED light sources

→ Monitoring optical properties

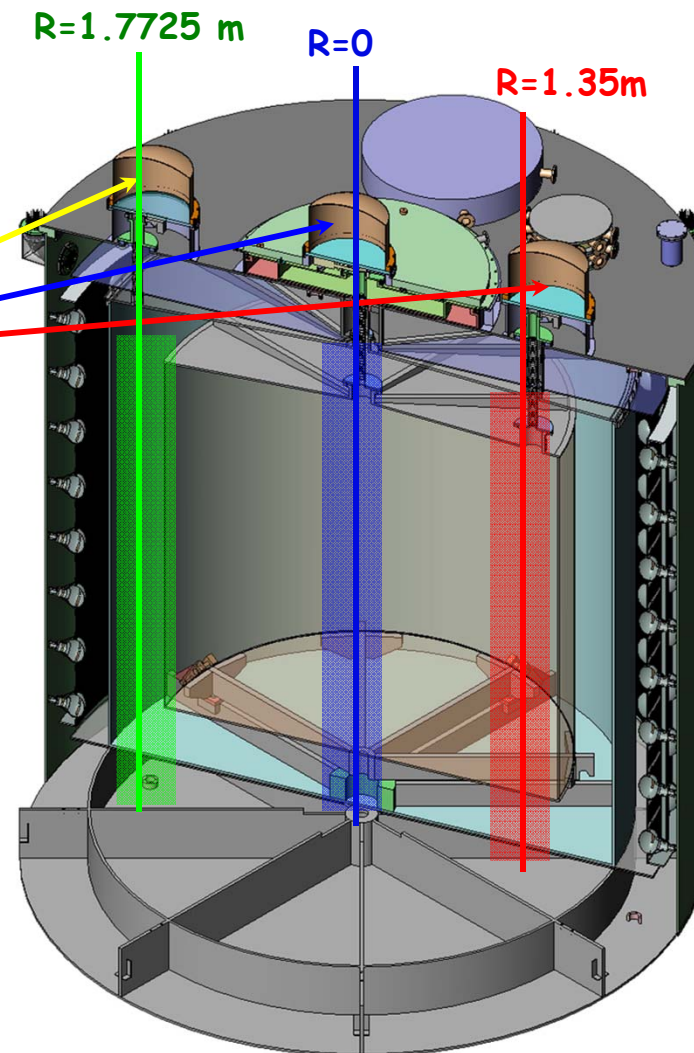
e^+ and n radioactive sources (fixed energy)

→ Energy calibration

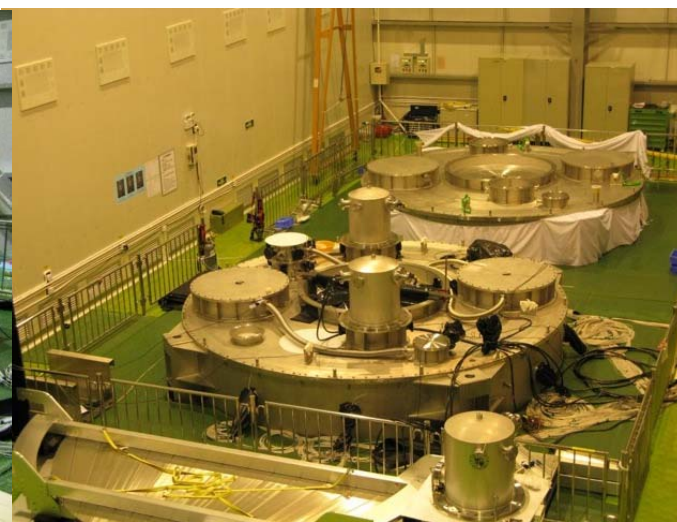
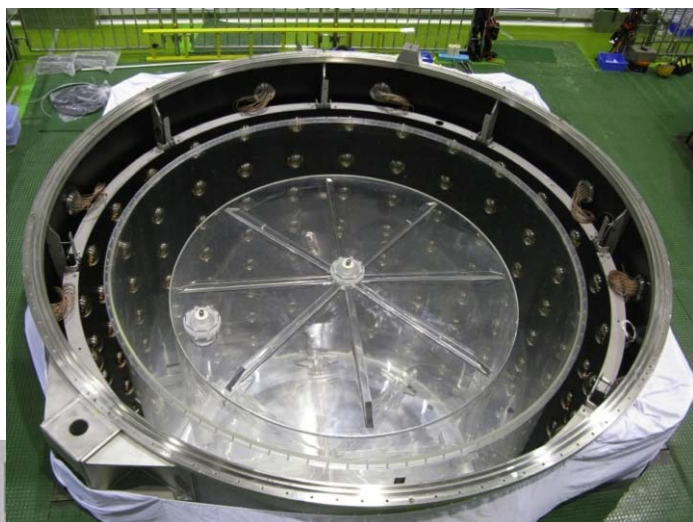
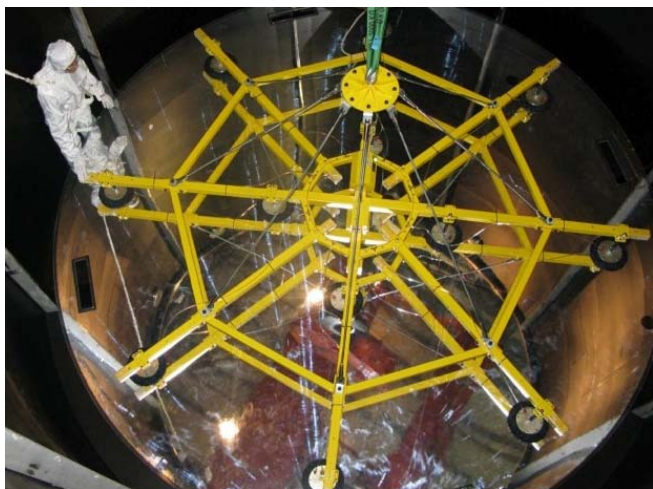
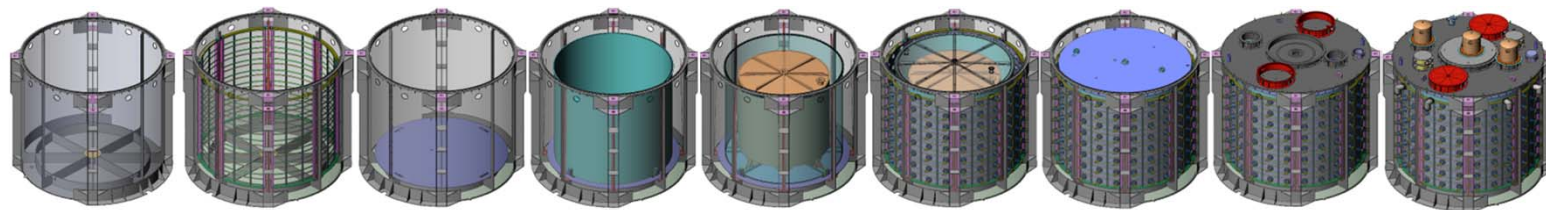
- ^{68}Ge source
- $\text{Am-}^{13}\text{C} + ^{60}\text{Co}$ source
- LED diffuser ball



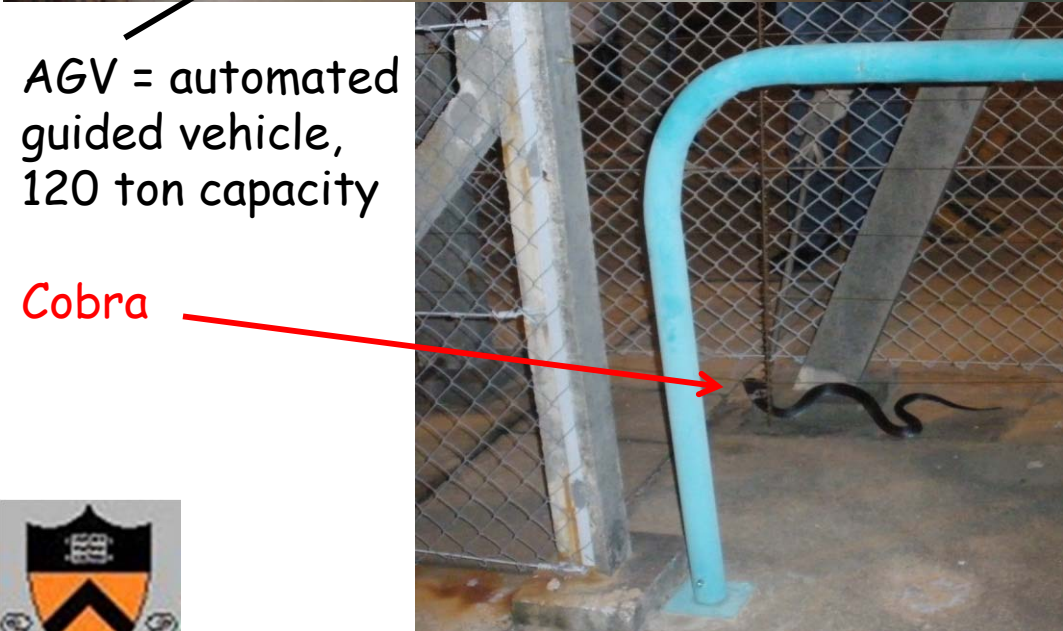
Automated calibration system



Antineutrino Detector Assembly



Transporting the Antineutrino Detectors



AGV = automated guided vehicle, 120 ton capacity

Cobra



4/24/2012

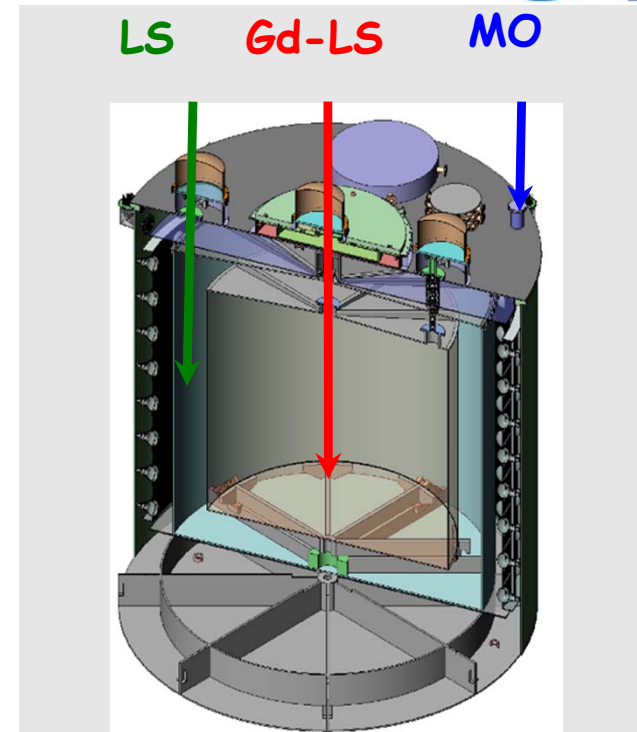
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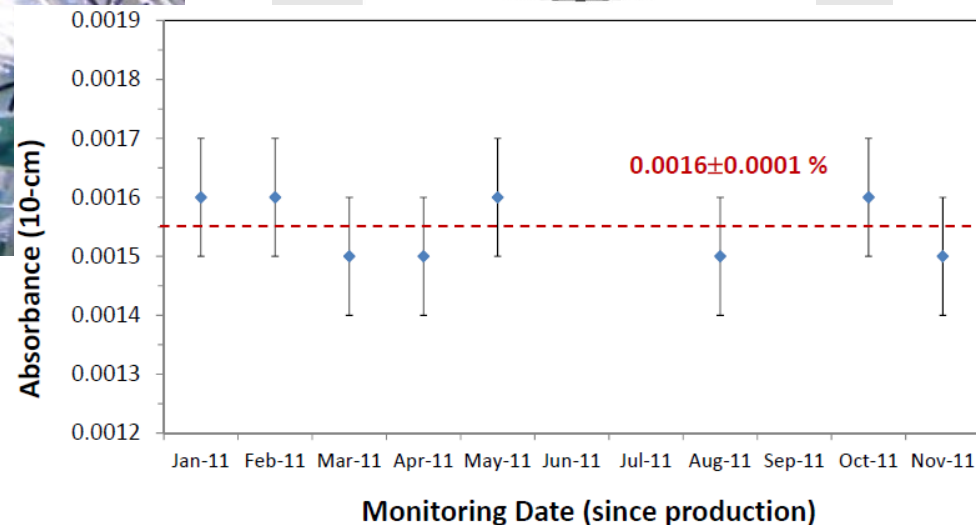


Liquid Scintillator production hall, underground



Detectors are filled from same reservoirs
“in-pairs” within < 2 weeks.

Target mass determined to $\pm 3\text{kg}$
out of 20,000 or < 0.02%.



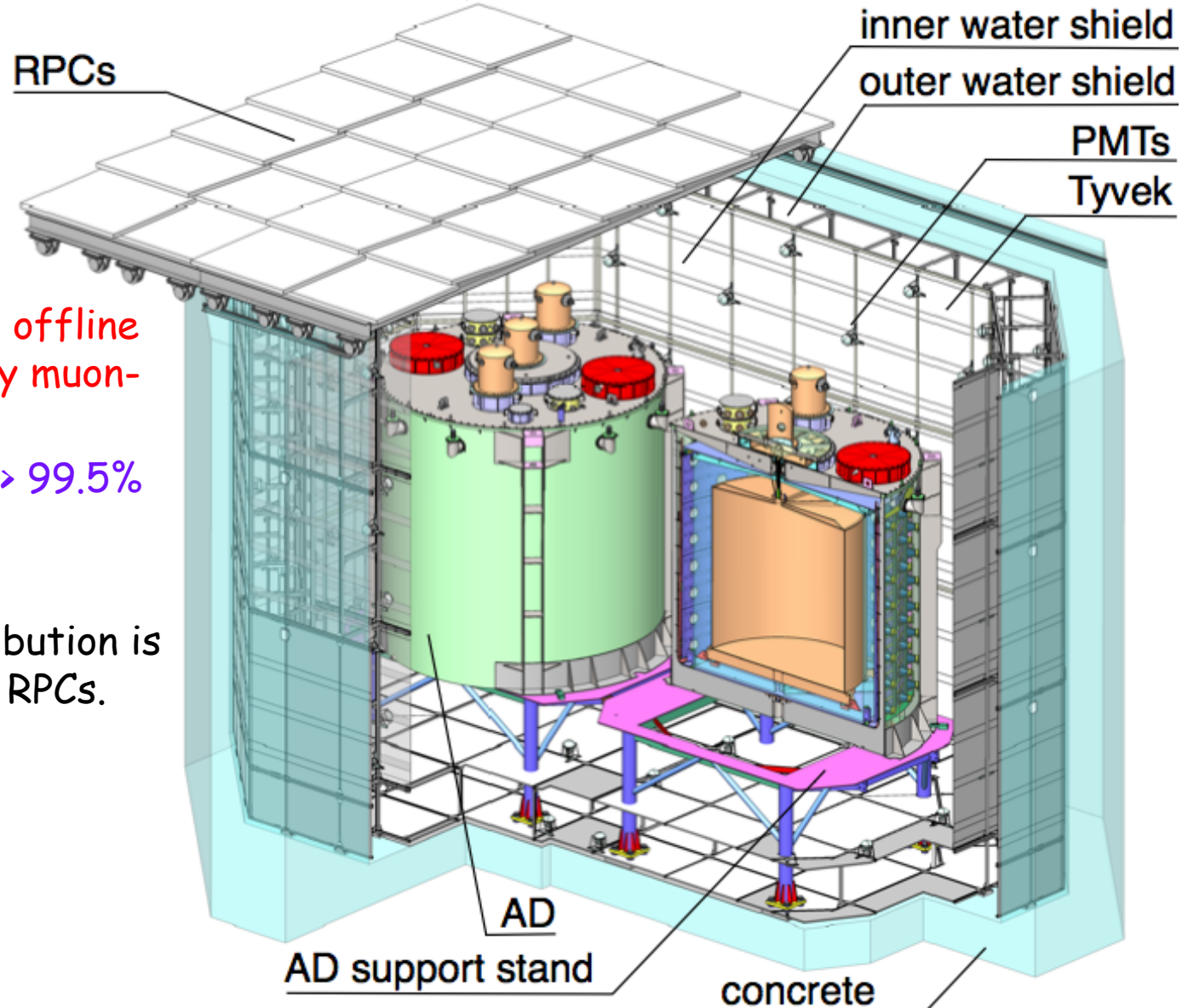
Gd-loaded scintillator shows good stability with time.

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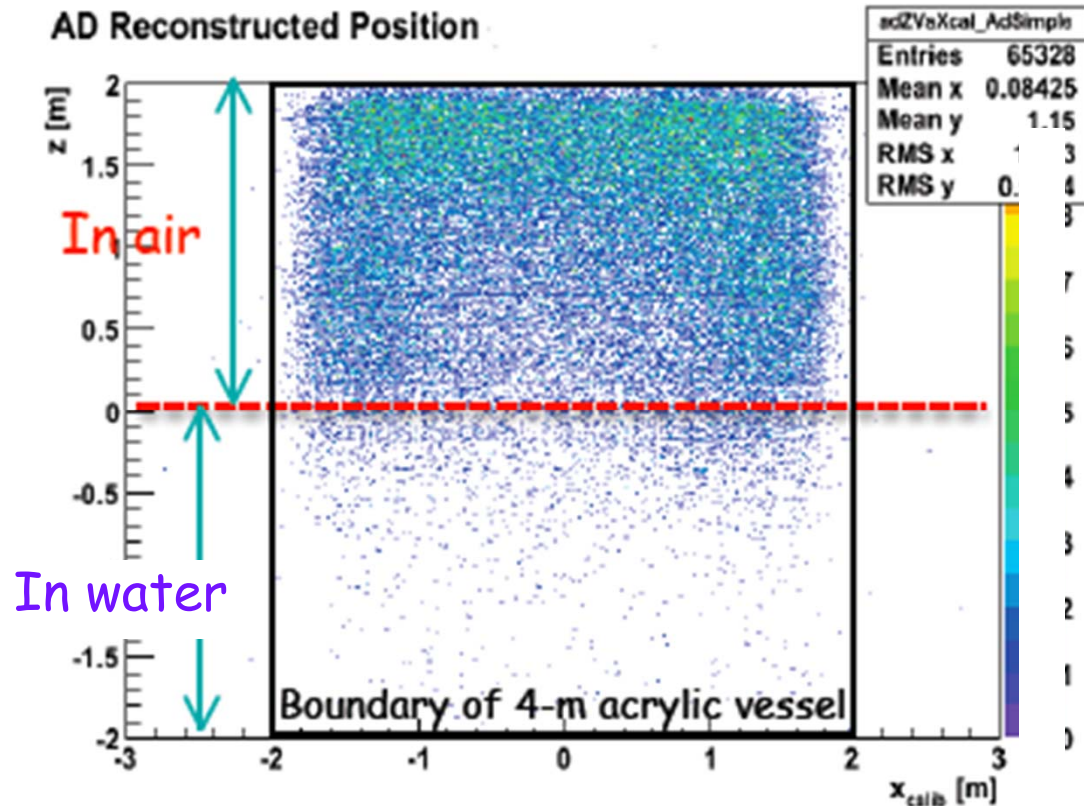
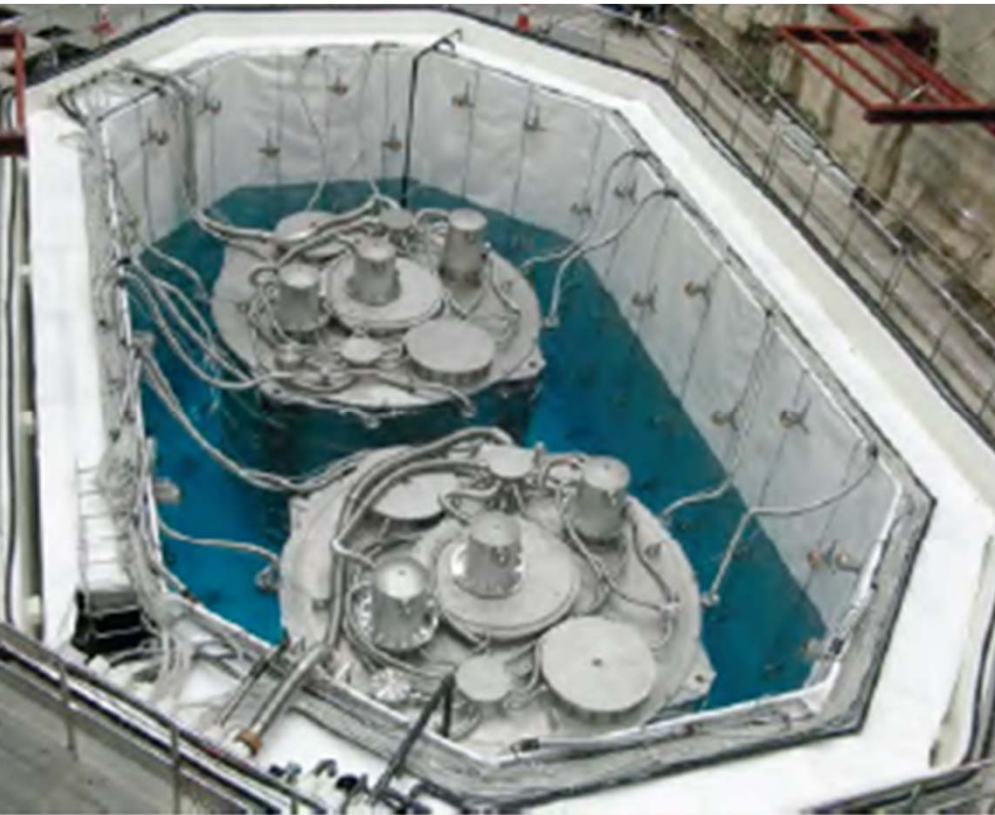
The Muon Tagging System

Dual tagging systems: 2.5-m thick, two-zone water shield, instrumented with PMTs, + resistive plate chambers (RPCs) above the water pool.



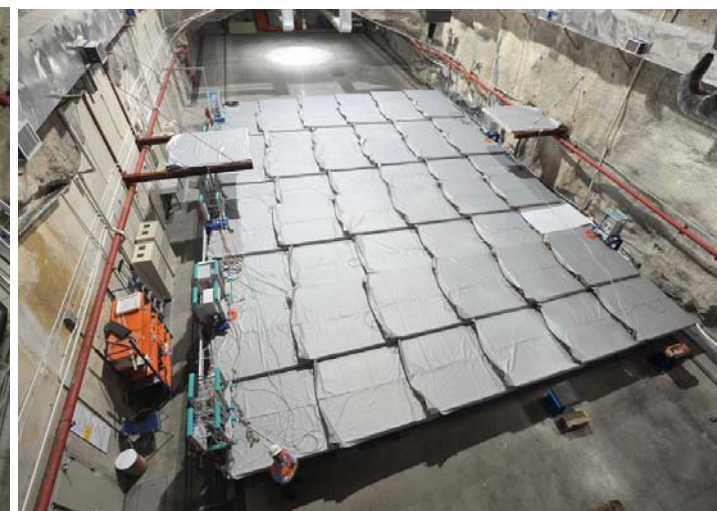
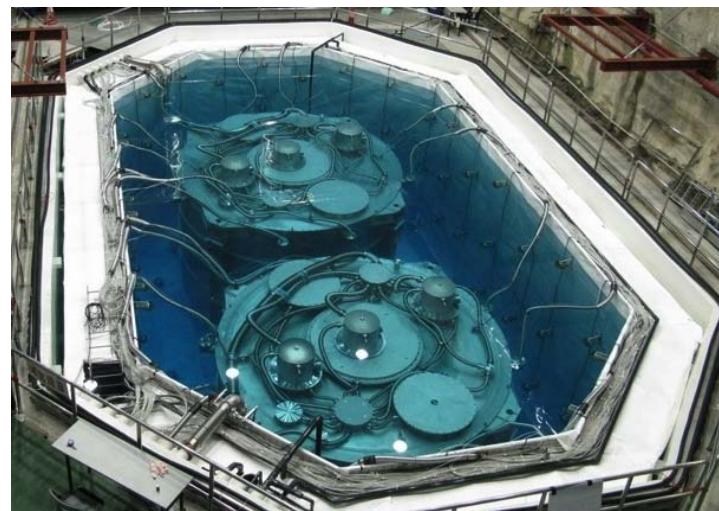
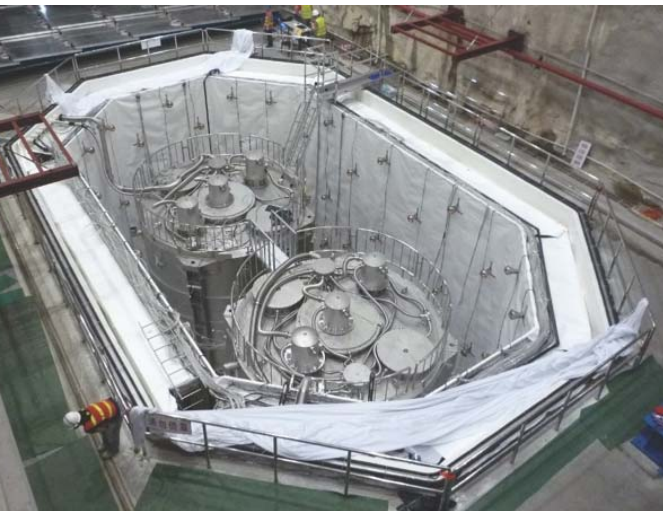
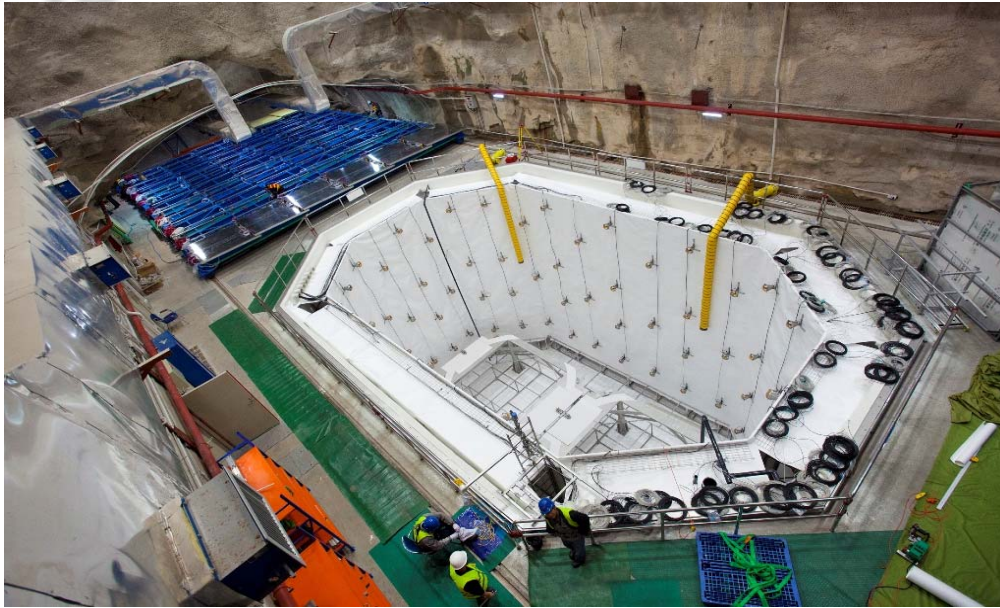
- No online veto, but rather offline tagging/veto of cosmic-ray muon-induced events.
- Design tagging efficiency $> 99.5\%$ with uncertainty $< 0.25\%$.
- Princeton hardware contribution is to the gas system for the RPCs.

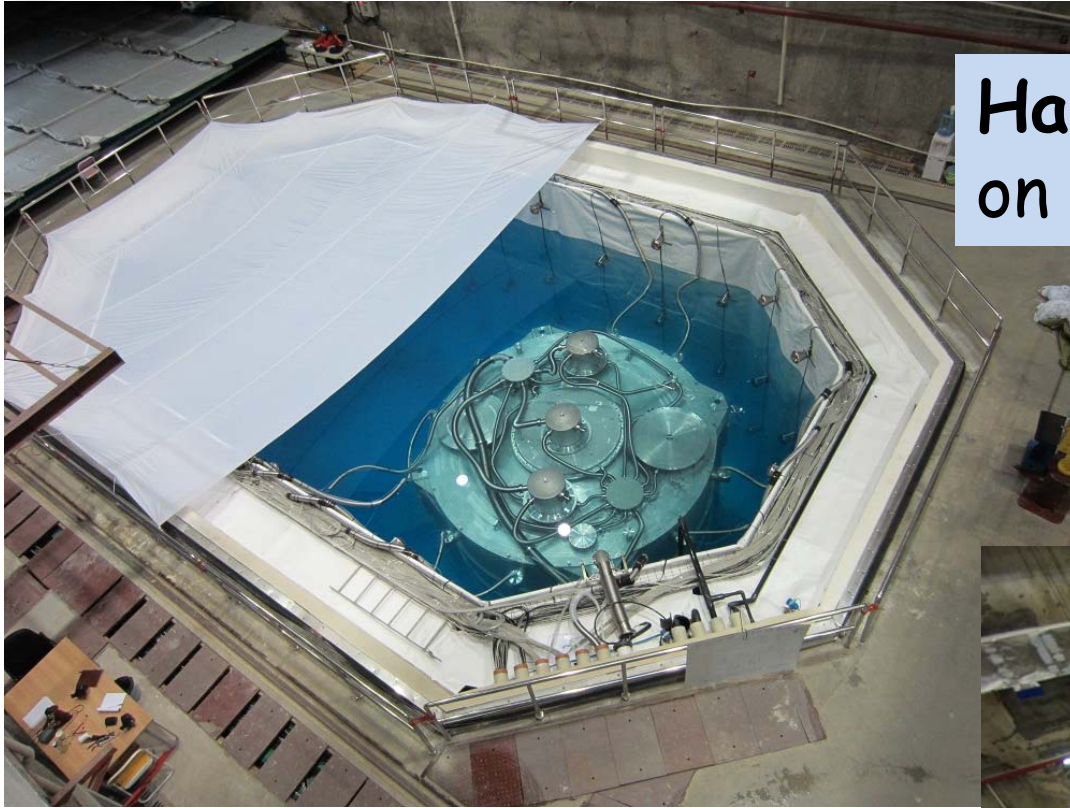




Singles rates vs. height in the partially filled pool show the suppression of radioactive backgrounds (due to the rock and/or radon) by the water shield.

[PMT coverage in the water pool is poor at the top, where the RPC array augments the muon coverage - but not that for rock radioactivity.]





Hall 2: Began 1 AD operation on Nov. 5, 2011

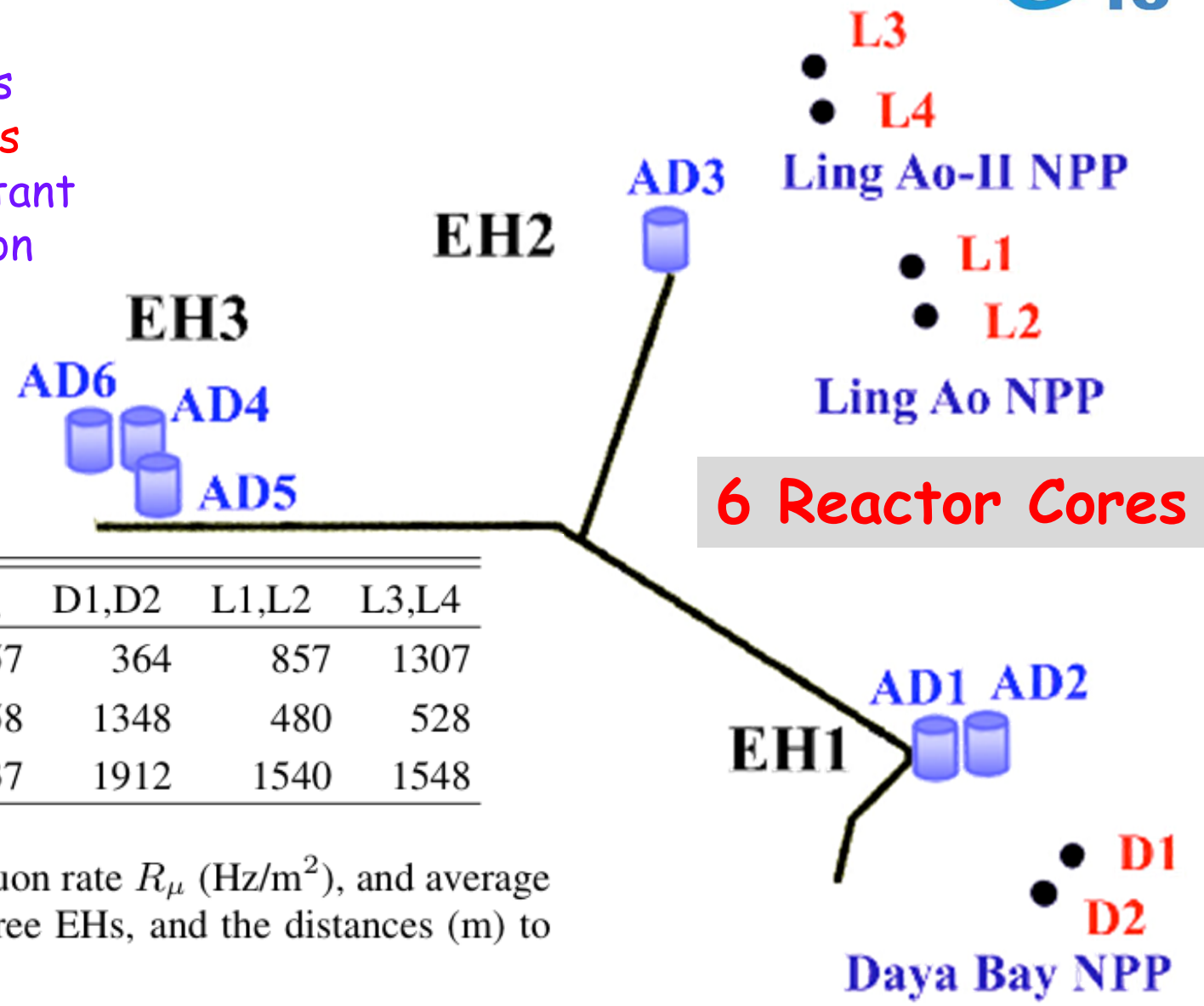
Hall 3: Began 3 AD operation on Dec. 24, 2011



2 more ADs still in assembly; installation planned for Summer 2012

6 Antineutrino Detectors in 3 Halls

EH1 is close to 2 reactors
 EH2 is close to 4 reactors
 ⇒ EH2 much more important
 in the near/far comparison



	Overburden	R_μ	E_μ	D1,D2	L1,L2	L3,L4
EH1	280	1.27	57	364	857	1307
EH2	300	0.95	58	1348	480	528
EH3	880	0.056	137	1912	1540	1548

TABLE I. Overburden (m.w.e), muon rate R_μ (Hz/m²), and average muon energy E_μ (GeV) of the three EHs, and the distances (m) to the reactor pairs.

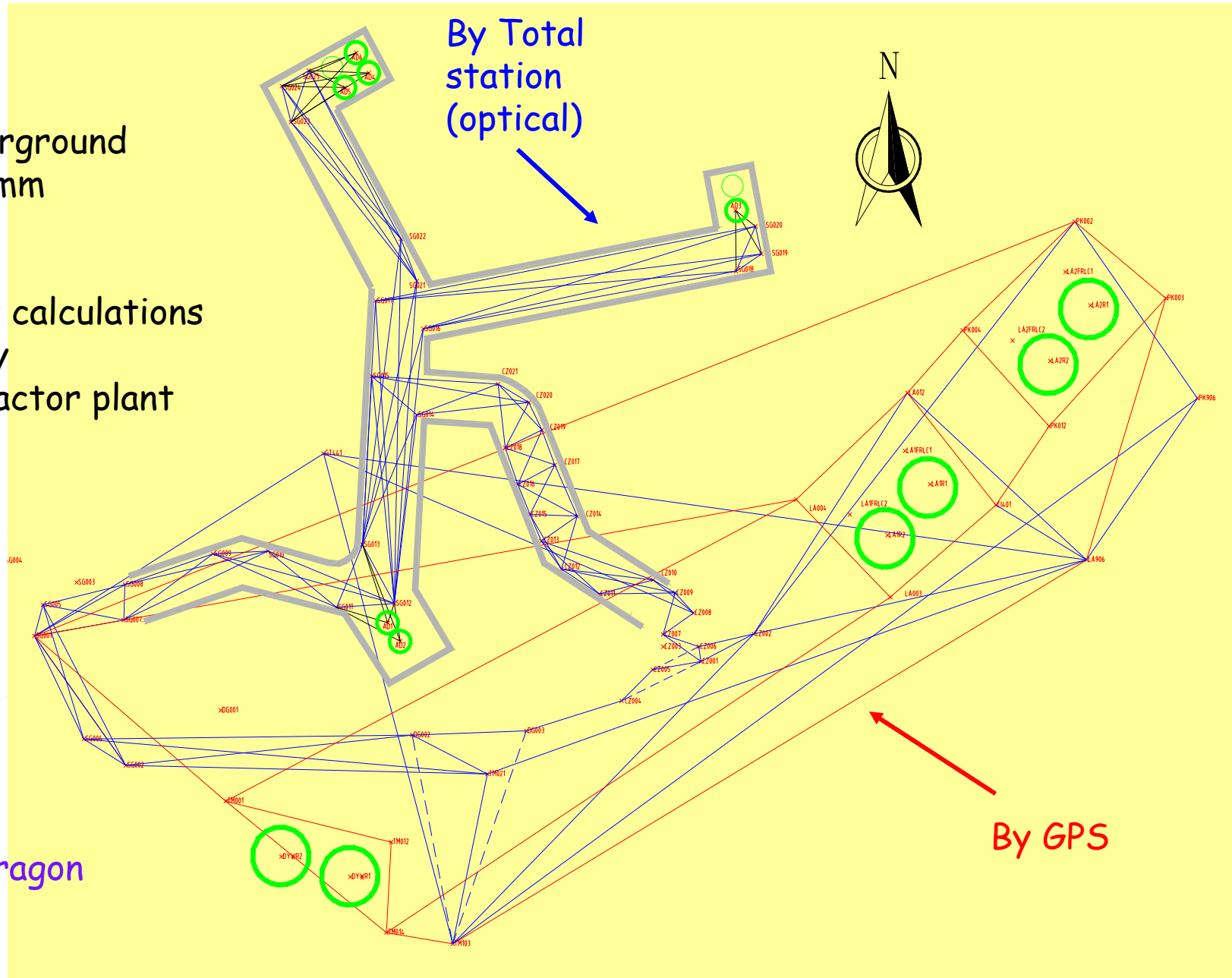


Detailed Survey:

- GPS above ground
- Total Station underground
- Final precision: 28mm

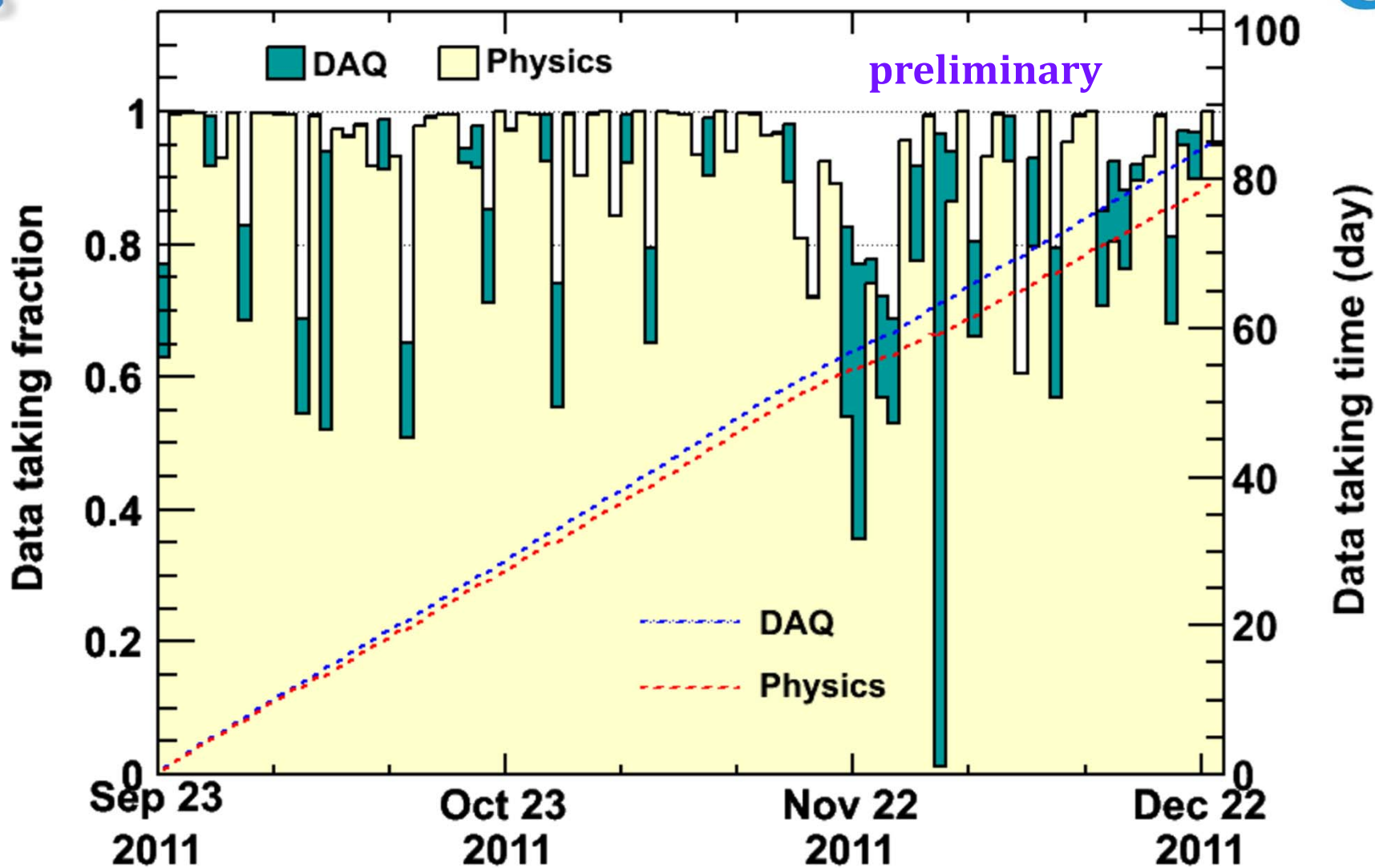
Validation:

- Three independent calculations
- Cross-check survey
- Consistent with reactor plant and design plans



2012 = year of the dragon

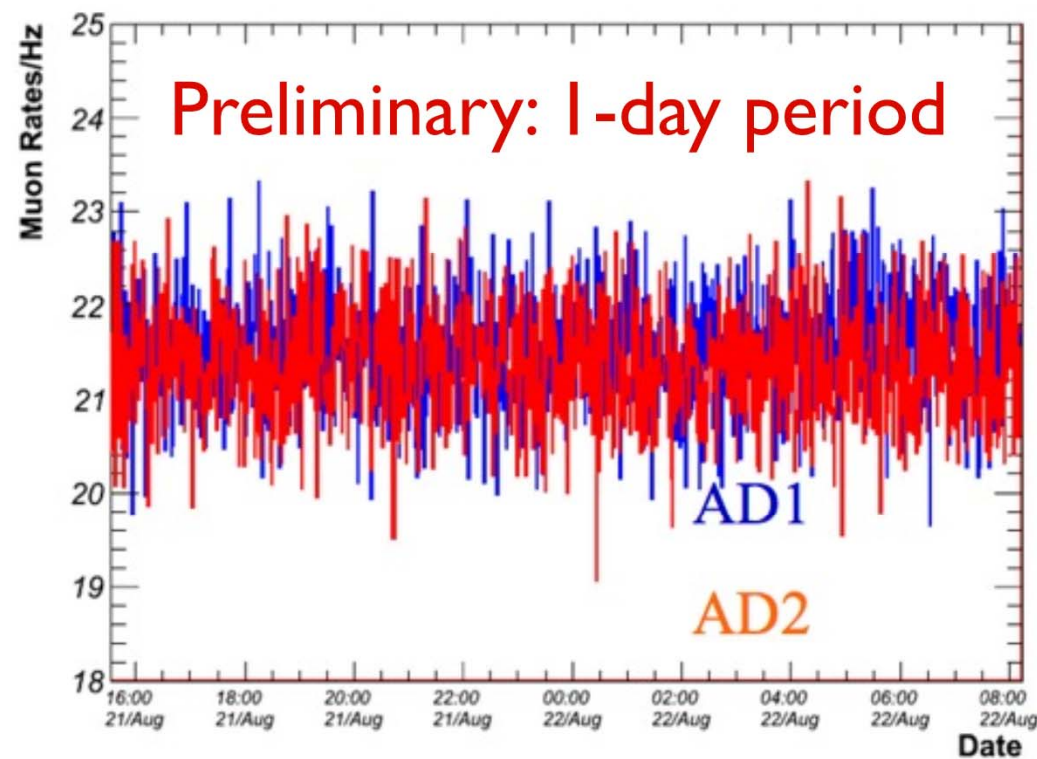
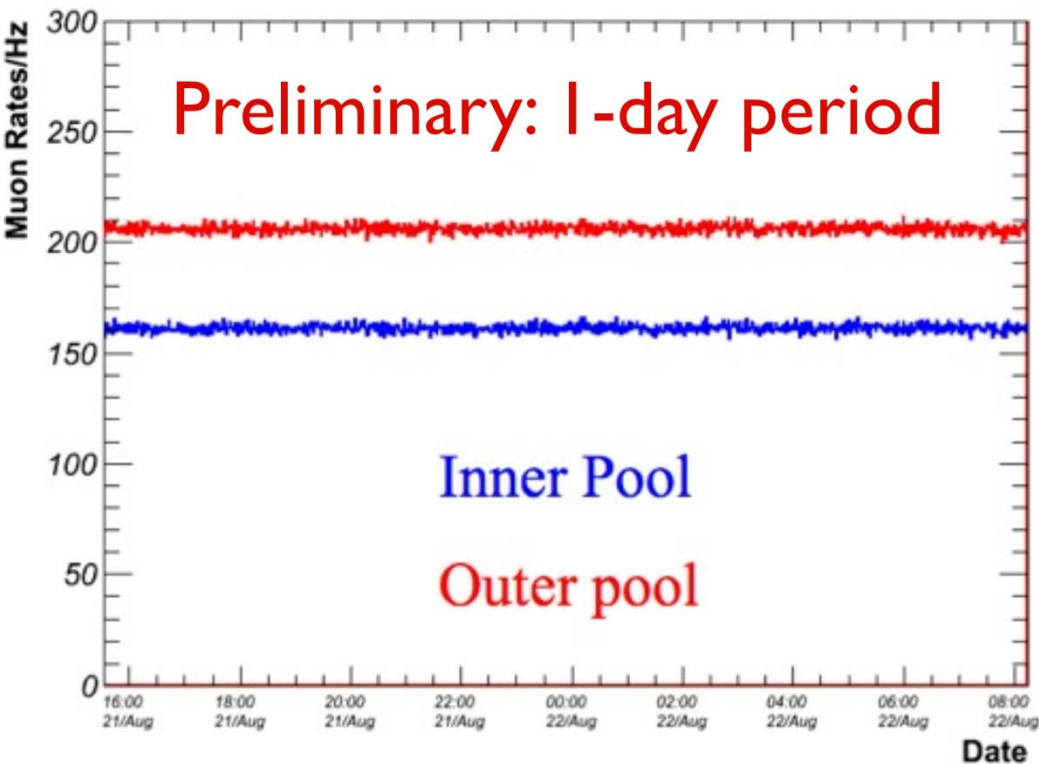




We have used the 3 months of data for side-by-side comparison of first two detectors.

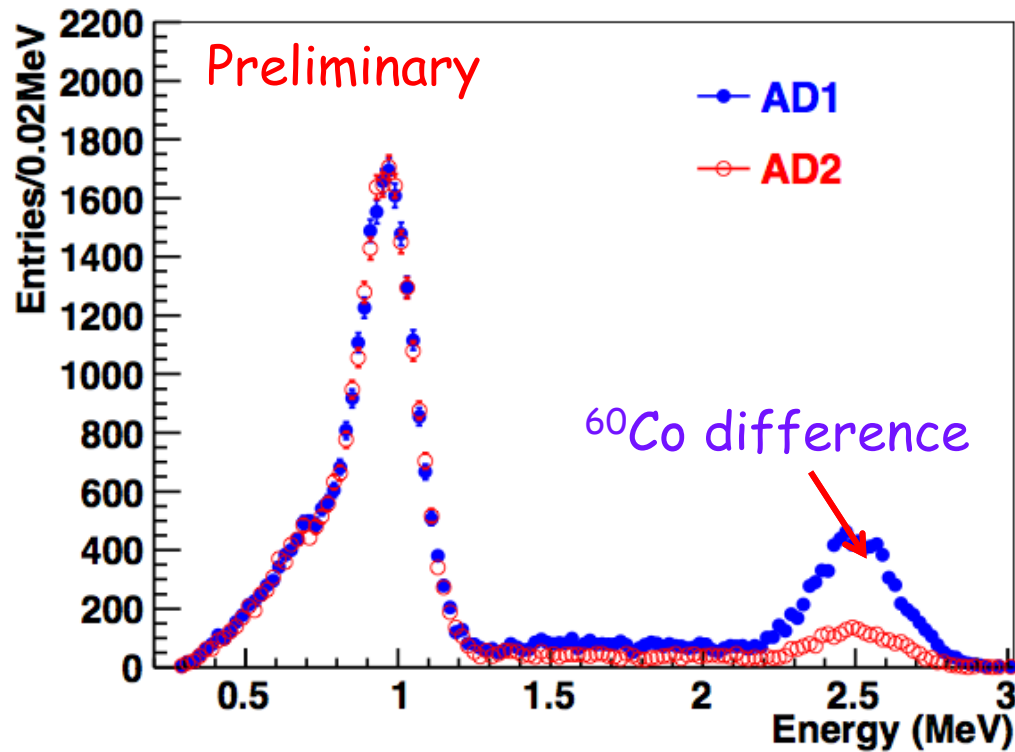
Detailed comparison of AD1 and 2: F.P. An *et al.*, <http://arxiv.org/abs/1202.6181>



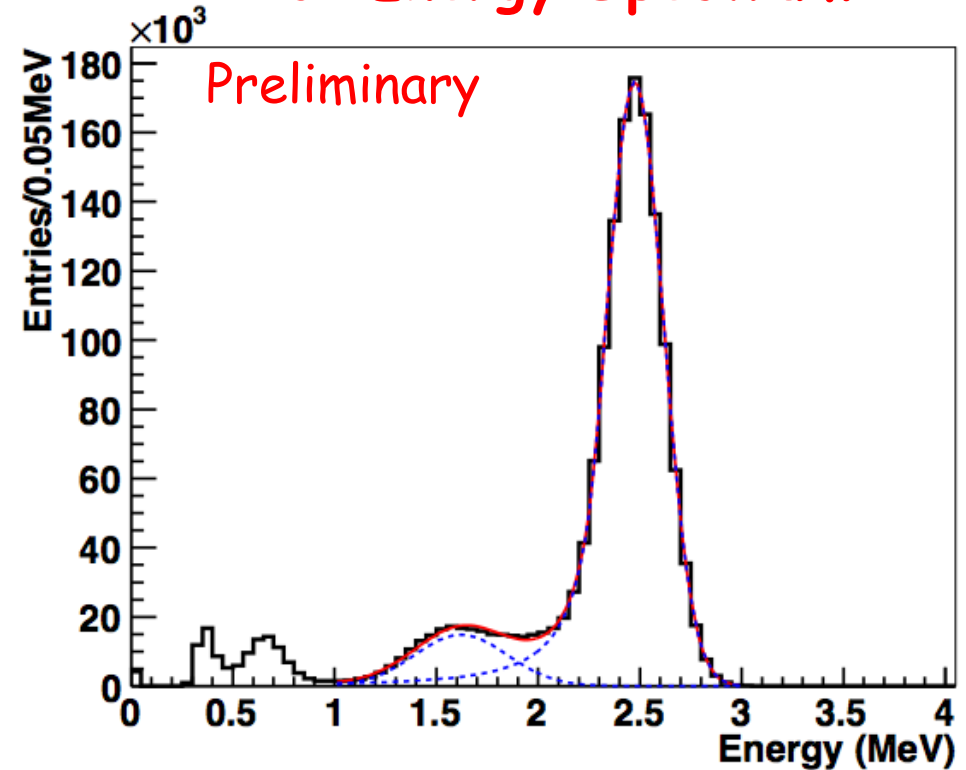


- Measured rates are consistent with expected ~ 20 Hz for each AD at the Daya Bay near site (EH1).

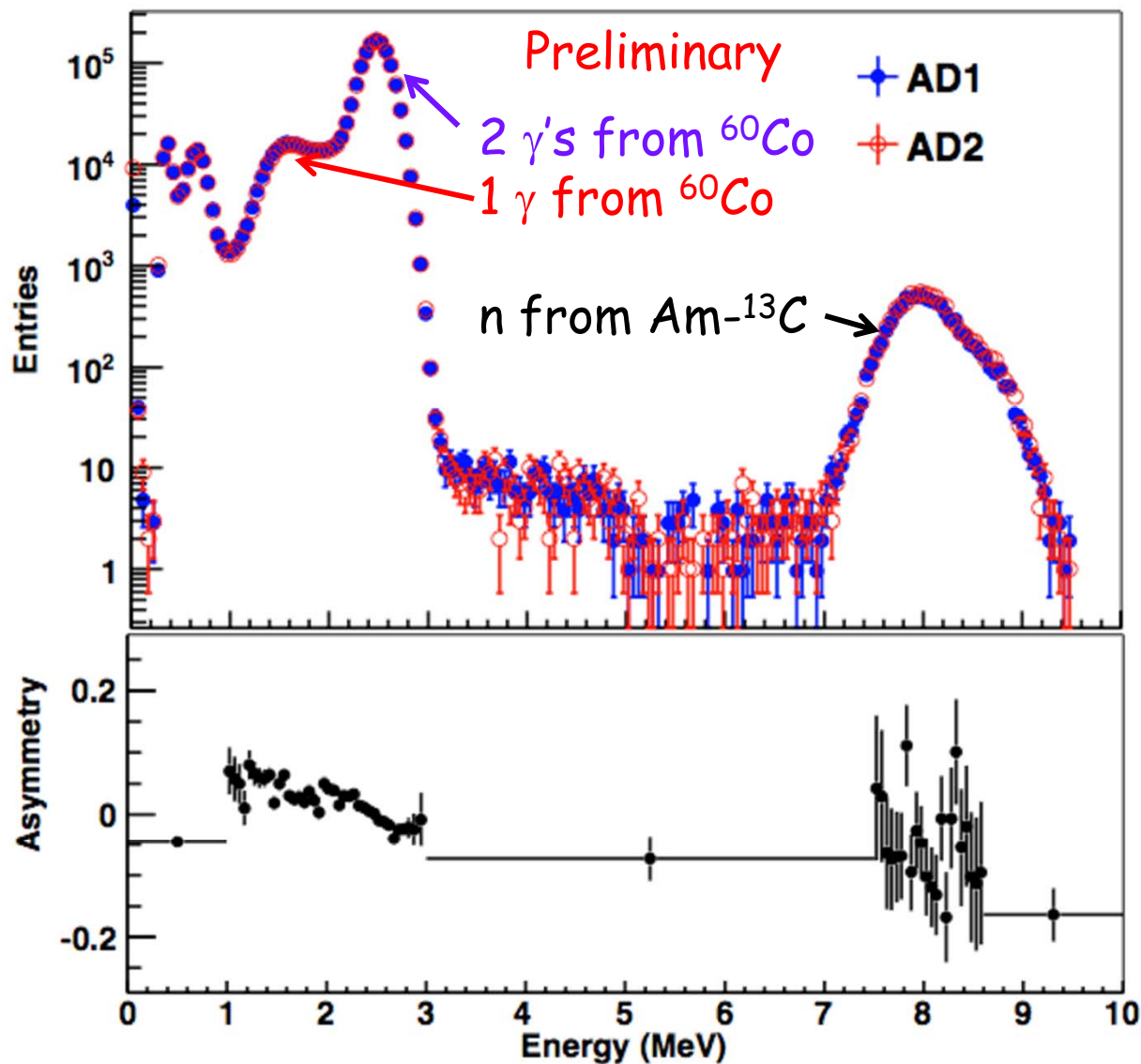
^{68}Ge Energy Spectrum



^{60}Co Energy Spectrum

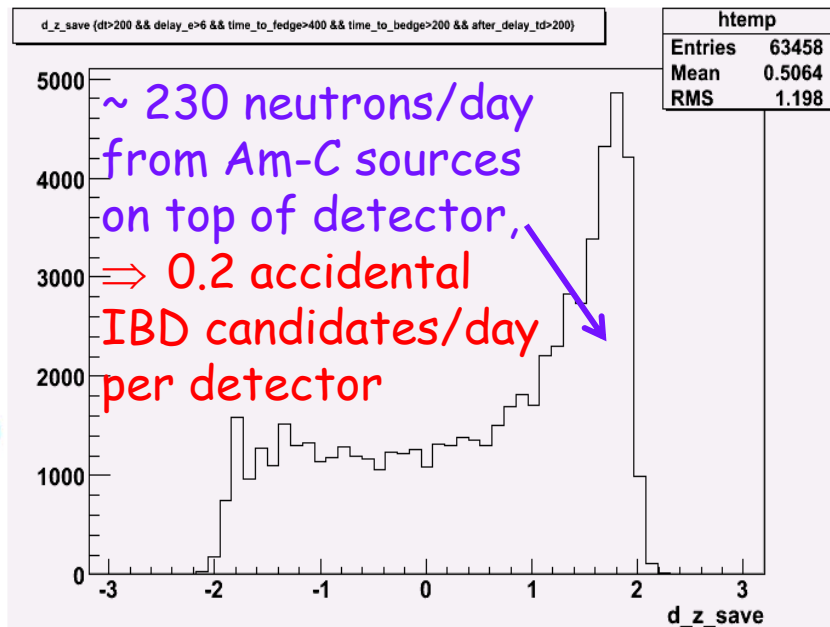
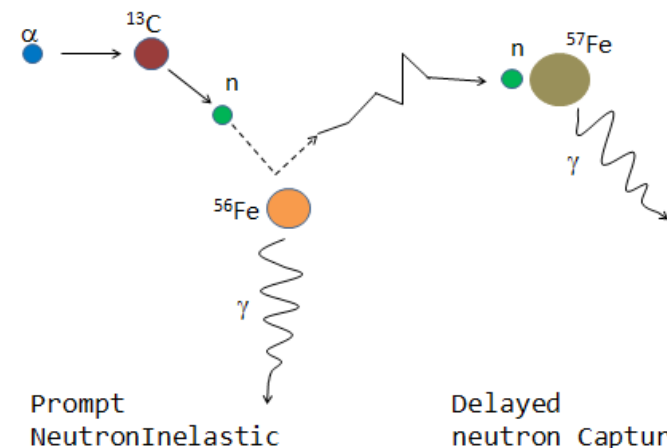


- Weekly automatic detector calibration.
- Energy responses of the detectors are studied with ^{68}Ge , ^{60}Co and $\text{Am-}^{13}\text{C}$.

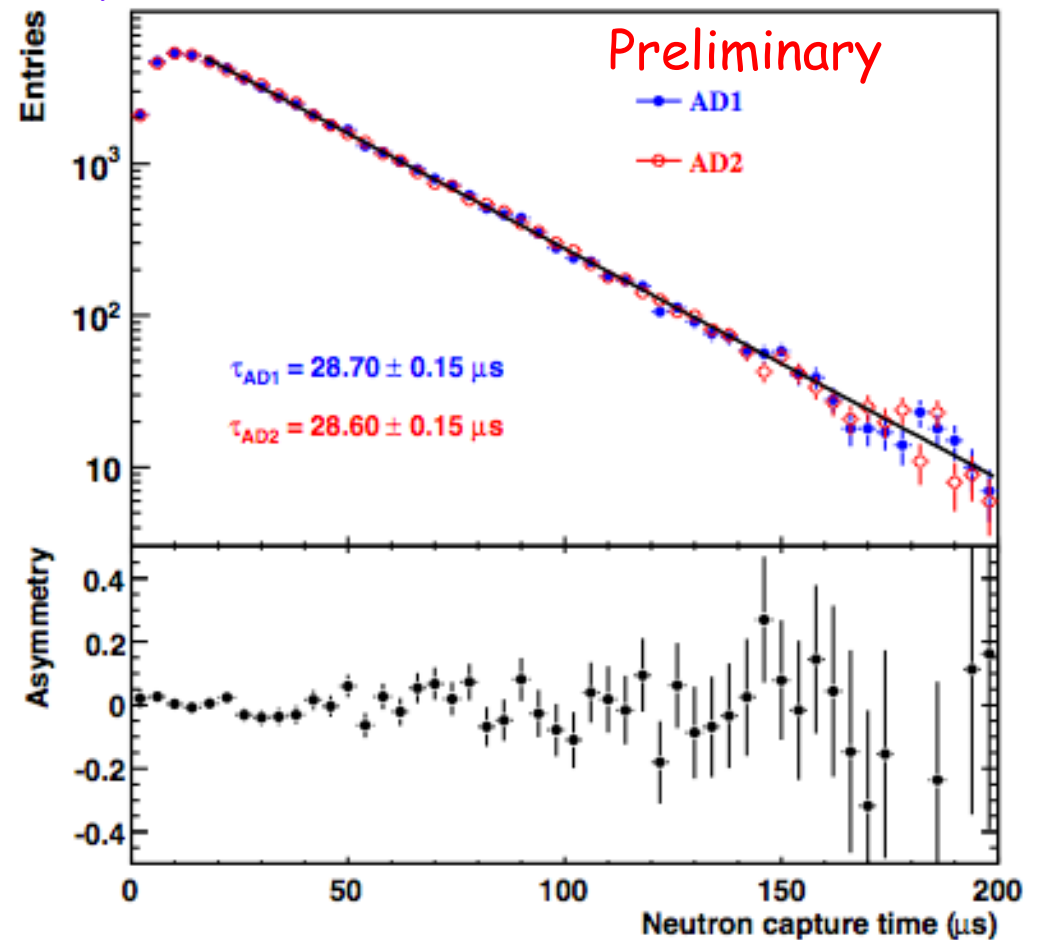
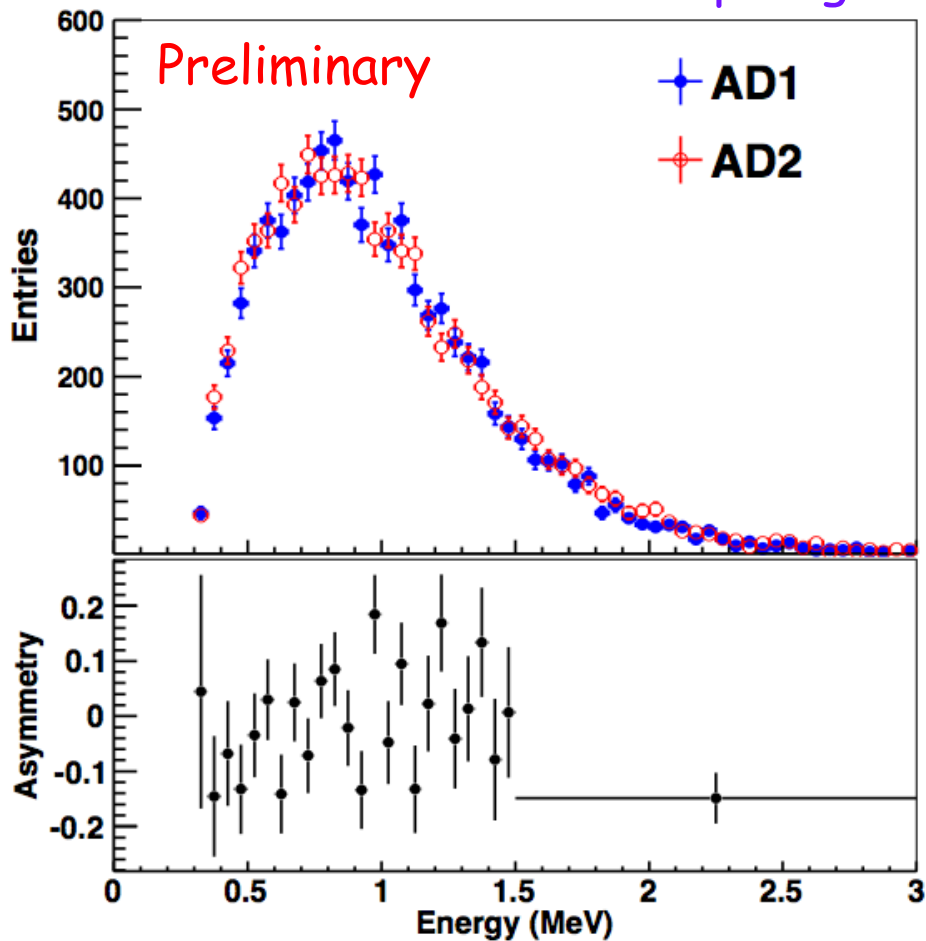


$$\text{Asymmetry} = \frac{N_{AD1} - N_{AD2}}{(N_{AD1} + N_{AD2})/2}$$

$N_{AD1,2}$ is the bin content for AD1 or 2

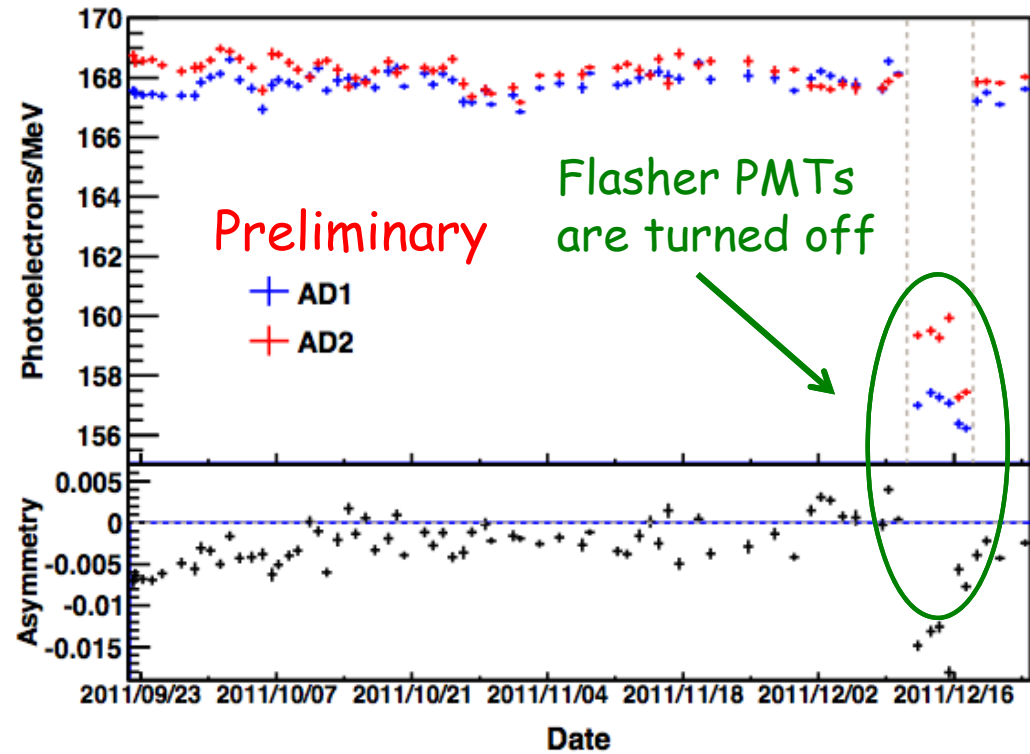
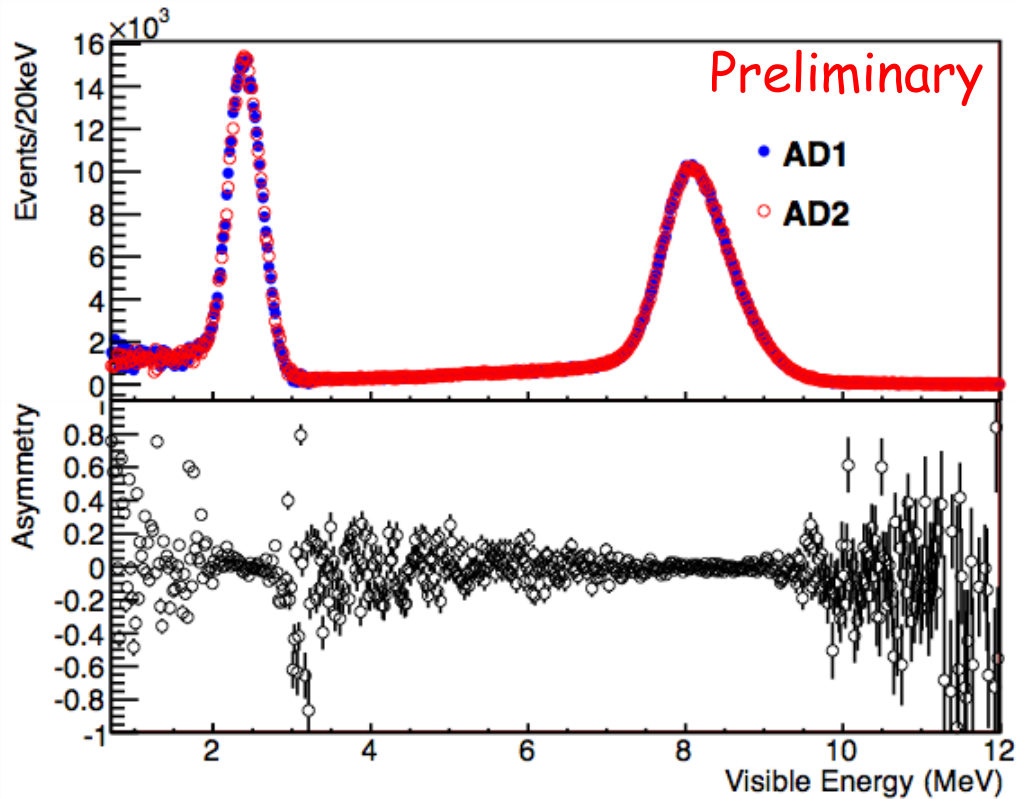


Prompt signal from proton recoil



- Two ADs with similar energy responses ($\sim 0.5\%$)
- Consistent response in capture time measurements



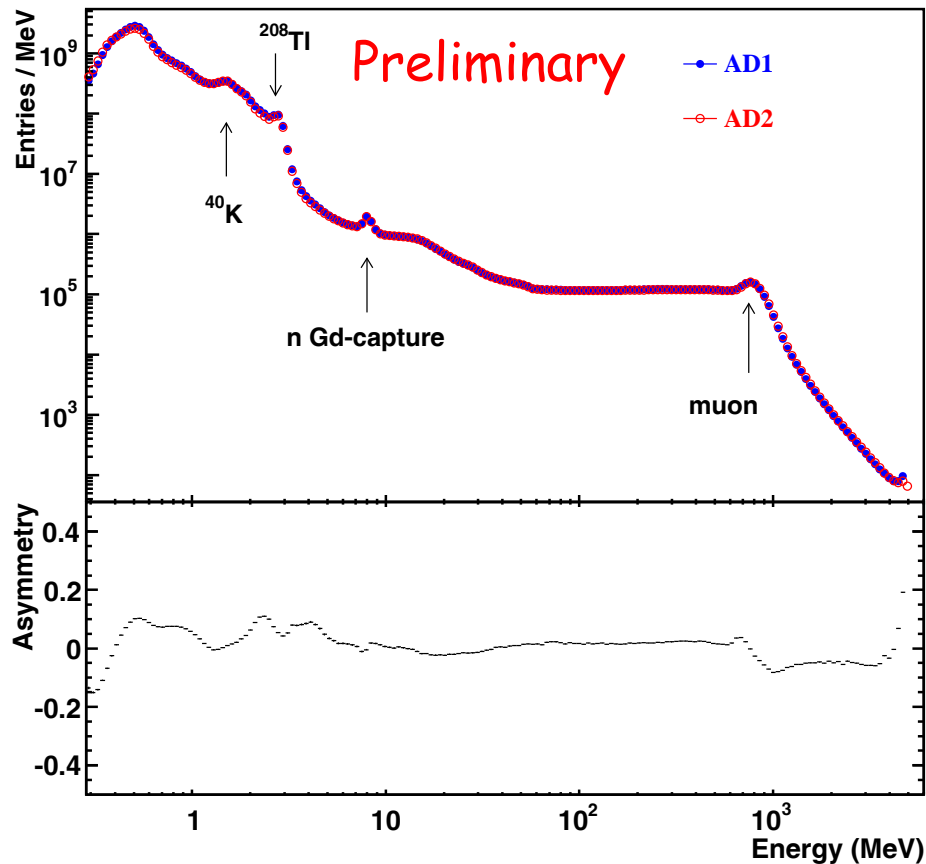


- Run-by-run calibration of detectors using spallation neutrons
 - ~ 168 p.e./MeV for AD1
 - ~ 169 p.e./MeV for AD2

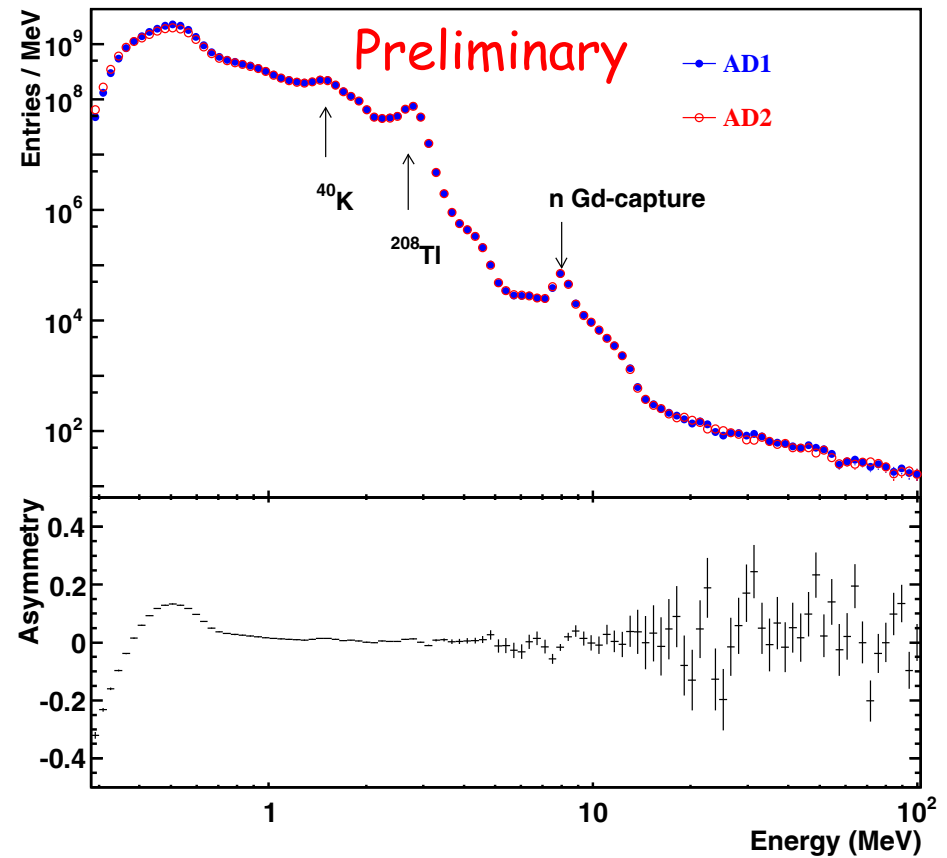
p.e. = photoelectrons detected in a photomultiplier tube.



Before muon (offline) veto,
 ~ 65 kHz with $E > 0.7$ MeV



After muon veto,
 \Rightarrow low-energy events due to radioactivity (mostly in PMTs)



- The difference in AD1 and AD2 triggers is mostly due to after-pulsing in the PMTs immediately following a muon event (which leads to a relatively large signal).



Reject "flashers" = emission of light by a PMT, due to secondary emission processes in its dynodes during photoelectron signal amplification.

Prompt "positron": $0.7 \text{ MeV} < E_p < 12 \text{ MeV}$

Delayed "neutron": $6.0 \text{ MeV} < E_d < 12 \text{ MeV}$

Neutron capture time: $1 \mu\text{s} < \Delta t < 200 \mu\text{s}$

Muon Veto:

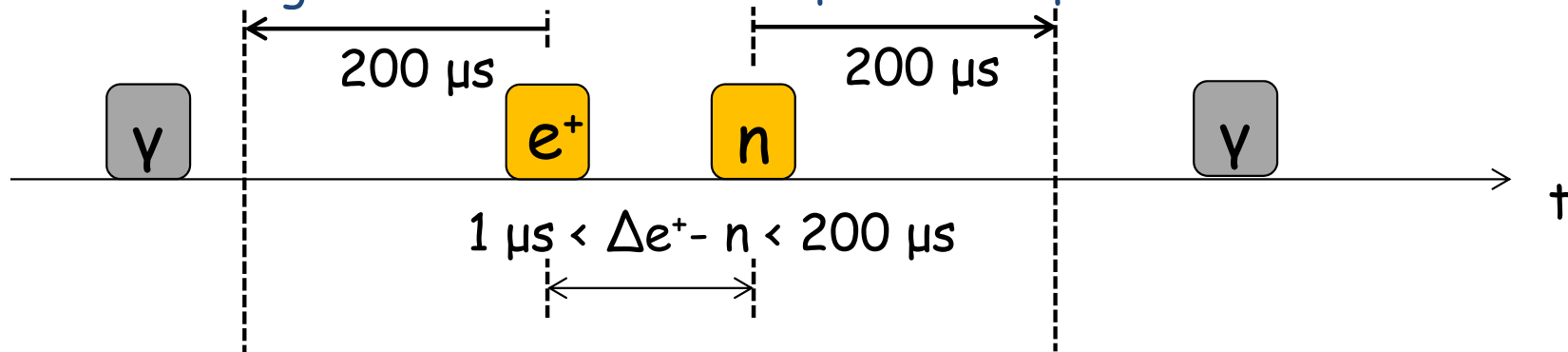
Pool Muon: Reject IBD candidate if $< 0.6 \text{ ms}$ after pool muon

AD Muon ($> 20 \text{ MeV}$): Reject if $< 1 \text{ ms}$ before IBD candidate

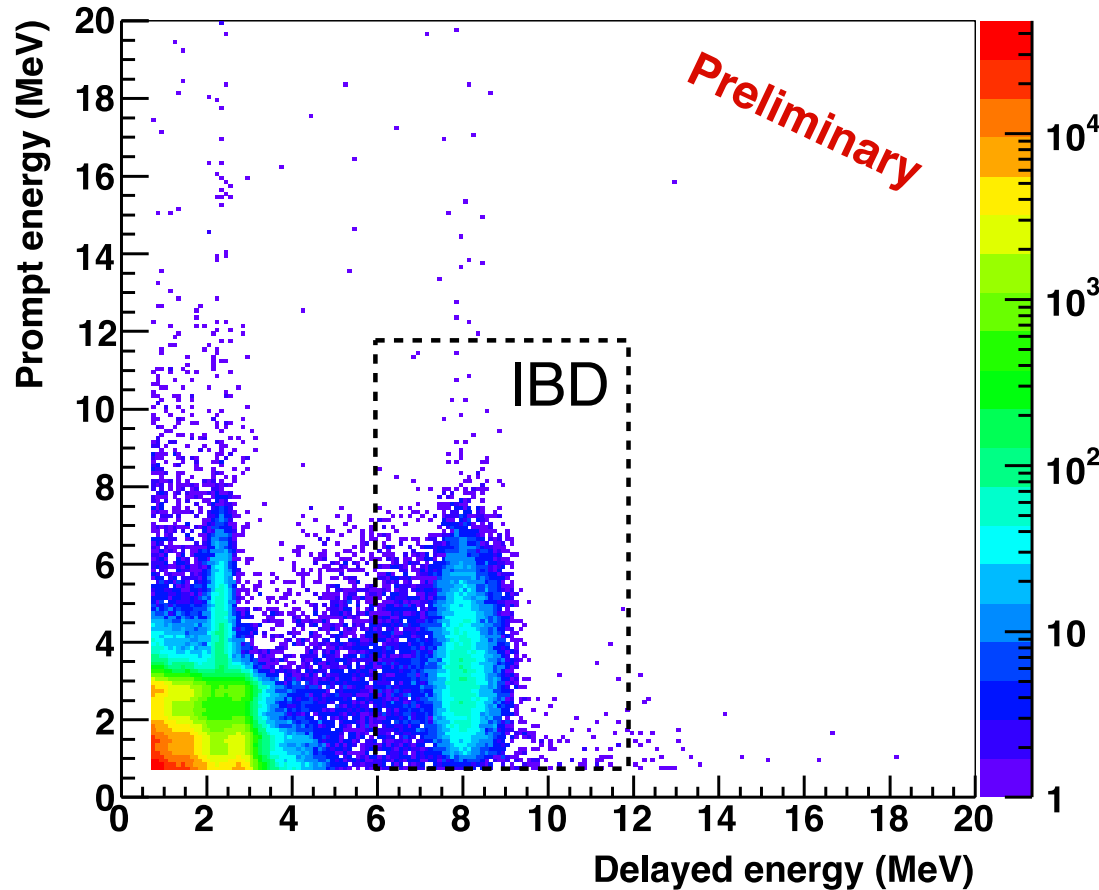
AD Shower Muon ($> 2.5 \text{ GeV}$): Reject if $< 1 \text{ s}$ before

Multiplicity cut:

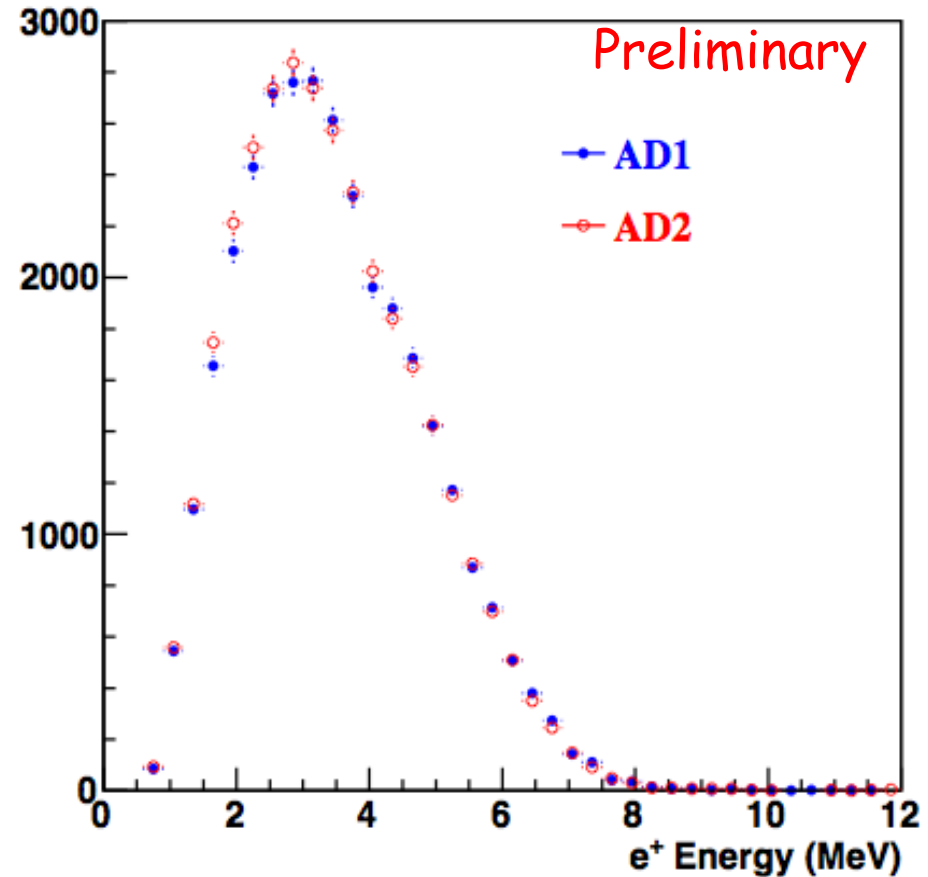
No other AD signal $> 0.7 \text{ MeV}$ in $-200 \mu\text{s}$ to $200 \mu\text{s}$ of IBD.



IBD = Inverse beta decay = $\bar{\nu} + p \rightarrow n + e$



Prompt Energy Spectrum



With Flasher cut, Muon veto, Prompt Energy, Delayed Energy, Time correlation and Multiplicity cuts.



A PMT "flash" leads to light collected near that PMT, and on the opposite side of the detector.

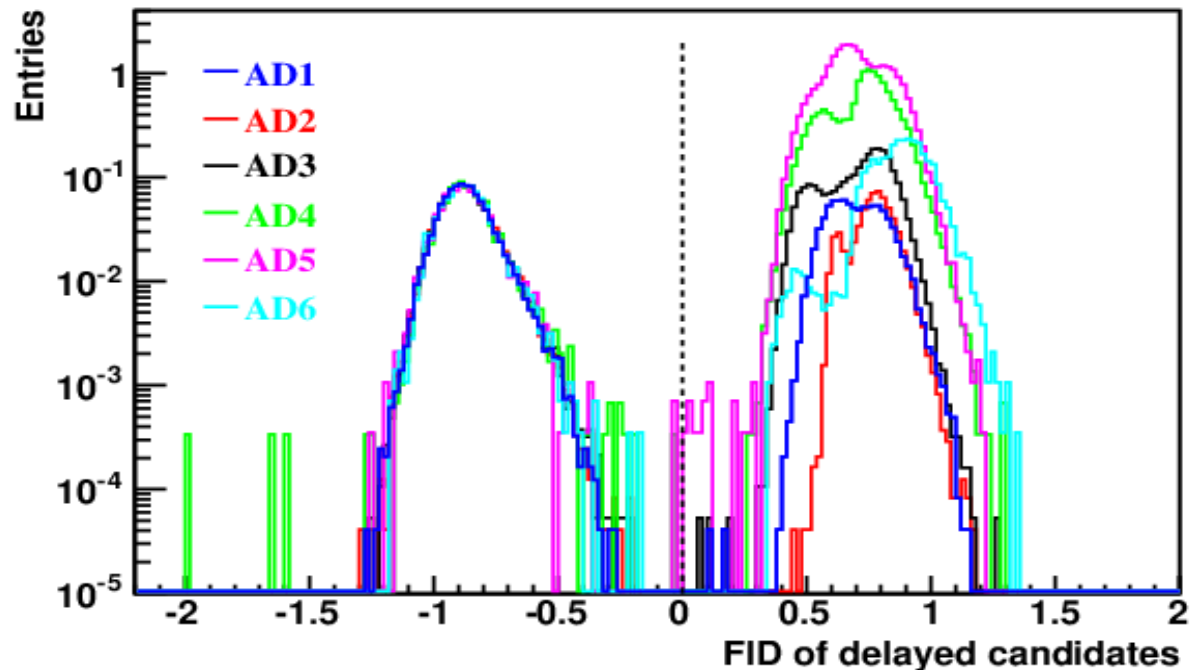
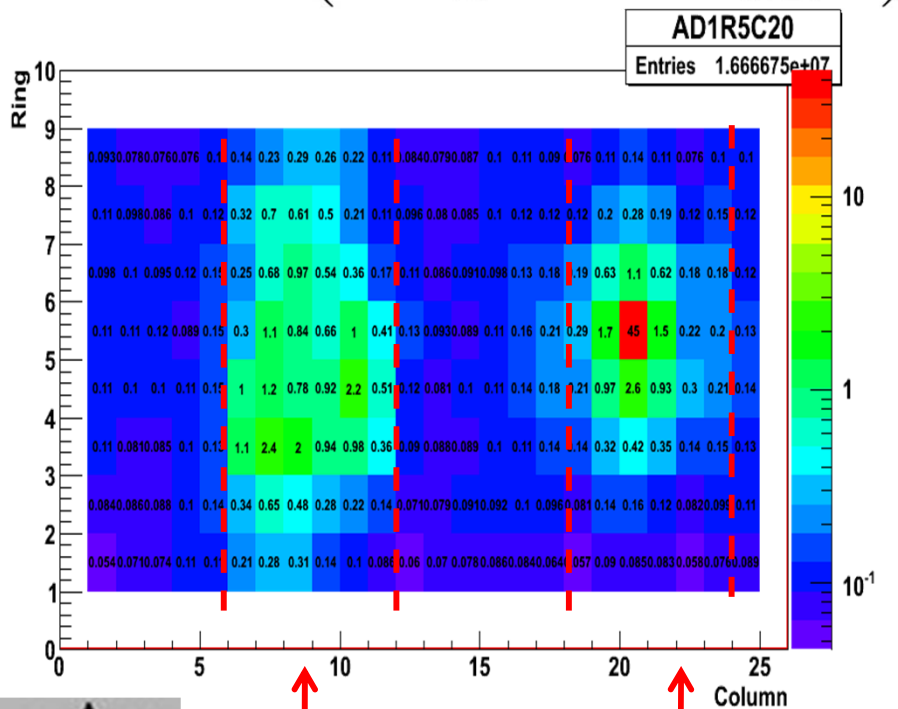
Q = charge collected from PMT anode

$MaxQ = Q_{max} / \sum Q$

Define Quadrant 1 as that which contains the "hottest" PMT

$Quadrant = Q_{Quadrant3} / (Q_{Quadrant2} + Q_{Quadrant4})$

$FID = \log \left(\left(\frac{Quadrant}{1.} \right)^2 + \left(\frac{MaxQ}{0.45} \right)^2 \right) < 0$ for IBD candidates, > 0 for "flashers"



Quadrant 3

Quadrant 1 with the flasher PMT

4/24/2012

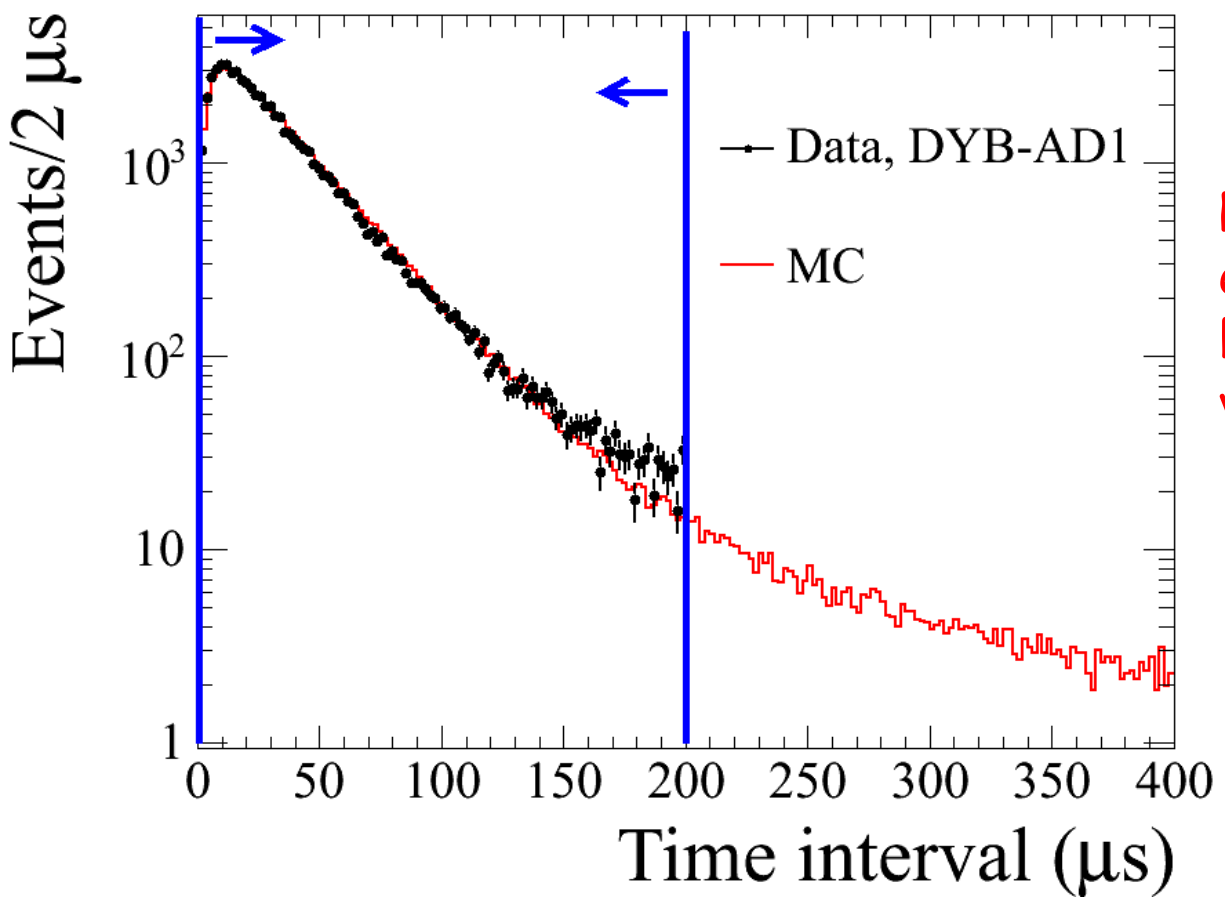
KT McDonald

Seminar at Oxford U

31

Neutron Capture Time

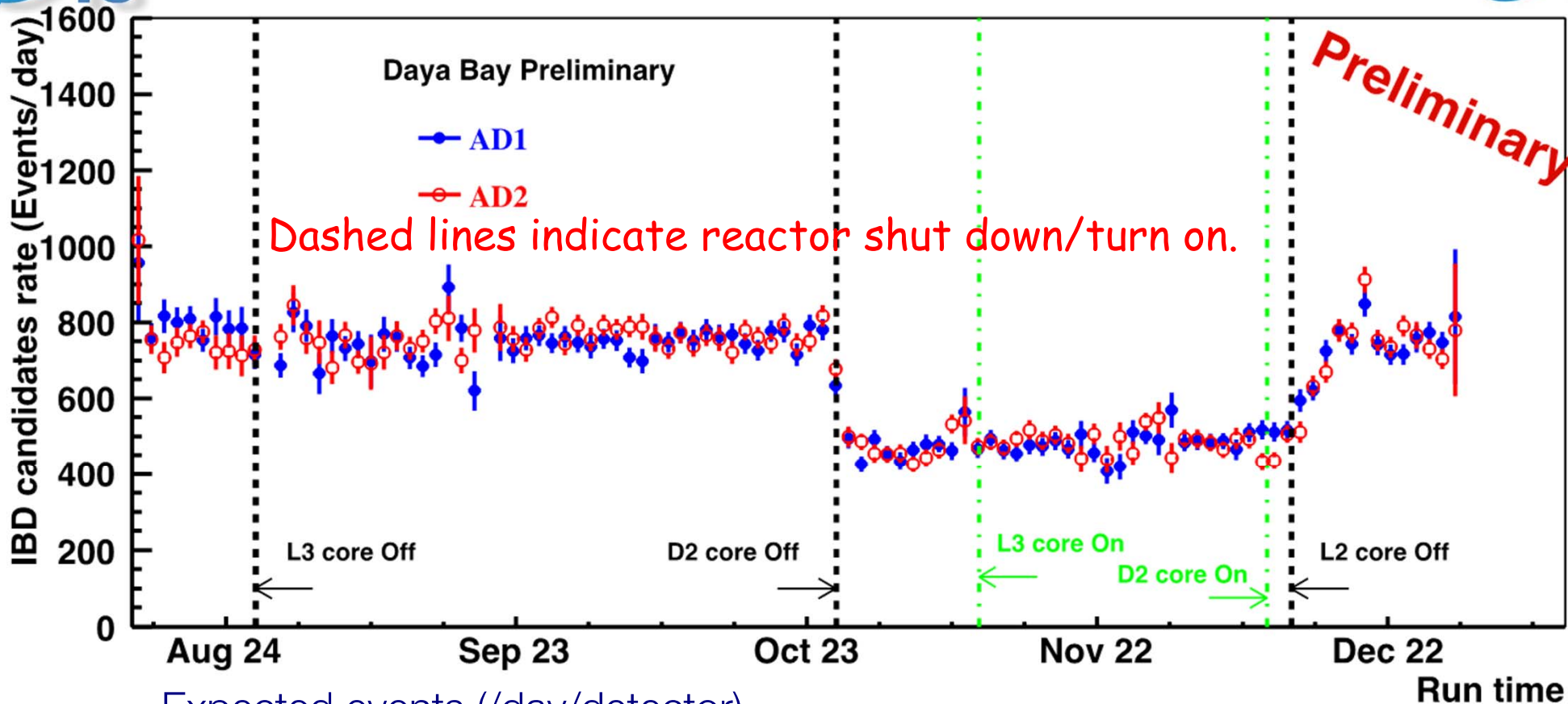
Consistent IBD capture time measured in all detectors



Relative detector efficiency estimated within 0.02% by considering possible variations in Gd concentration.

Simulation contains no background (deviates from data at $> 150 \mu\text{s}$)





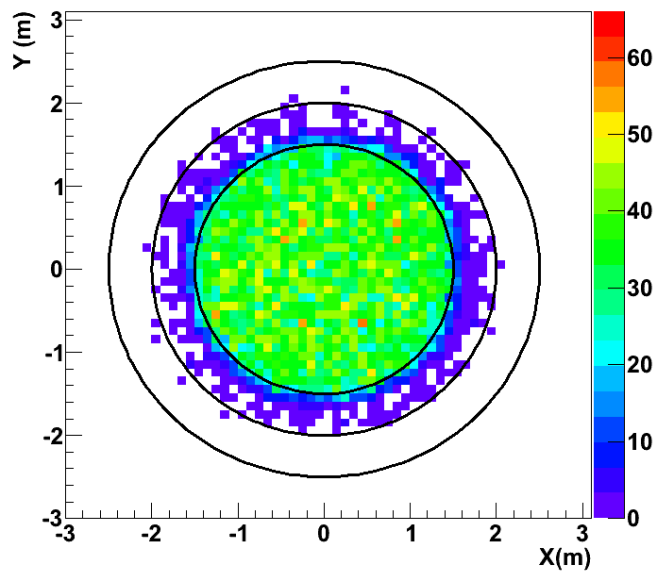
Expected events (/day/detector)

	DYB Site	LA Site	Far Site
IBD Evts	840	760	90
BKG Evts	<0.6%	<0.5%	<0.4%

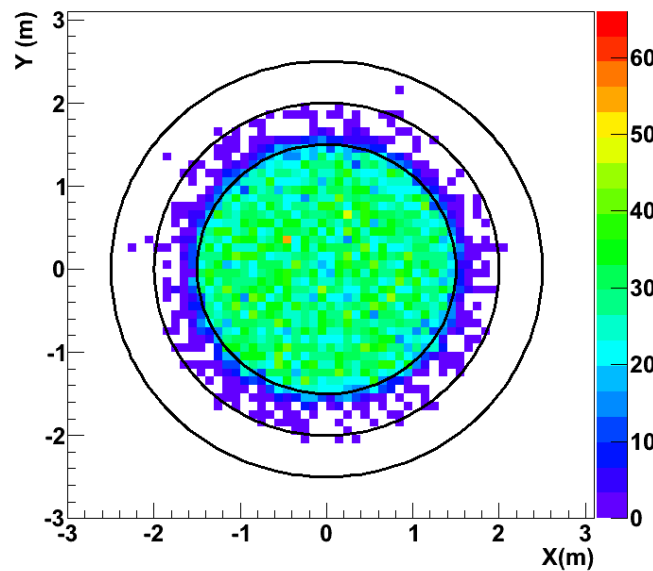
Expected events with all cores ON



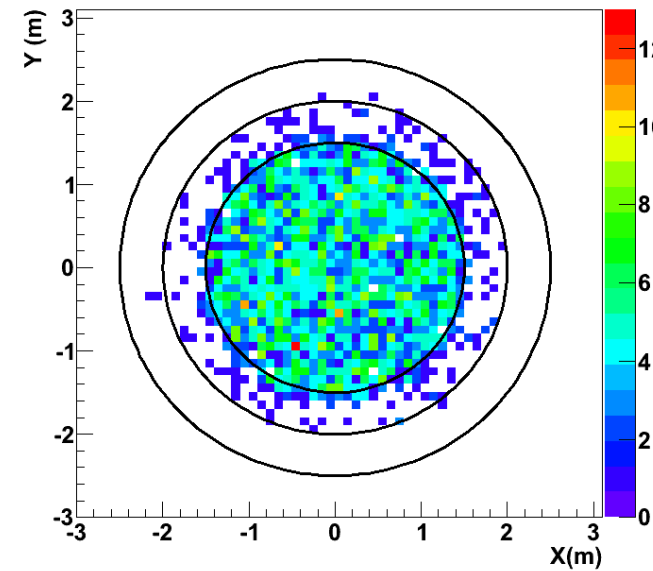
AD1



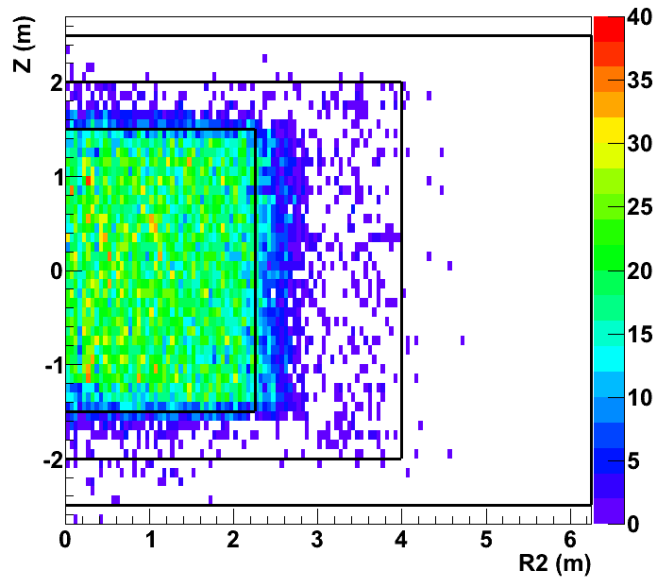
AD3



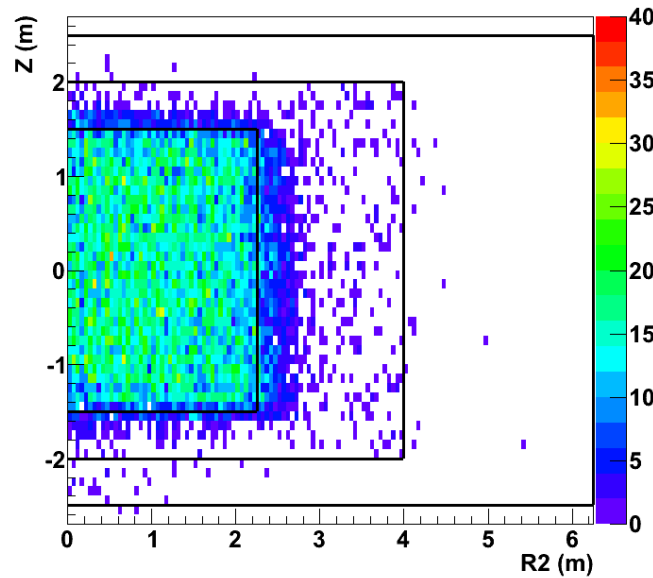
AD6



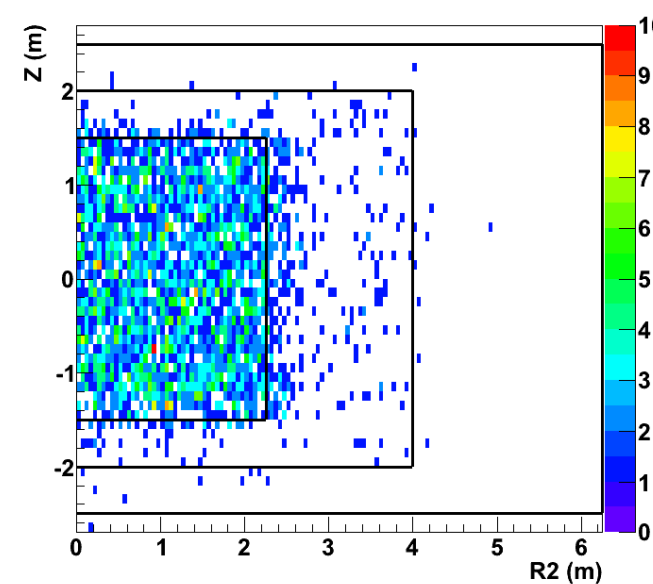
AD1



AD3

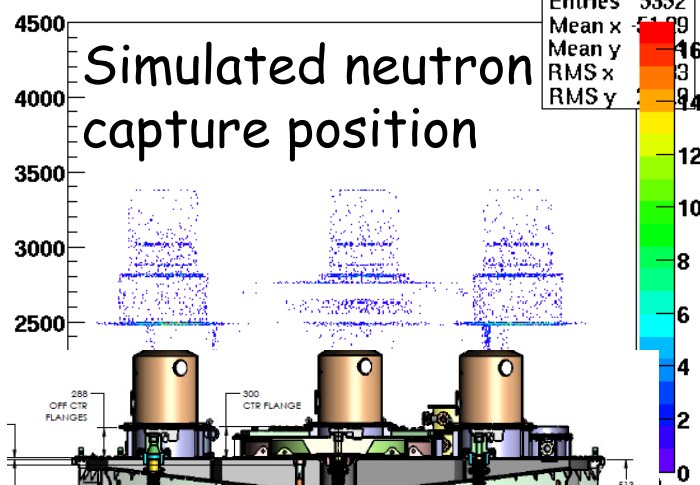


AD6



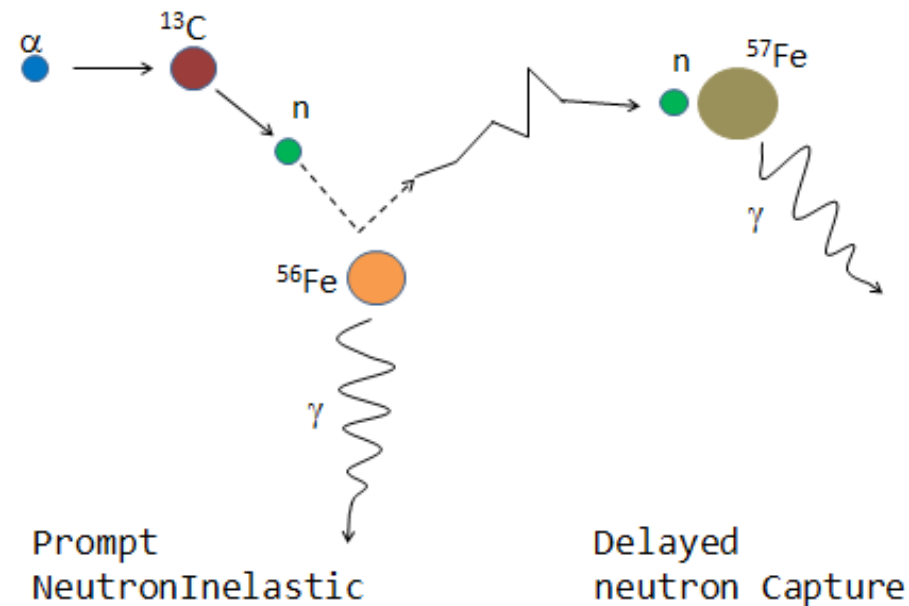
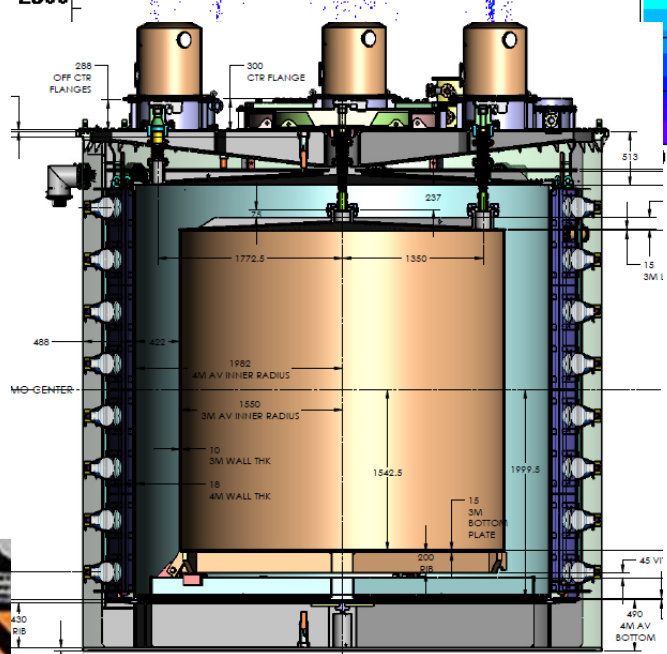
Weak (0.5 Hz) neutron calibration source in ACU can mimic IBD via inelastic scattering and capture on iron.

All Three ACU Capture Position z vs ξ (recE > 6 MeV)



Simulated neutron capture position

hCap	
Entries	5352
Mean x	51.99
Mean y	16
RMS x	13
RMS y	14

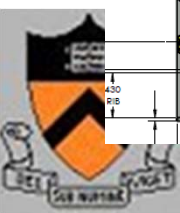


Prompt Neutron Inelastic

Delayed neutron Capture

Constrain far site B/S to $0.3 \pm 0.3\%$:

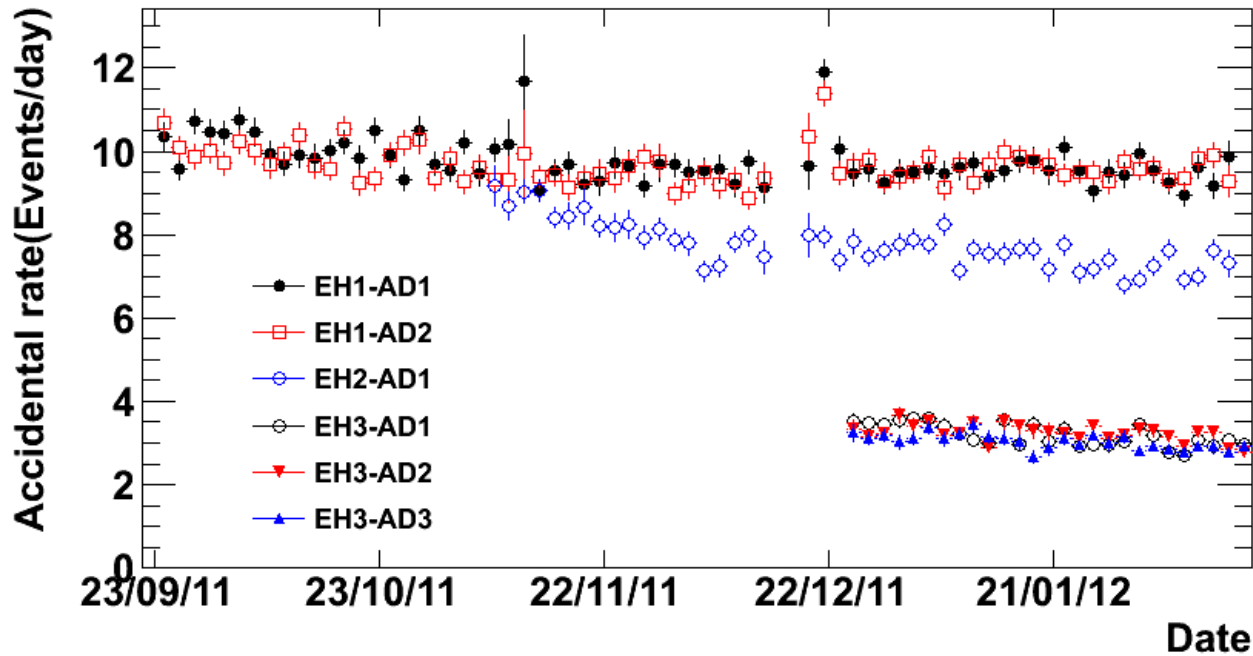
- Measure uncorrelated gamma rays from ACU in data
- Estimate ratio of correlated/uncorrelated rate using simulation
- Assume 100% uncertainty from simulation



Two random single signals can accidentally mimic a prompt-delayed IBD signal

Accidental rate and spectrum can be accurately predicted from singles data.

Multiple analyses/methods estimate consistent rates.



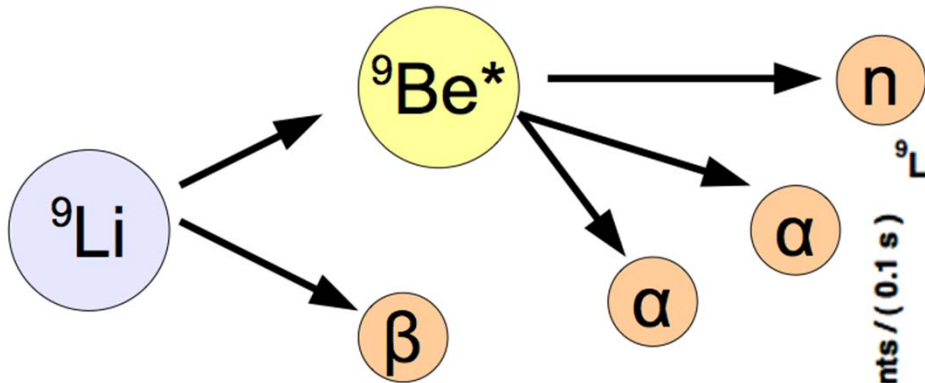
	EH1-AD1	EH1-AD2	EH2-AD1	EH3-AD1	EH3-AD2	EH3-AD3
Accidental rate(/day)	9.82 ± 0.06	9.88 ± 0.06	7.67 ± 0.05	3.29 ± 0.03	3.33 ± 0.03	3.12 ± 0.03
B/S	1.37%	1.38%	1.44%	4.58%	4.77%	4.43%

Lower rate of accidentals in EH3 is due to the lower rate of neutrons from cosmic-ray muons (although rate of neutrons from the retracted Am-C source is the same).



β -n decay:

- Prompt: β -decay
- Delayed: neutron capture



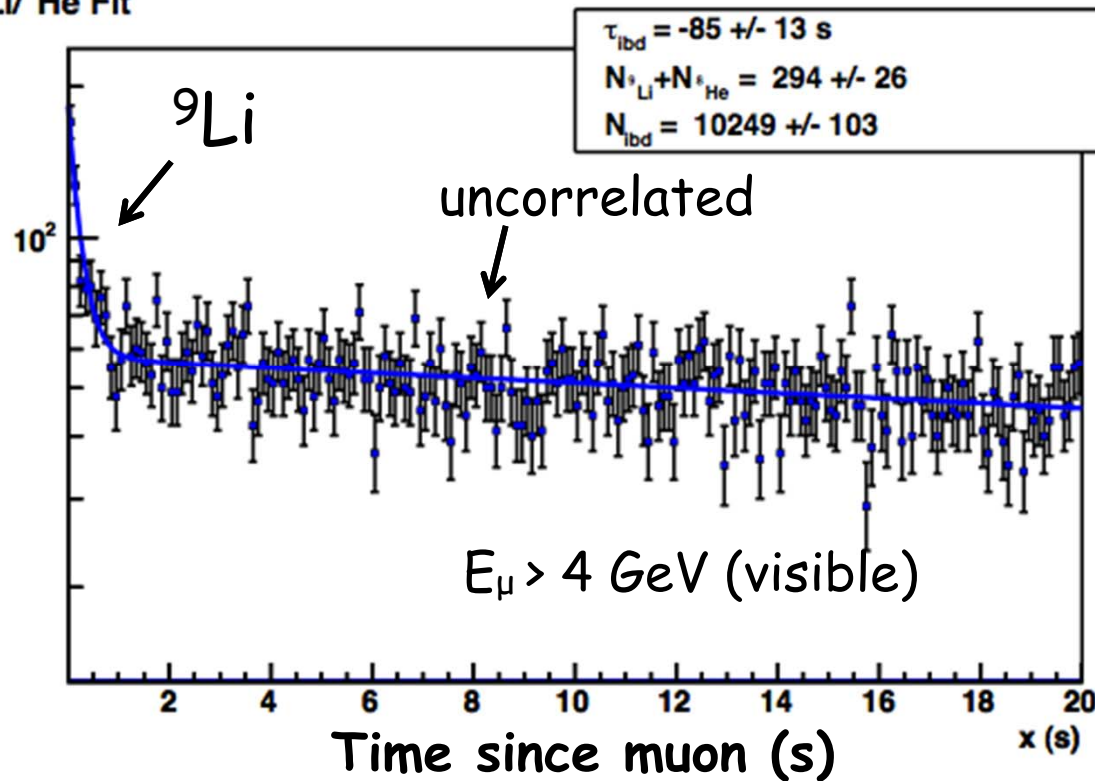
${}^9\text{Li}$: $\tau_{1/2} = 178$ ms, $Q = 13.6$ MeV
 ${}^8\text{He}$: $\tau_{1/2} = 119$ ms, $Q = 10.6$ MeV

Analysis muon-veto cuts
 control B/S to $\sim 0.4 \pm 0.2\%$.

- Generated by cosmic rays
- **Long-lived**
- Mimics antineutrino signal

${}^9\text{Li}/{}^8\text{He}$ Fit

Events / (0.1 s)



Fast Neutrons:

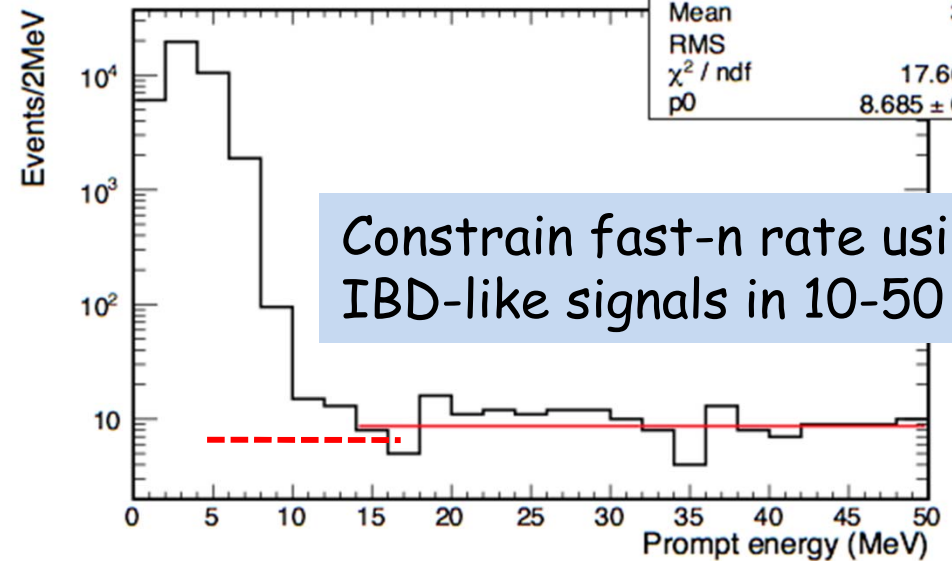
Energetic neutrons produced by cosmic rays (inside and outside of muon-veto system)

Mimics antineutrino (IBD) signal:

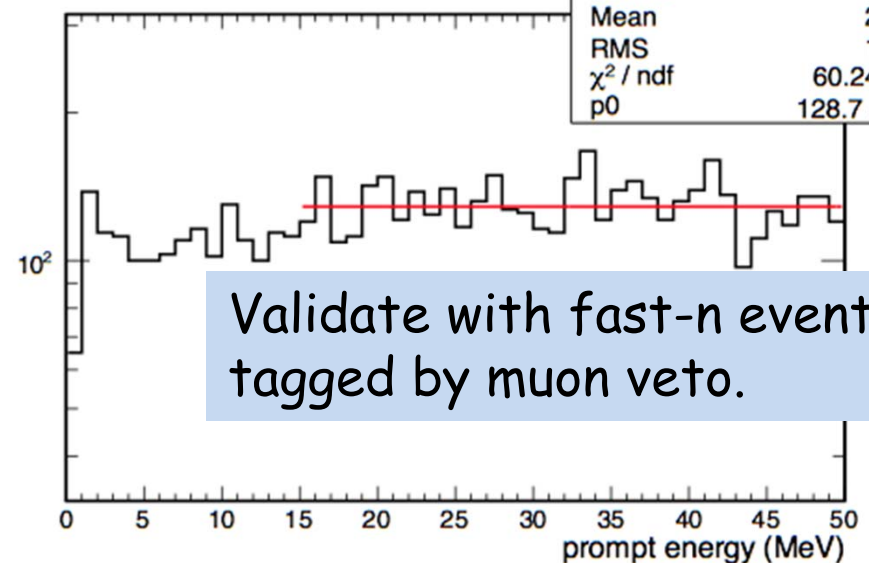
- Prompt: Neutron collides/stops in target
- Delayed: Neutron captures on Gd

Analysis muon-veto cuts control B/S to 0.06% (0.1%) of far (near) signal.

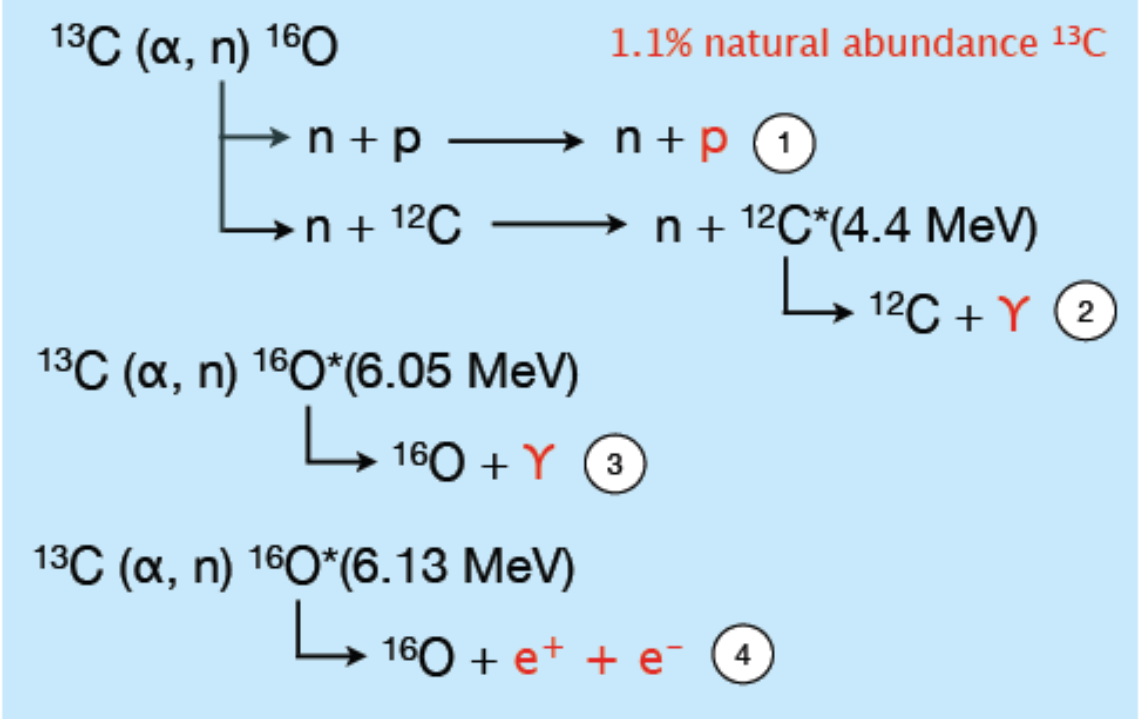
EH1 Prompt energy, AD#1



prompt energy of fast neutron candidate



Background: $^{13}\text{C}(\alpha, n)^{16}\text{O}$



Potential alpha sources:
 ^{238}U , ^{232}Th , ^{235}U , ^{210}Po :

Each of them are measured in-situ:

U & Th: cascading decay of
 Bi (or Rn) - Po - Pb

^{210}Po : spectrum fitting

Example alpha rate in AD1	^{238}U	^{232}Th	^{235}U	^{210}Po
Bq	0.05	1.2	1.4	10

Combining (α, n) cross-section, correlated background rate is determined.

Near Site: 0.04+-0.02 per day, B/S (0.006±0.004)%
Far Site: 0.03+-0.02 per day, B/S (0.04±0.02)%



	AD1	AD2	AD3	AD4	AD5	AD6
Antineutrino candidates	28935	28975	22466	3528	3436	3452
DAQ live time (day)	49.5530		49.4971	48.9473		
Veto time (day)	8.7418	8.9109	7.0389	0.8785	0.8800	0.8952
Efficiency	0.8019	0.7989	0.8363	0.9547	0.9543	0.9538
Accidentals (/day)	9.82 ± 0.06	9.88 ± 0.06	7.67 ± 0.05	3.29 ± 0.03	3.33 ± 0.03	3.12 ± 0.03
Fast neutron (/day)	0.84 ± 0.28	0.84 ± 0.28	0.74 ± 0.44	0.04 ± 0.04	0.04 ± 0.04	0.04 ± 0.04
$^8\text{He}/^9\text{Li}$ (/day)	3.1 ± 1.6		1.8 ± 1.1	0.16 ± 0.11		
Am-C corr. (/day)	0.2 ± 0.2					
$^{13}\text{C}(\alpha, n)^{16}\text{O}$ (/day)	0.04 ± 0.02	0.04 ± 0.02	0.035 ± 0.02	0.03 ± 0.02	0.03 ± 0.02	0.03 ± 0.02
Antineutrino rate (/day)	714.17 ± 4.58	717.86 ± 4.60	532.29 ± 3.82	71.78 ± 1.29	69.80 ± 1.28	70.39 ± 1.28



Antineutrino flux is estimated for each reactor core

Flux estimated using:

$$S(E_\nu) = \frac{W_{th}}{\sum_i (f_i/F) e_i} \sum_i^{istopes} (f_i/F) S_i(E_\nu)$$

Reactor operators provide:

- Thermal power data: W_{th}
- Relative isotope fission fractions: f_i

Energy released per fission: e_i

V. Kopekin et al., Phys. Atom. Nucl. 67, 1892 (2004)

Antineutrino spectra per fission: $S_i(E_\nu)$

K. Schreckenbach et al., Phys. Lett. B160, 325 (1985)

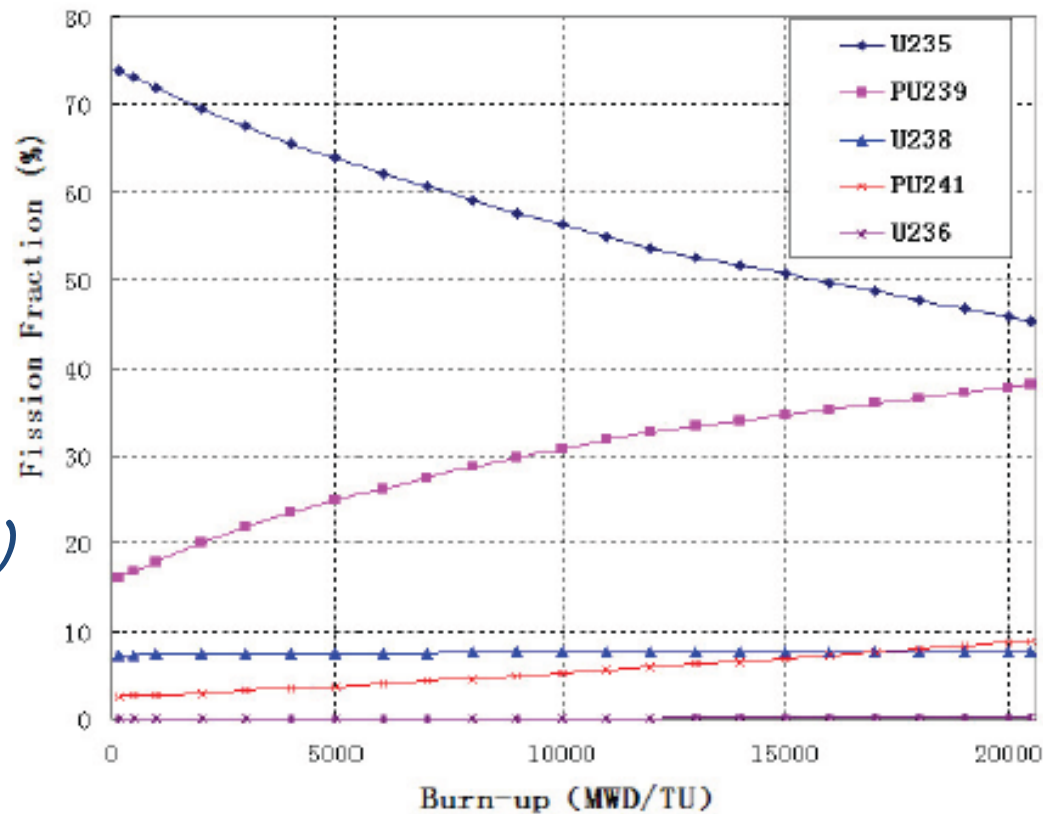
A. A. Hahn et al., Phys. Lett. B218, 365 (1989)

P. Vogel et al., Phys. Rev. C24, 1543 (1981)

T. Mueller et al., Phys. Rev. C83, 054615 (2011)

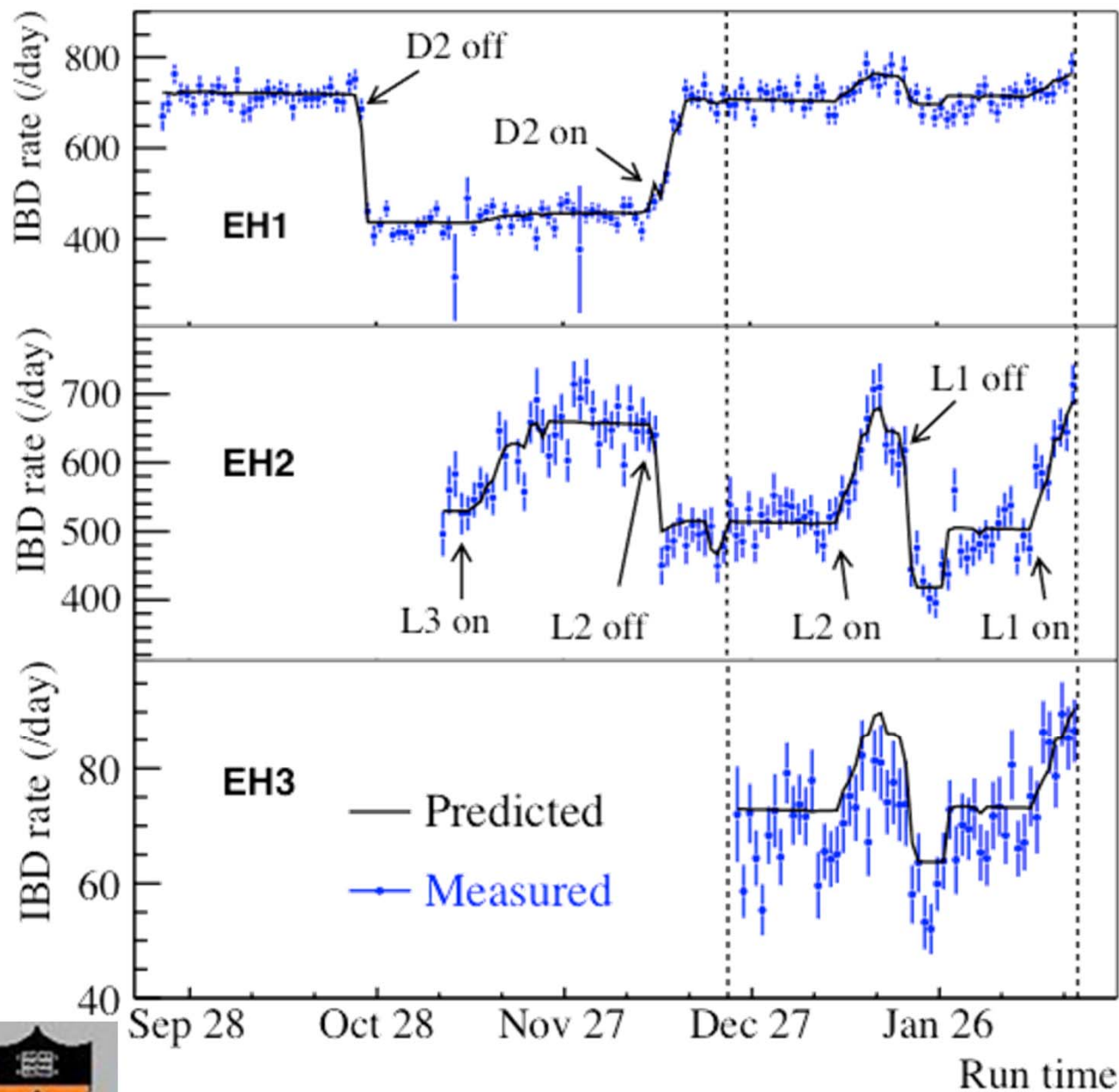
P. Huber, Phys. Rev. C84, 024617 (2011)

Isotope fission rates vs. reactor burnup



Flux model has negligible impact on far vs. near oscillation measurement





Detected rates strongly correlated with reactor flux expectations.

- Predicted Rates:** (in figure)
- Assumes no oscillation.
 - Normalization is determined by fit to data.
 - Absolute normalization is within a few percent of expectations.



Uncertainty Summary

Detector		
	Efficiency	Correlated
Target Protons		0.47%
Flasher cut	99.98%	0.01%
Delayed energy cut	90.9%	0.6%
Prompt energy cut	99.88%	0.10%
Multiplicity cut		0.02%
Capture time cut	98.6%	0.12%
Gd capture ratio	83.8%	0.8%
Spill-in	105.0%	1.5%
Livetime	100.0%	0.002%
Combined	78.8%	1.9%

For near/far rate analysis , only uncorrelated uncertainties are relevant.

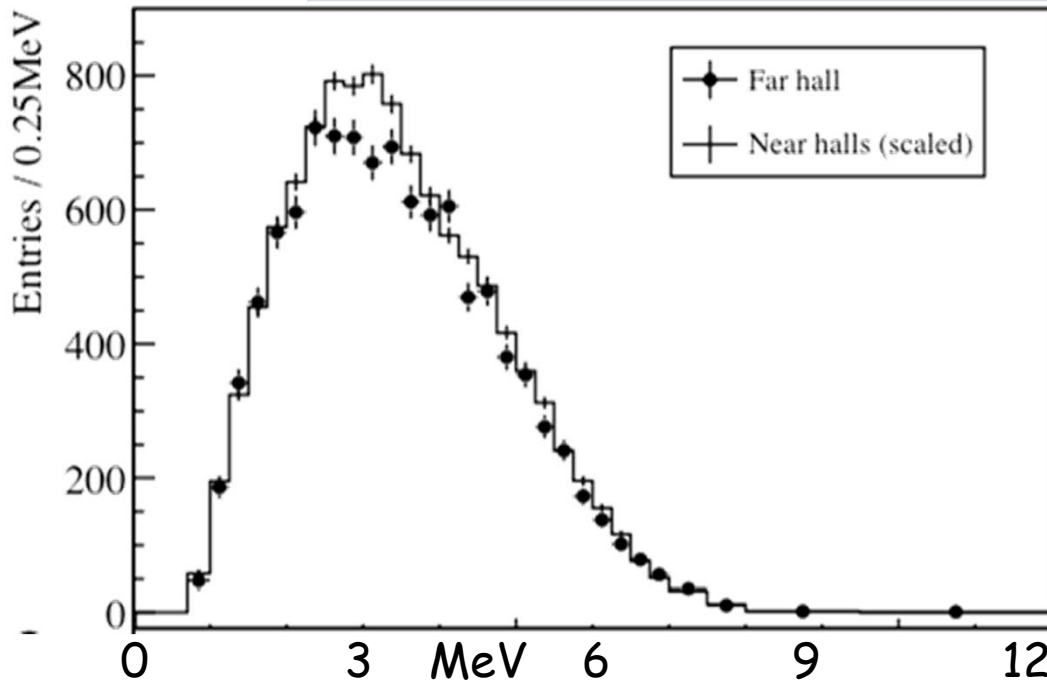
Largest systematics are smaller than far site statistics (~1%)

Reactor			
	Correlated		Uncorrelated
Energy/fission	0.2%	Power	0.5%
$\bar{\nu}_e$ /fission	3%	Fission fraction	0.6%
		Spent fuel	0.3%
Combined	3%	Combined	0.8%

Influence of uncorrelated reactor systematics reduced (~1/20) by far vs. near measurement.



Compare the far/near measured rates (and spectra)



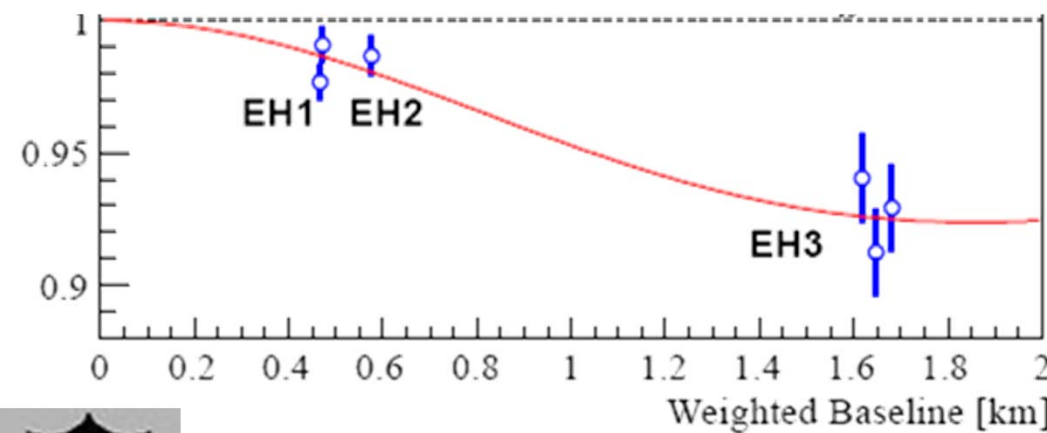
$$R = \frac{Far_{measured}}{Far_{expected}} = \frac{M_4 + M_5 + M_6}{\sum_{i=4}^6 (\alpha_i(M_1 + M_2) + \beta_i M_3)}$$

M_n is the measured IBD rate in detector n.

Weights α_i, β_i are determined from baselines and reactor fluxes.

$\alpha \sim 0.014, \beta \sim 0.10$

\Rightarrow EH2 weighted 7 times EH1



$$R = 0.940 \pm 0.011 \text{ (stat)} \pm 0.004 \text{ (syst)}$$

Clear observation of far-site deficit.

Spectral distortion consistent with oscillation.*

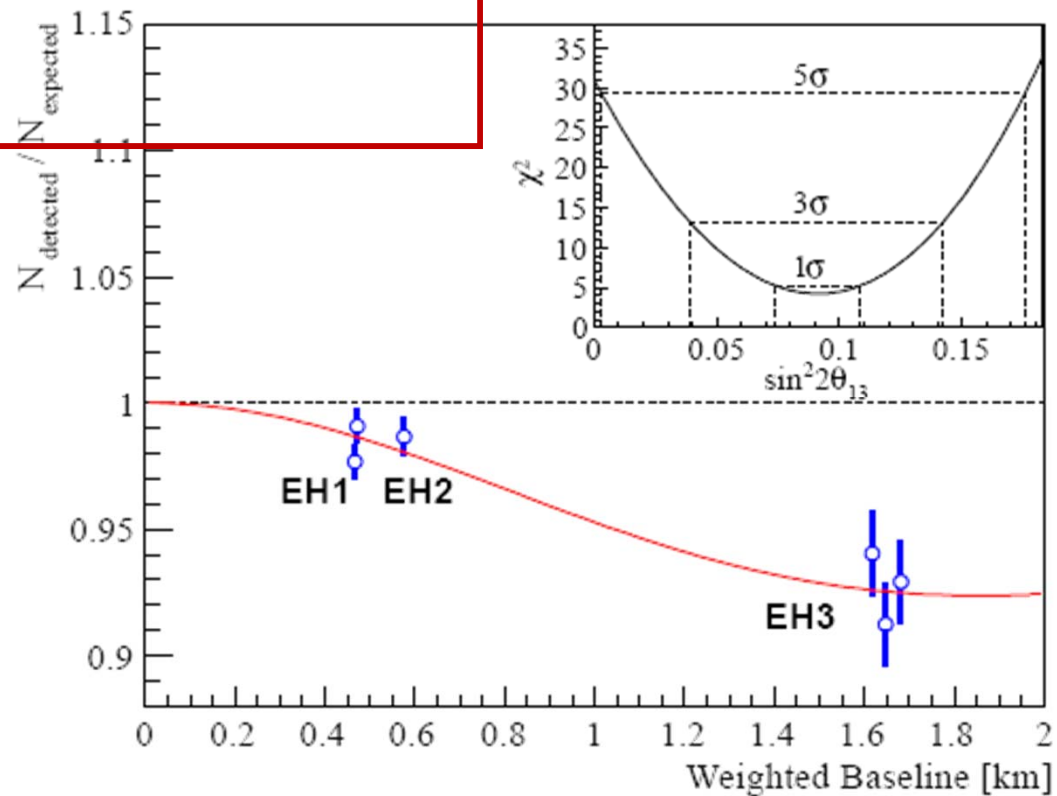
- Caveat: Spectral systematics not fully studied; θ_{13} value from shape analysis is not recommended.



$$\chi^2 = \min_{\gamma} \sum_{d=1}^6 \frac{[M_d - T_d (1 + \sum_r \omega_r^d \epsilon_r + \epsilon + \epsilon_d) - B_d (1 + \eta_d)]^2}{M_d} + \sum_r \frac{\epsilon_r^2}{\sigma_r^2} + \sum_{d=1}^6 \left[\frac{\epsilon_d^2}{\sigma_d^2} + \left(\frac{\eta_d}{\sigma_B^d} \right)^2 \right] + \cancel{\frac{\epsilon^2}{\sigma^2}}$$

"Pull" χ^2 analysis.
 No constraint on absolute normalization.
 Fit on the near-far relative measurement.

$\text{Sin}^2 2\theta_{13} = 0.092 \pm 0.016(\text{stat}) \pm 0.005(\text{syst})$
 $\chi^2/\text{NDF} = 4.26/4$
 5.2 σ for non-zero θ_{13}



- The Daya Bay reactor neutrino experiment has made an unambiguous observation of reactor electron-antineutrino disappearance at ~ 2 km for $E_\nu \sim 3$ MeV:

$$R_{\text{far/near}} = 0.940 \pm 0.011 \text{ (stat)} \pm 0.004 \text{ (syst)}$$

- Interpretation of this disappearance as neutrino oscillation in a 3-neutrino context yields, via a χ^2 analysis:

$$\sin^2 2\theta_{13} = 0.092 \pm 0.016 \text{ (stat)} \pm 0.005 \text{ (syst)}$$

ruling out zero at 5.2 standard deviations.

$$\theta_{13} = 8.8^\circ \text{ (or } 81.2^\circ\text{?)}, \sin^2 \theta_{13} = 0.024$$

- Installation of final pair of antineutrino detectors this Summer



In the 3- ν context, taking $\Delta m_{13}^2 \sim \Delta m_{23}^2$, the ν_e survival probability is

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) = 1 - \sin^2 2\theta_{13} \sin^2 \left(\Delta m_{13}^2 L / 4E_\nu \right) - \cos^4 \theta_{13} \sin^2 2\theta_{12} \sin^2 \left(\Delta m_{12}^2 L / 4E_\nu \right)$$

For large L , the 2nd term averages to $(1/2) \sin^2 2\theta_{13} \sim 0.04$, such that

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) \approx 0.96 - \cos^4 \theta_{13} \sin^2 2\theta_{12} \sin^2 \left(\Delta m_{12}^2 L / 4E_\nu \right)$$

If $\theta_{13} \sim 81^\circ$, then no oscillations could be seen in long baseline experiments with reactor or solar neutrinos.

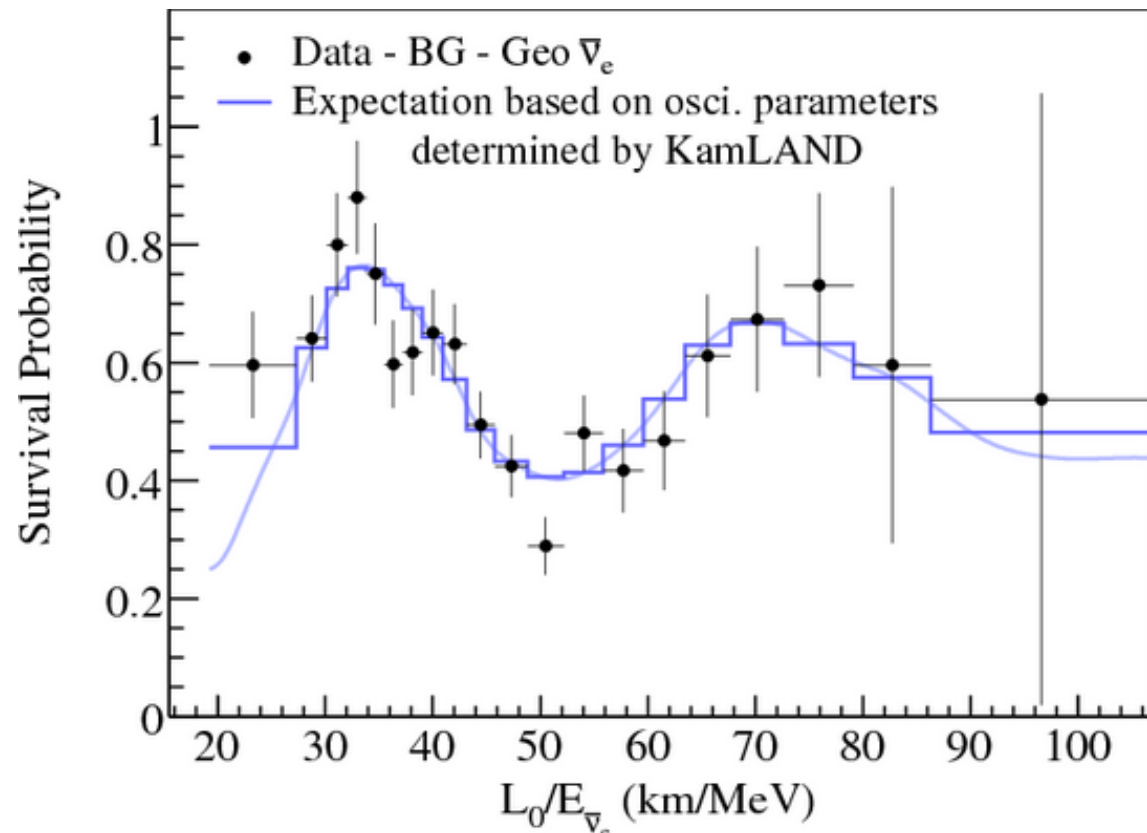
The clear evidence for oscillations in the KamLAND experiment excludes that θ_{13} is near 90° .

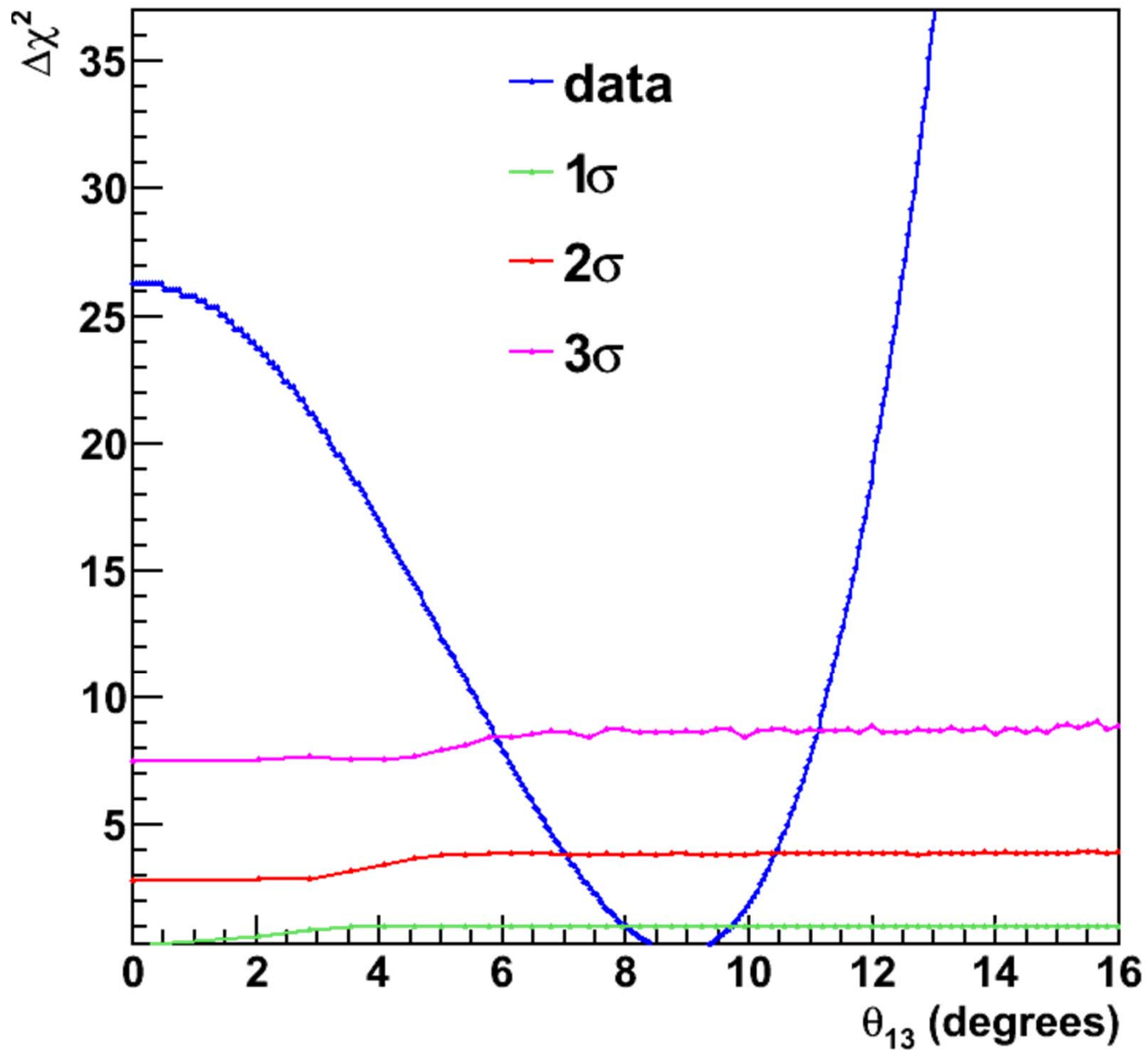
Phys. Rev. Lett. **100**, 221803 (2008)

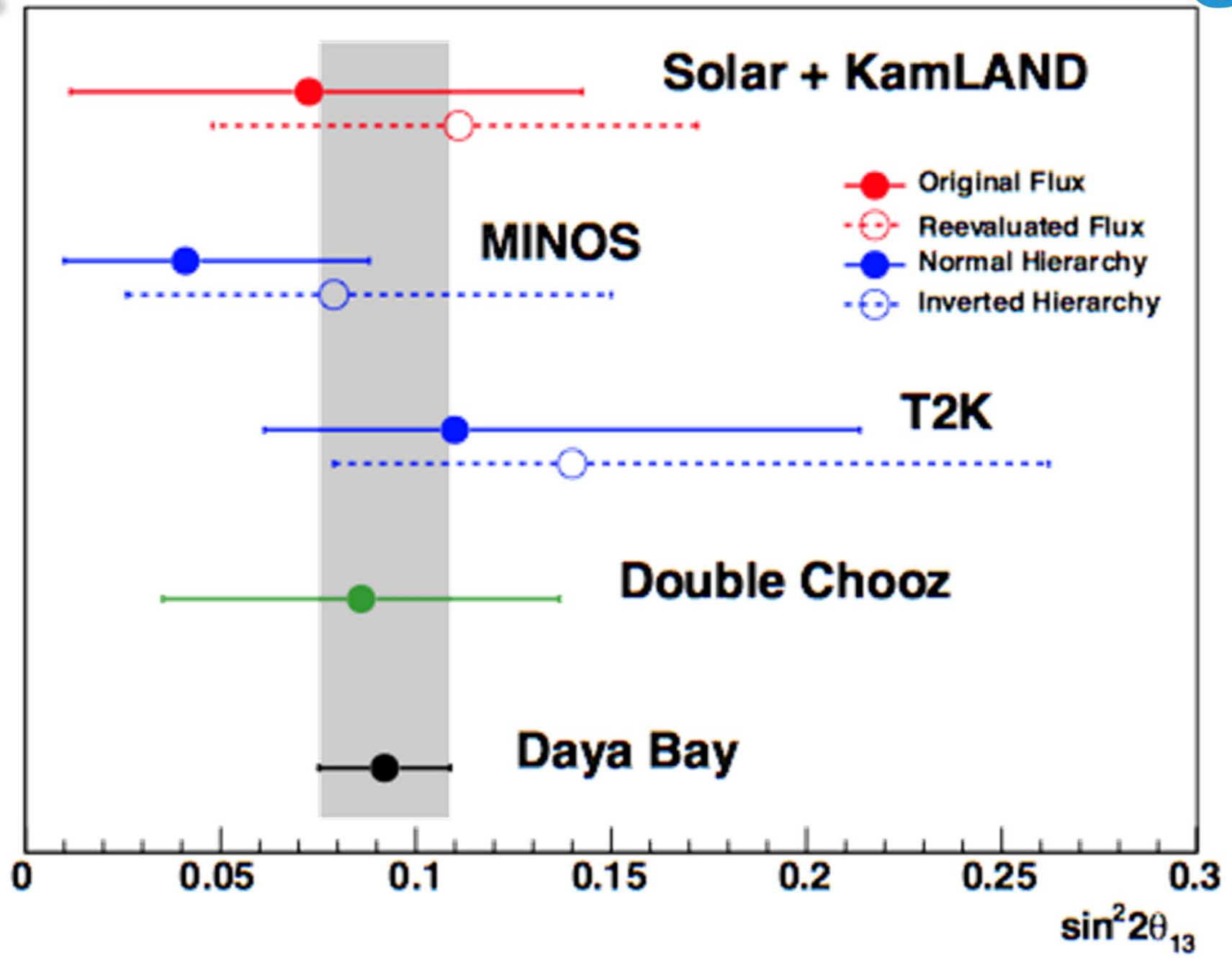
Using our value for $\theta_{13} = 8.8^\circ$,

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) \approx 0.96 - 0.95 \sin^2 2\theta_{12} \sin^2 \left(\Delta m_{12}^2 L / 4E_\nu \right)$$

\Rightarrow 5% increase in $\sin^2 \theta_{12}$ compared to previous fits.





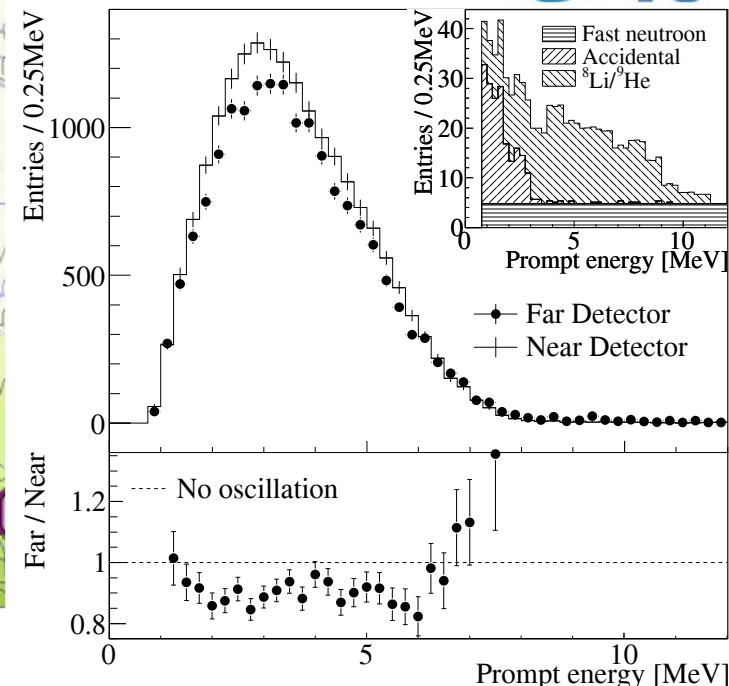
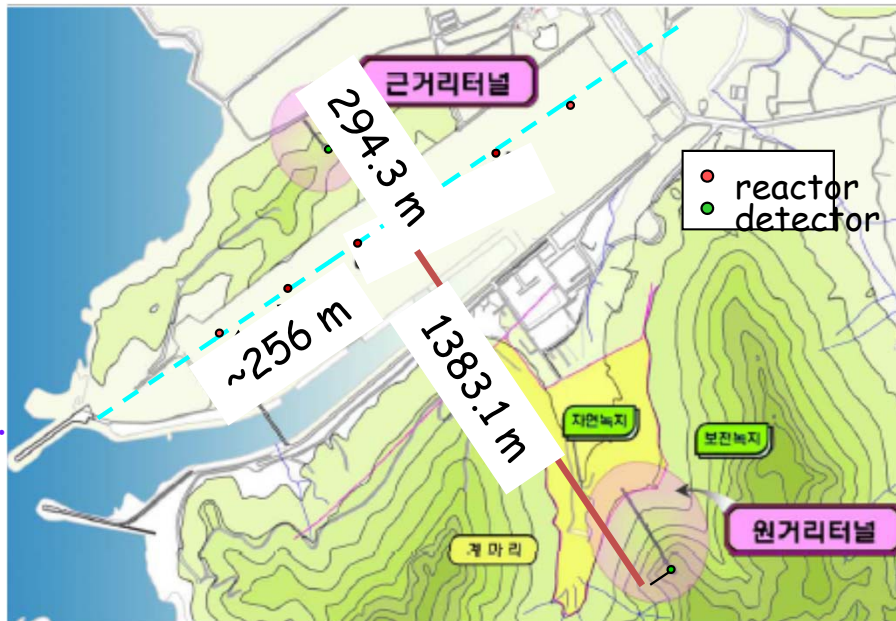


Yonggwang Nuclear Power Plant

Six ~ 3 GW_e reactors along a line.

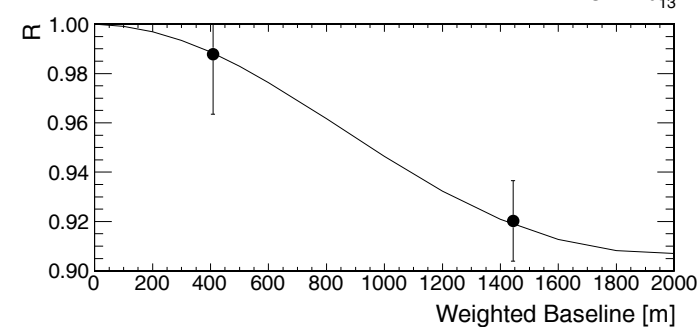
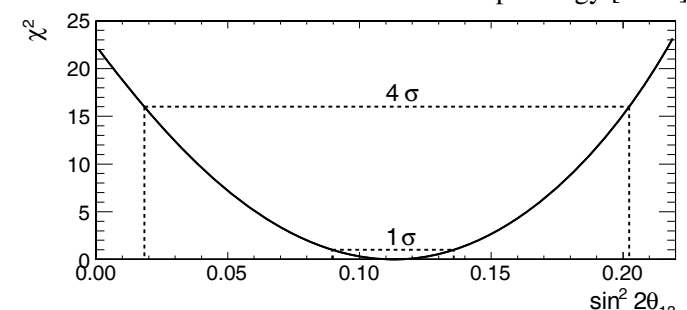
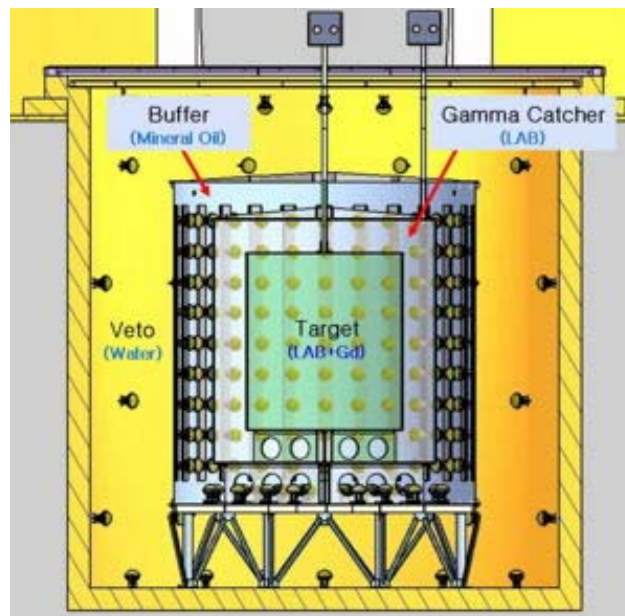
1 Near and 1 Far detector, each 16 tons.

229 days of data since August 2011.

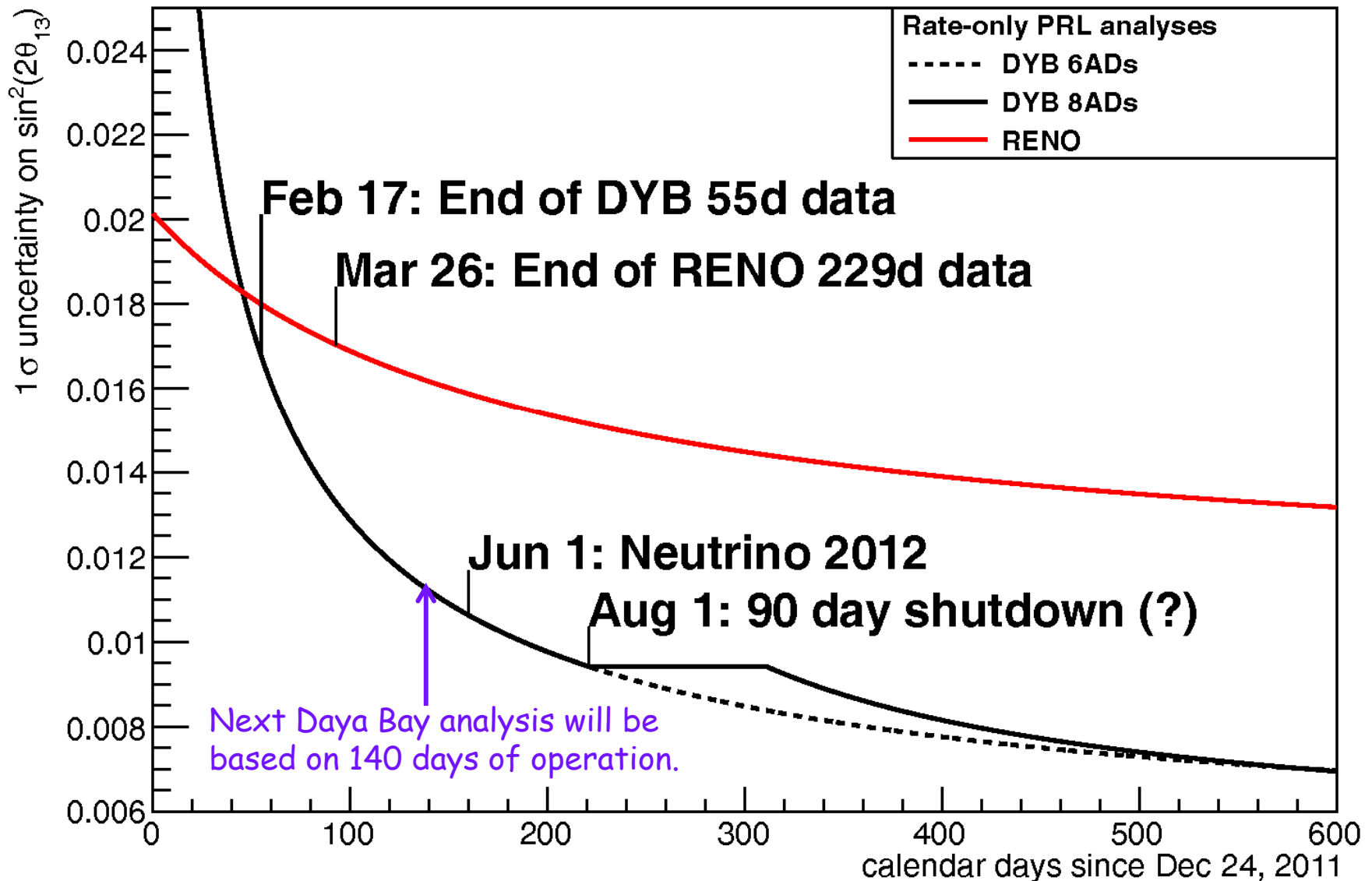


$$\sin^2 2\theta_{13} = 0.113 \pm 0.013 \text{ (stat)} \pm 0.019 \text{ (syst.)}$$

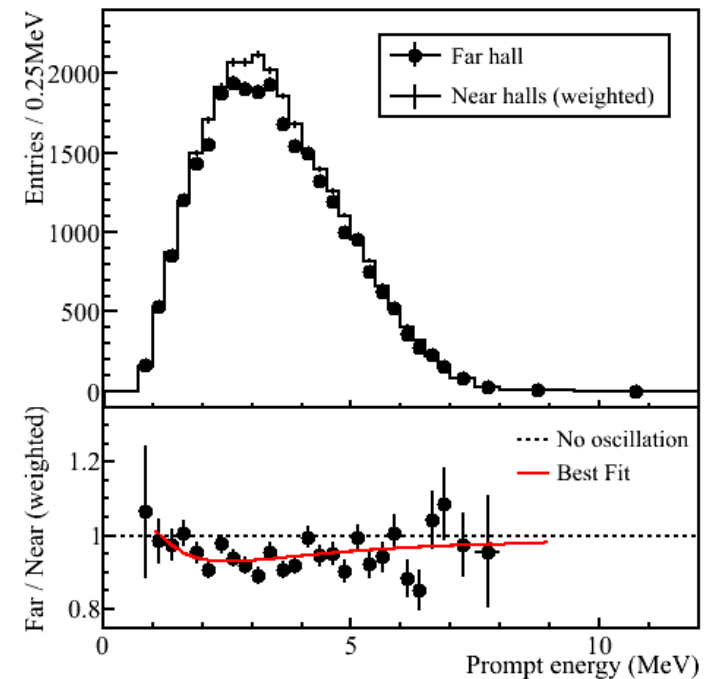
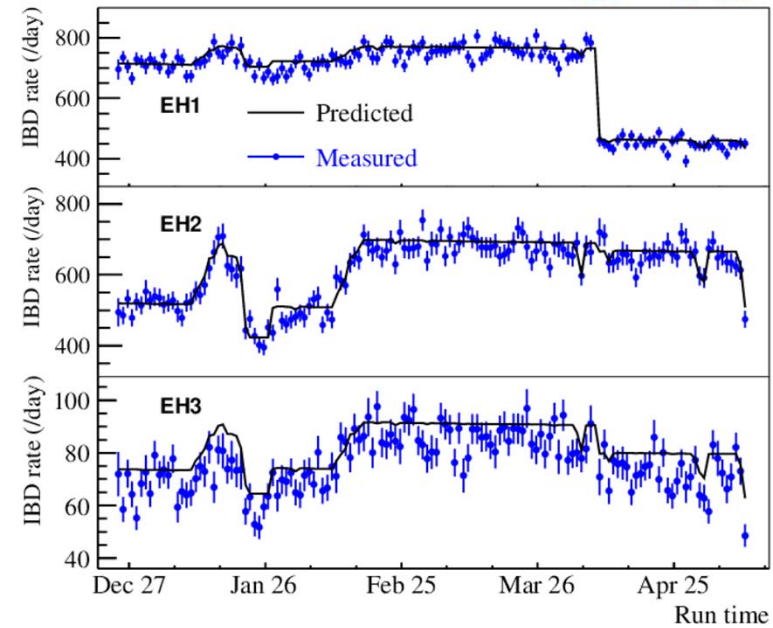
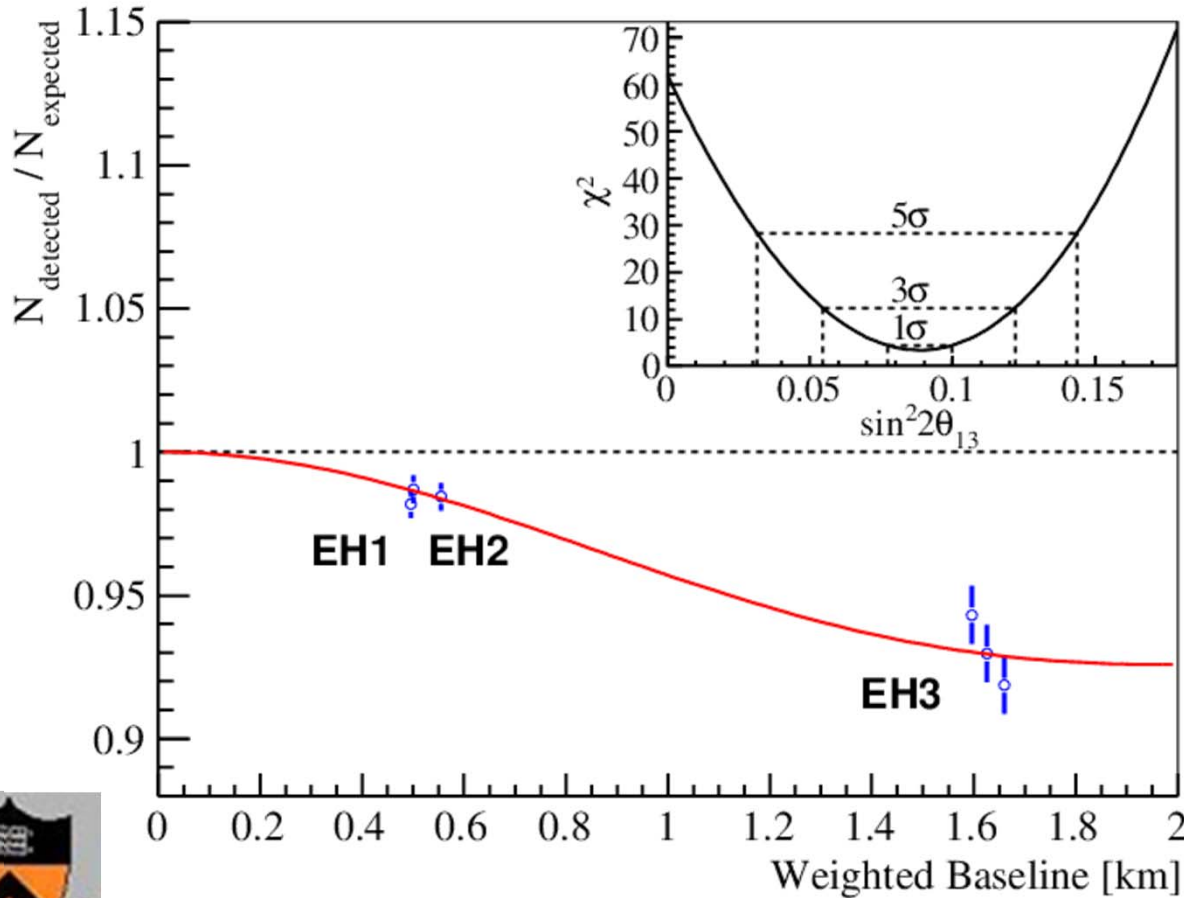
<http://arxiv.org/abs/1204.0626>
(Apr. 8, 2012)

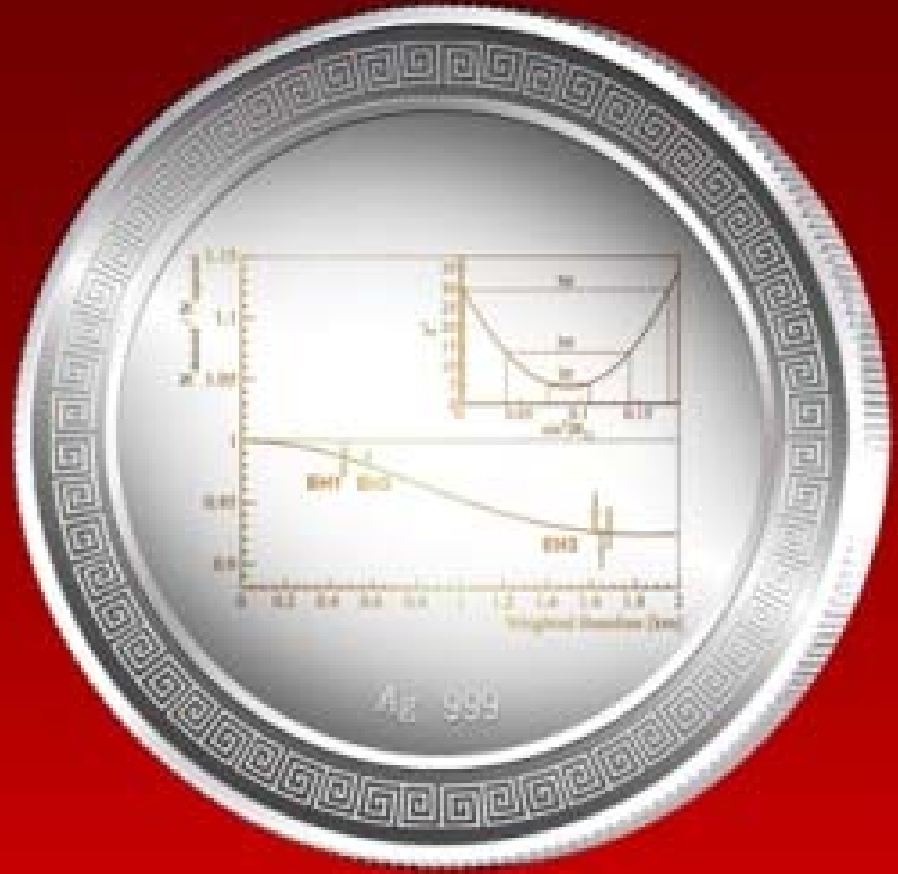


1σ Uncertainty vs. time assuming \sqrt{N} statistics



- χ^2 fit result: $\sin^2 2\theta_{13} = 0.089 \pm 0.010(\text{stat}) \pm 0.005(\text{syst})$
- $\chi^2_{\min} = 3.4$
- 7.7σ non-zero





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