## The Fields in a Box with Resistive Walls

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## 1 Problem

A uniform electric field is desired throughout a cube of edge a. This can be arranged by constructing a cubical box in which the face at z = 0 is a (perfect) conductor at potential 0, the face at z = a is a conductor at potential  $V_0$ , and the four remaining faces are made of a poor conductor, say of resistivity  $\rho_0$  Ohms/square.

Suppose, however, the material of the faces between z = 0 and a has a nonuniform resistivity which varies as,

$$\rho(z) = \rho_0 \left( 1 + \epsilon \sin \frac{2\pi z}{a} \right),$$

where  $\epsilon \ll 1$ .

An exact solution for the potential inside the box can be given, but is very cumbersome.

Calculate the potential V(z) on the resistive walls, and **estimate** the potential at the center of the cube.

Estimate the maximum ratio of the transverse electric field to the longitudinal field  $(E_z)$ . Where does this maximum occur?

You may assume the current in the resistive faces flows parallel to the z axis.

## 2 Solution

We first find the current in the resistive walls, and then calculate the potential as a function of z on the walls. To estimate the potential  $\phi$  as the center of the cube, we use the fact that  $\nabla^2 \phi = 0$  implies that the potential in the center of a small volume element is the average of the potential over the surrounding surface.

First, the wall current  $I = V_0/R$ .

The total resistance  $R = \int_0^a R(z) dz$ , where  $R(z) = \rho(z) dz/4a$ , since the perimeter of the wall is 4a. Hence,

$$R = \frac{\rho_0}{4a} \int_0^a \left(1 + \epsilon \sin \frac{2\pi z}{a}\right) dz = \frac{\rho_0}{4a} \left[z - \frac{\epsilon a}{2\pi} \cos \frac{2\pi z}{a}\right]_0^a = \frac{\rho_0}{4}.$$

Then on the wall,

$$V(z) = \int_0^z IdR = \frac{V_0}{R} \frac{\rho_0}{4a} \int_0^z \left(1 + \epsilon \sin\frac{2\pi z}{a}\right) dz = V_0 \left[\frac{z}{a} + \frac{\epsilon}{2\pi} \left(1 - \cos\frac{2\pi z}{a}\right)\right]$$

The average potential on the wall is,

$$\langle \phi \rangle_{\text{wall}} = \frac{V_0}{2} \left( 1 + \frac{\epsilon}{\pi} \right).$$

The average potential over the entire surface of the cube is,

$$\langle \phi \rangle = \frac{1}{6} \left[ 0 + V_0 + 4 \frac{V_0}{2} \left( 1 + \frac{\epsilon}{\pi} \right) \right] = \frac{V_0}{2} \left( 1 + \frac{2\epsilon}{3\pi} \right) \approx \phi_{\text{center}}.$$

The electric field in the z direction is  $E_z \approx V_0/a$ .

The transverse electric field is greatest in the midplane, z = a/2, where the wall potential is,

$$\phi_{\text{wall}} = \frac{V_0}{2} \left( 1 + \frac{2\epsilon}{\pi} \right).$$

The electric field in the x direction is roughly,

$$E_x \approx \frac{\phi_{\text{wall}} - \phi_{\text{center}}}{a/2} = \frac{4\epsilon V_0}{3\pi a} \approx \frac{4\epsilon}{3\pi} E_z.$$

One might argue that the variation of the potential with x is sine-like, and so the maximum slope is  $\pi/2$  times the simplified estimate. Then,

$$E_{x,\max} \approx \frac{2\epsilon}{3} E_z.$$

Similarly for  $E_y$ .

The transverse field is therefore greatest in the corners of the wall, at the midplane in z, and,

$$\frac{E_{\perp,\max}}{E_z} \approx \frac{2\sqrt{2}}{3}\epsilon \approx \epsilon.$$

No doubt,  $E_{\perp}/E_z \approx \epsilon$  could have been guessed at once.